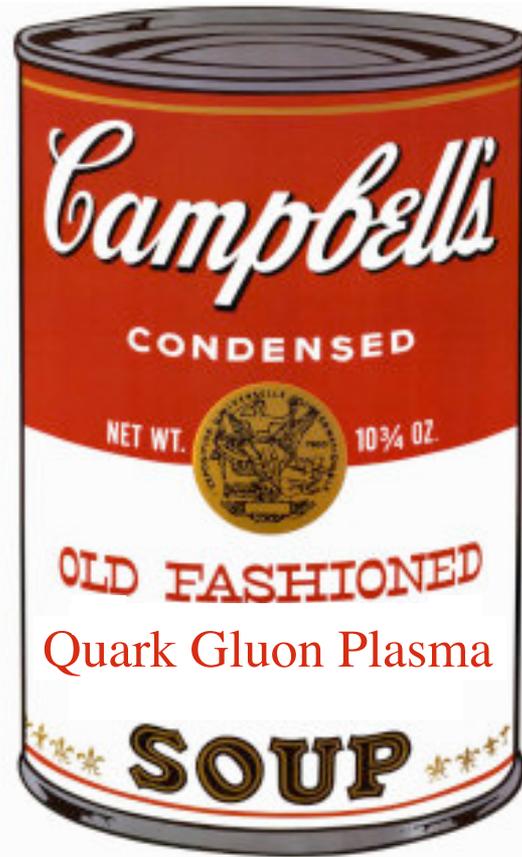


Creating Quark Soup



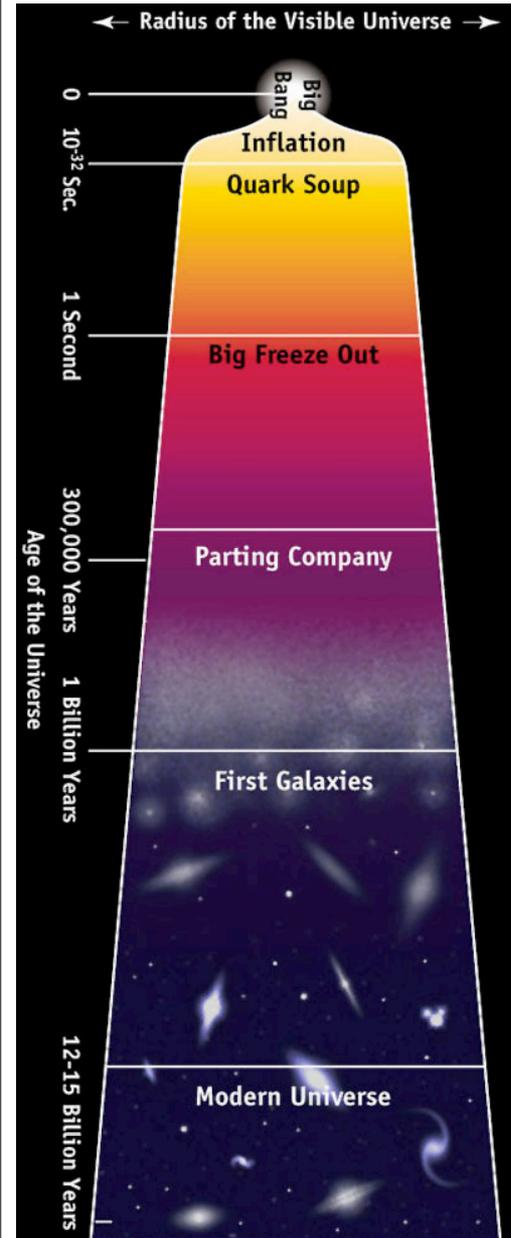
Andy Warhol

Helen Caines - Yale University

April 13th 2009

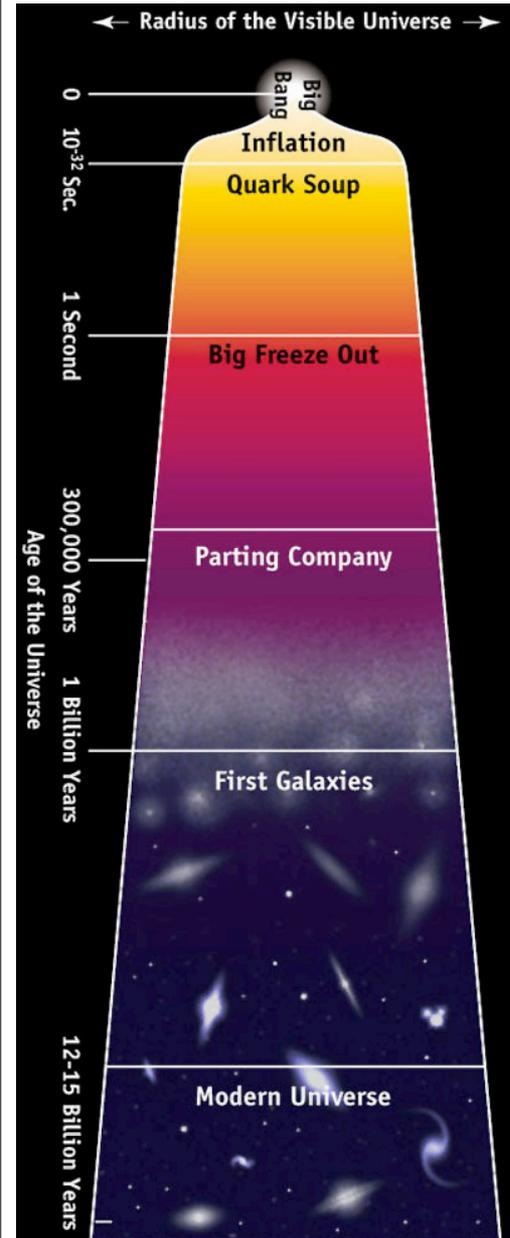
Franklin & Marshall College

Evolution of the universe



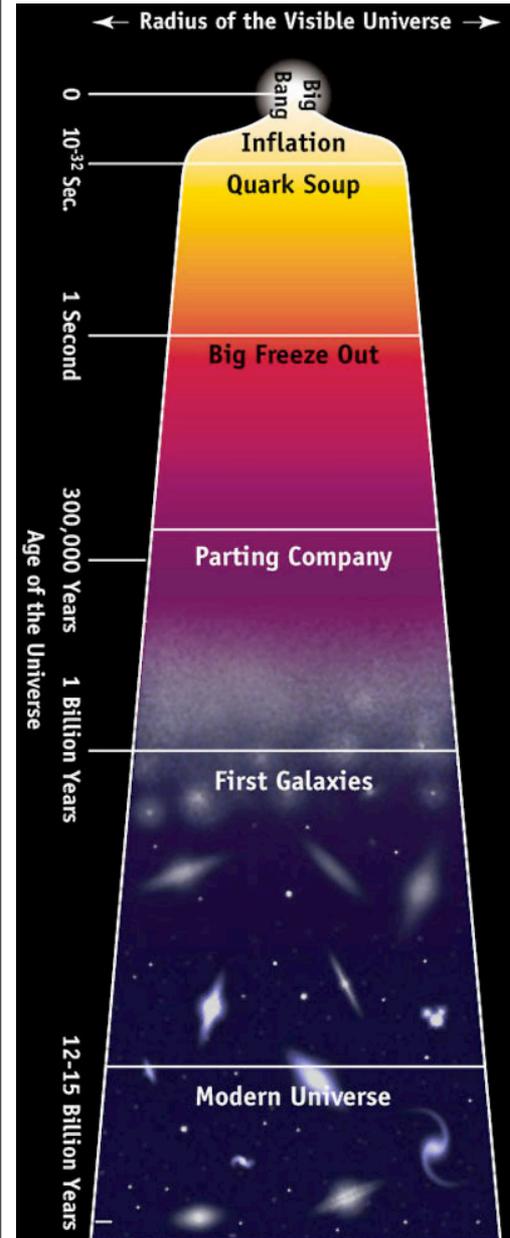
10^{-44} sec	Quantum Gravity	Unification of all 4 forces	10^{32} K
10^{-35} sec	Grand Unification	E-M/Weak = Strong forces	10^{27} K
10^{-35} sec ?	Inflation	universe exponentially expands by 10^{26}	10^{27} K
$2 \cdot 10^{-10}$ sec	Electroweak unification	E-M = weak force	10^{15} K
$2 \cdot 10^{-6}$ sec	Proton-Antiproton pairs	creation of nucleons	10^{13} K
6 sec	Electron-Positron pairs	creation of electrons	6×10^9 K
3 min	Nucleosynthesis	light elements formed	10^9 K
10^6 yrs	Microwave Background	recombination - transparent to photons	3000 K
10^9 yrs ?	Galaxy formation	bulges and halos of normal galaxies form	20 K

Evolution of the universe



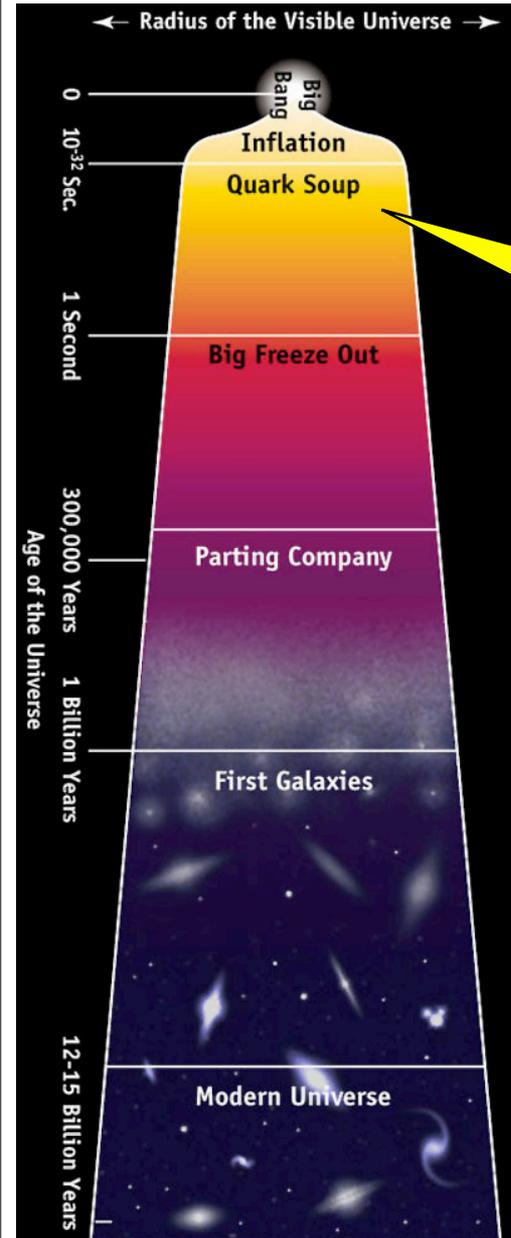
The universe gets cooler !

Evolution of the universe



Reheating Matter ?

Evolution of the universe

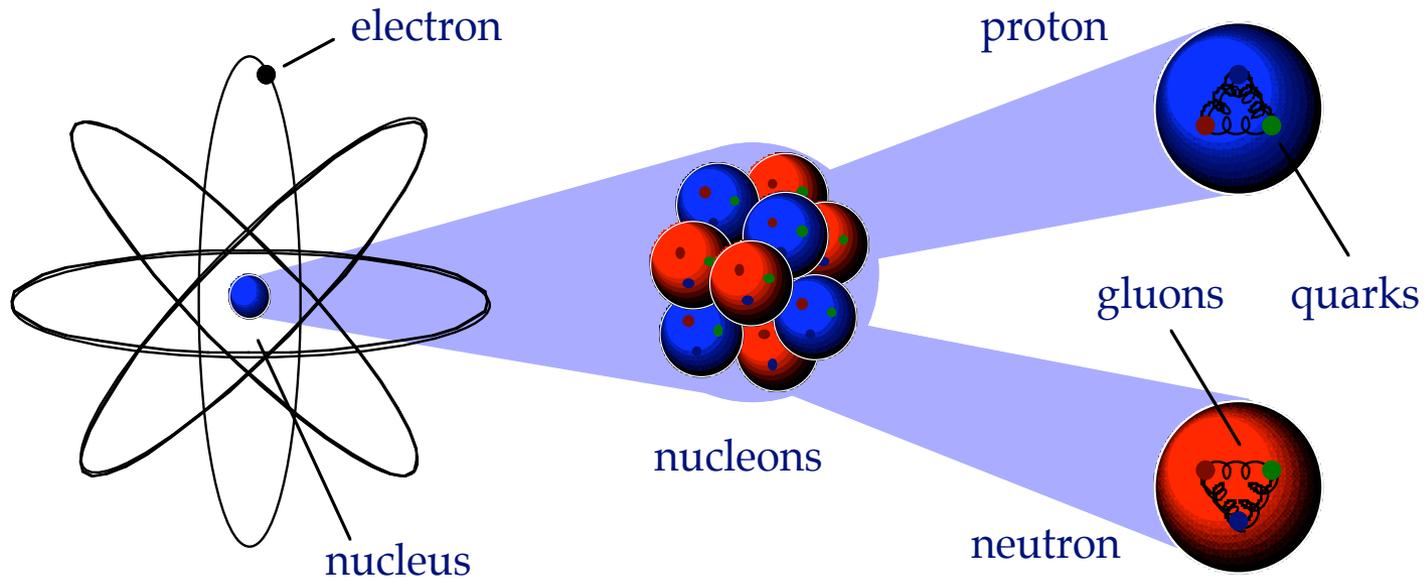


Reheating Matter ?

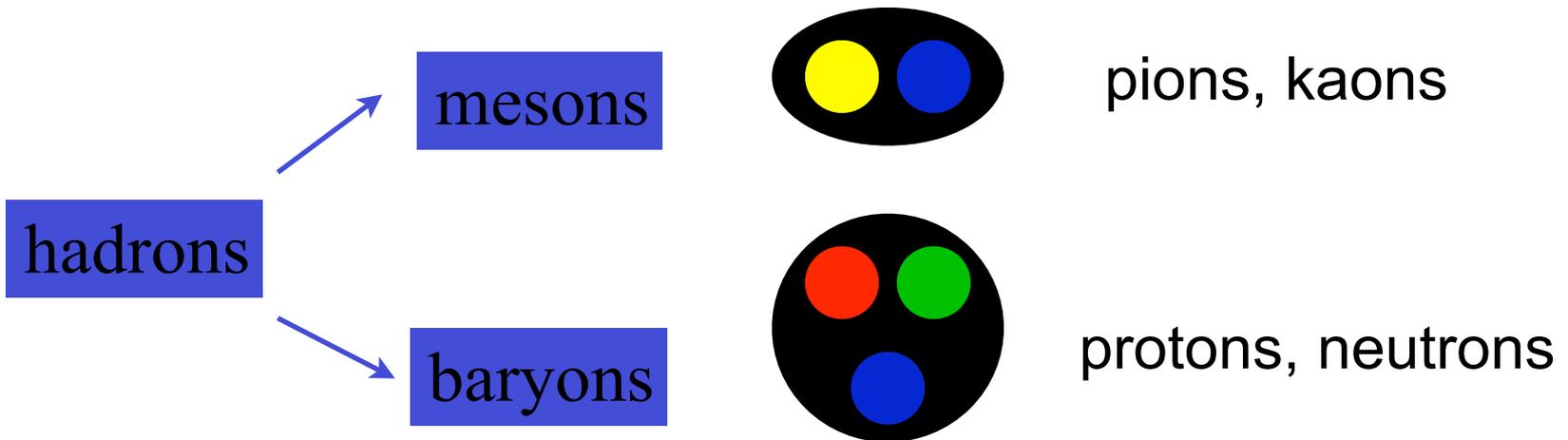
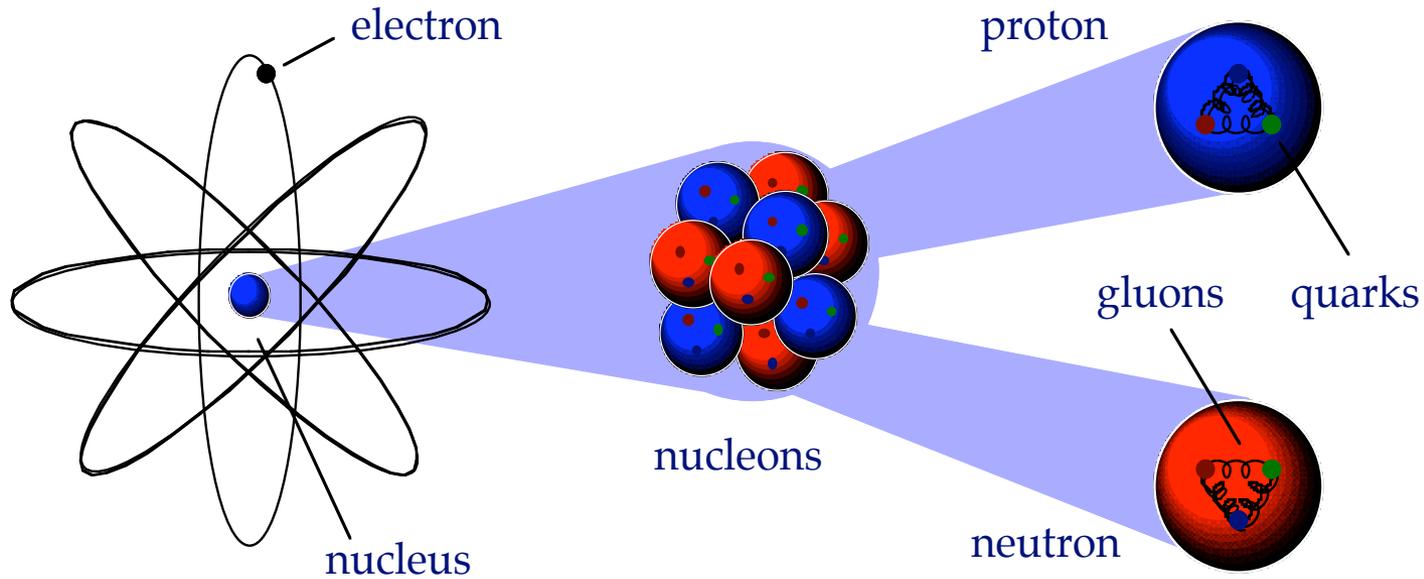
?

Need temperatures
around
 $1.5 \cdot 10^{12}$ K
(200 MeV)
far hotter than center of
the sun ($\sim 2 \cdot 10^7$ K)

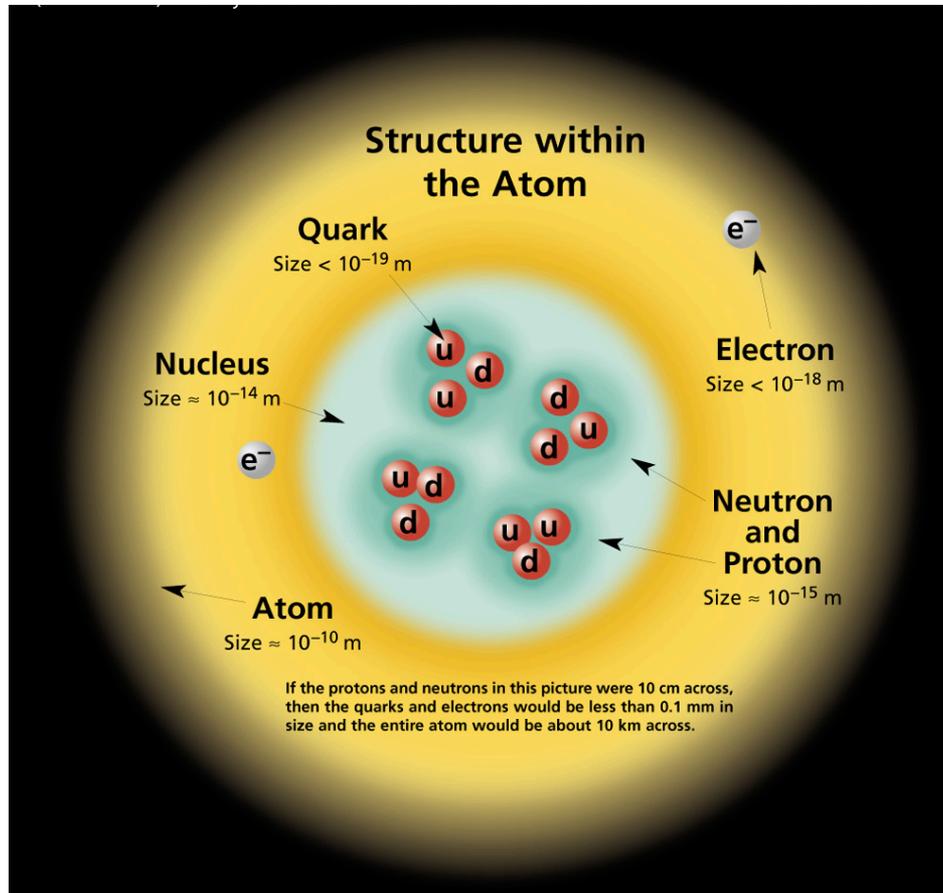
Essential ingredients of matter



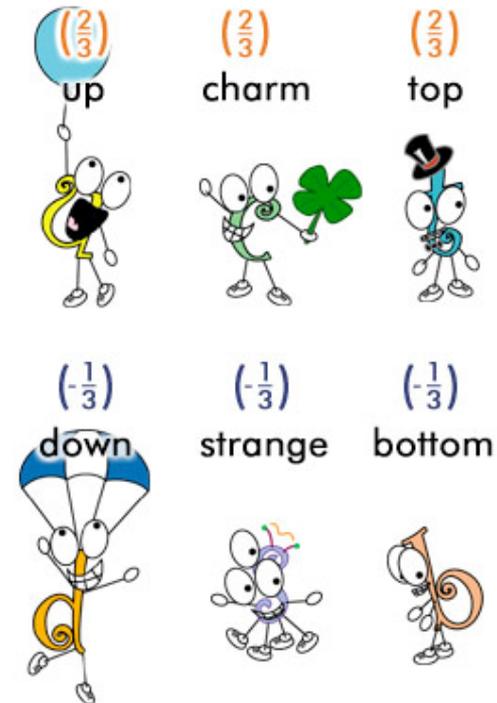
Essential ingredients of matter



More about partons

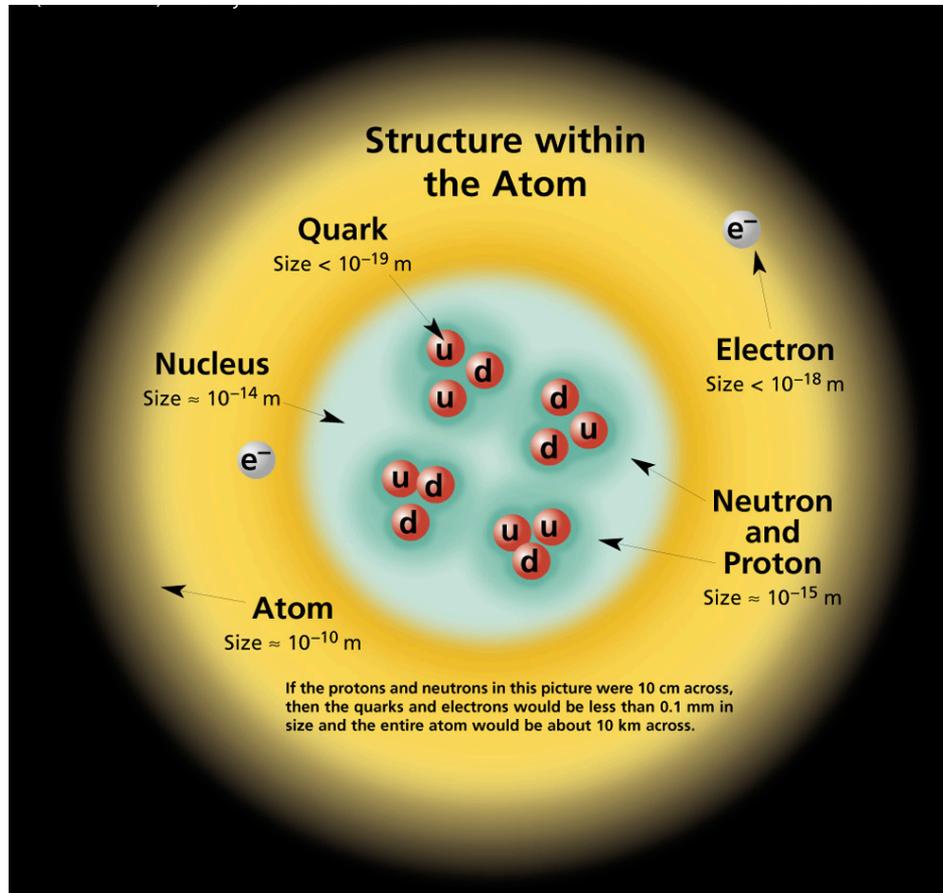


Ordinary matter made of **up** and **down** quarks

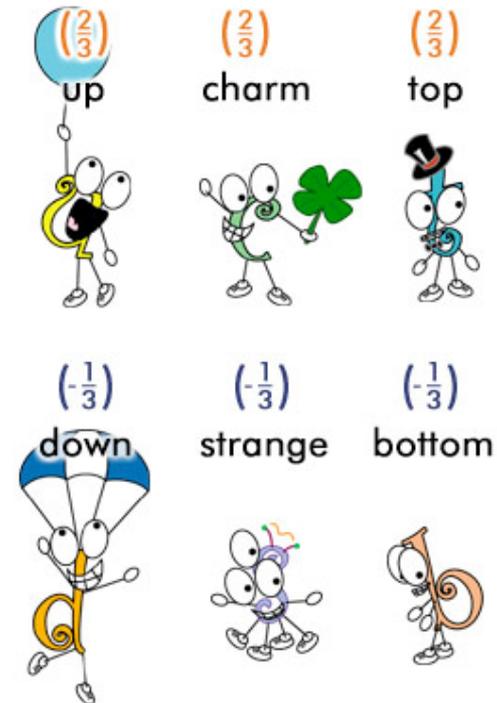


- Quarks interact by exchanging gluons
 - Nucleons are held together by gluons

More about partons



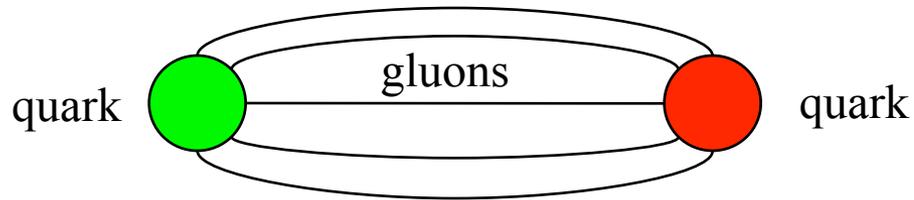
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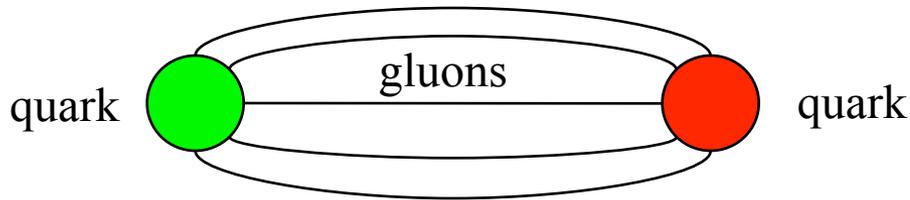
Free quarks have never been seen - distinctive non-integer charge

Why we don't see free quarks



The size of a nucleus is $1.2A^{1/3}$ fm where A is the mass number and a fm is 10^{-15} m

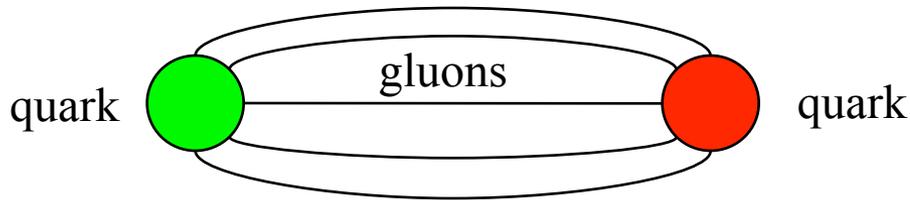
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$$1 \frac{GeV}{fm} = \frac{10^9 eV}{10^{-15} m} \times \frac{1.6 \times 10^{-19} J}{eV} = 1.6 \times 10^5 N$$

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Compare to gravitational force at Earth's surface

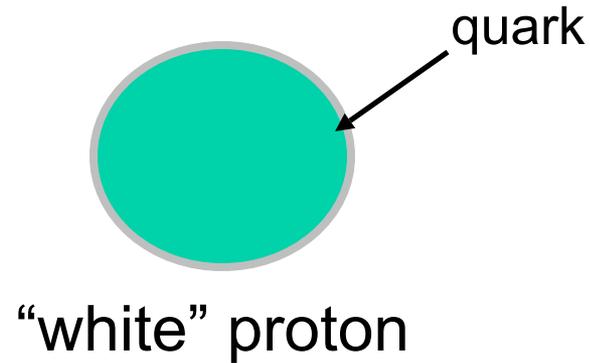
$$F = 1.6 \times 10^5 N = M \times g = M \times 9.8 m/s^2$$

 $M = 16,300 kg$

Quarks exert 16 metric tons of force on each other!

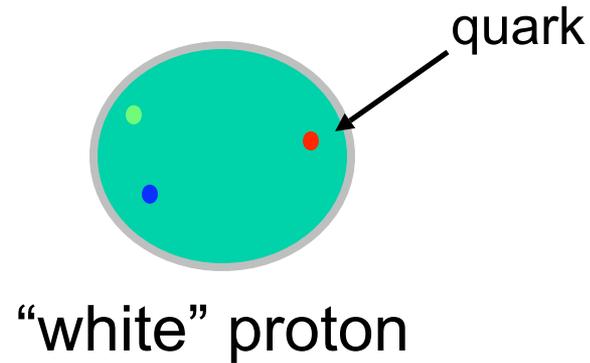
Confinement - QCD

Confinement: fundamental & crucial (but *not* understood!) feature of strong force - colored objects (quarks) have ∞ energy in normal vacuum



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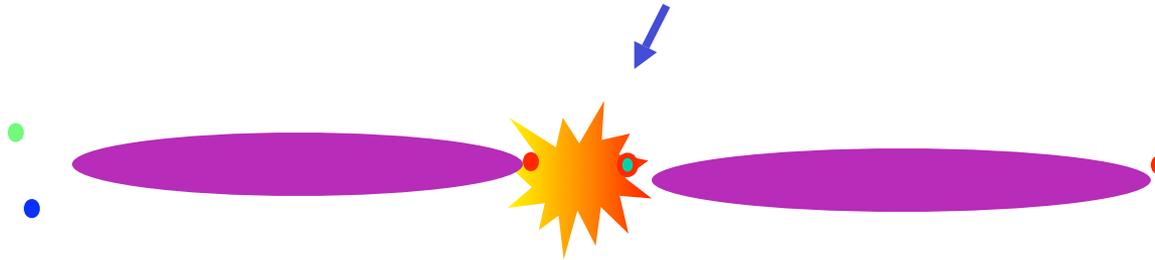


Strong **color** field
Force *grows* with
separation !!!

Confinement - QCD

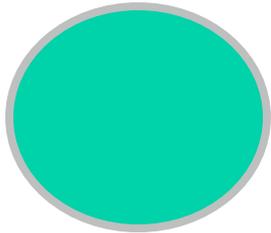
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quark-antiquark pair
created from vacuum

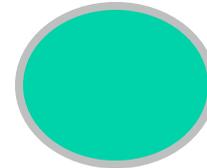


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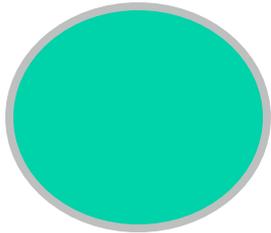
“white” proton
(confined quarks)



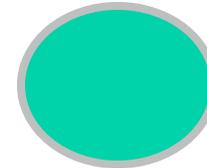
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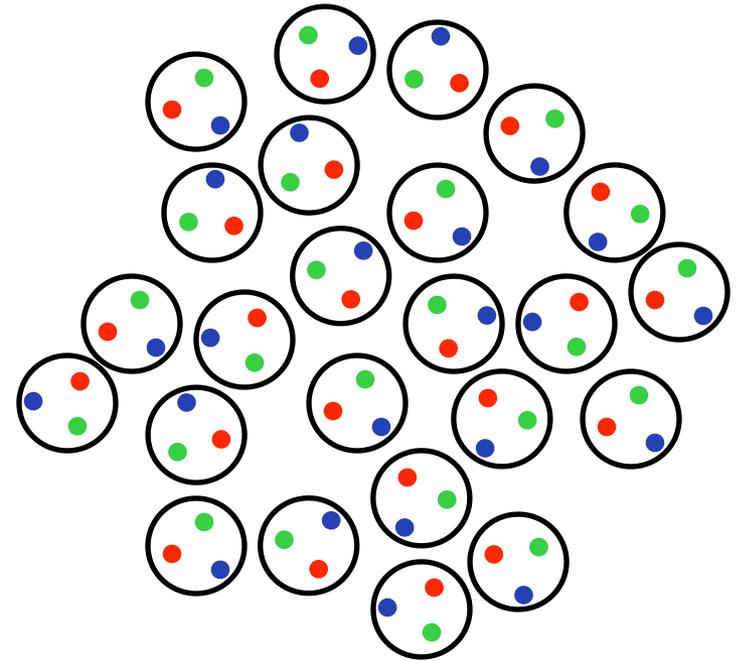
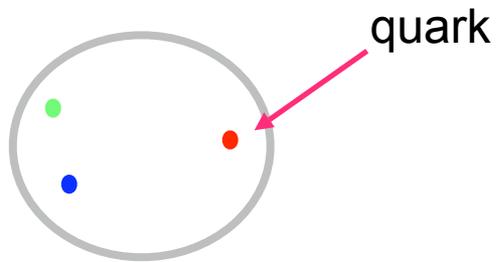
“white” π^0
(confined quarks)

To understand the strong force and confinement: Create and study a system of deconfined colored quarks and gluons

Recreating in the laboratory

We try to make a deconfined state of matter

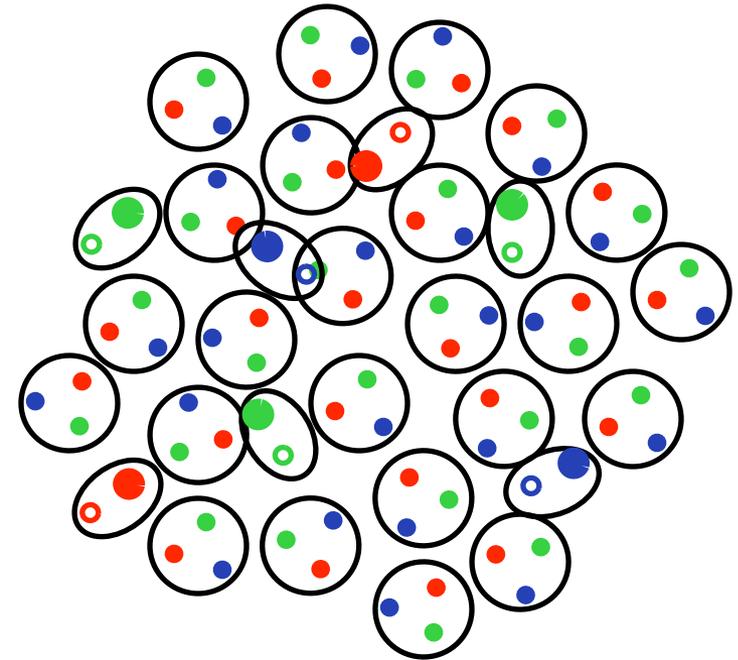
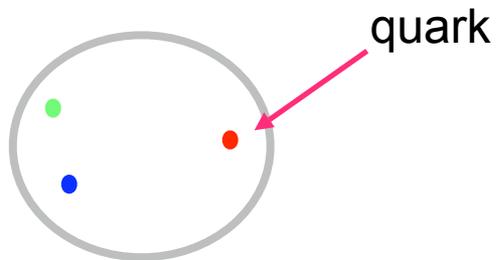
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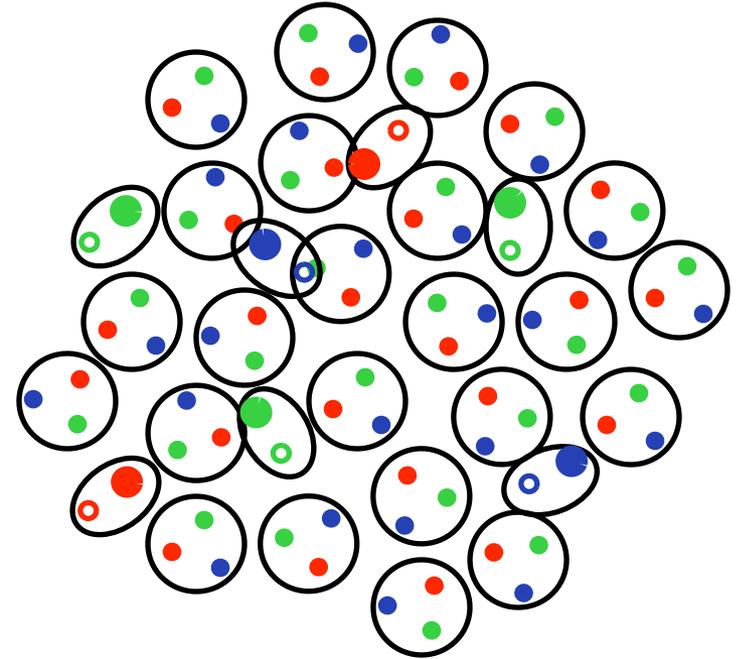
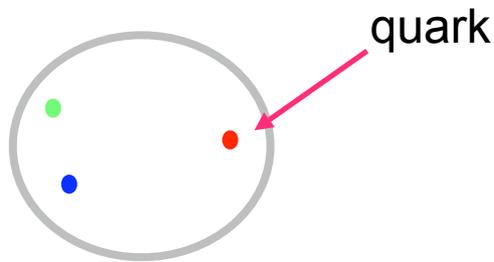


Quarks are also confined within
mesons

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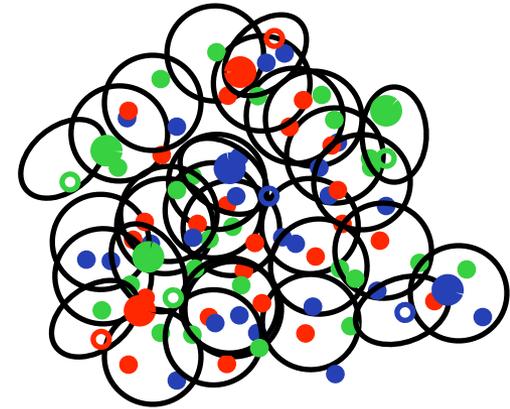
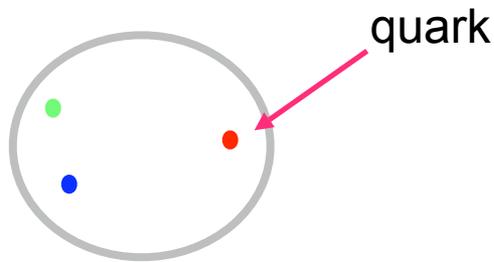
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How can we create a deconfined state of QCD matter?

Recreating in the laboratory

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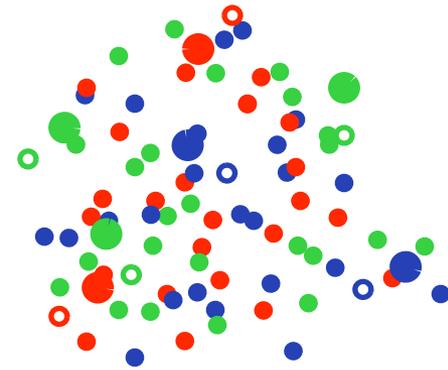
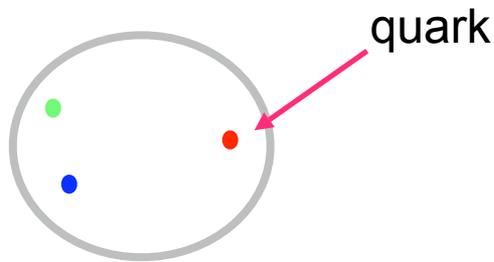
How can we create a deconfined
state of QCD matter?

by heating or compressing

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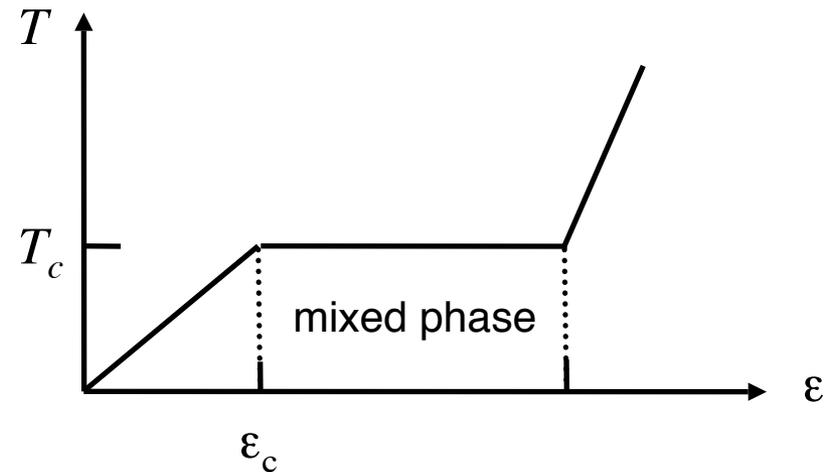
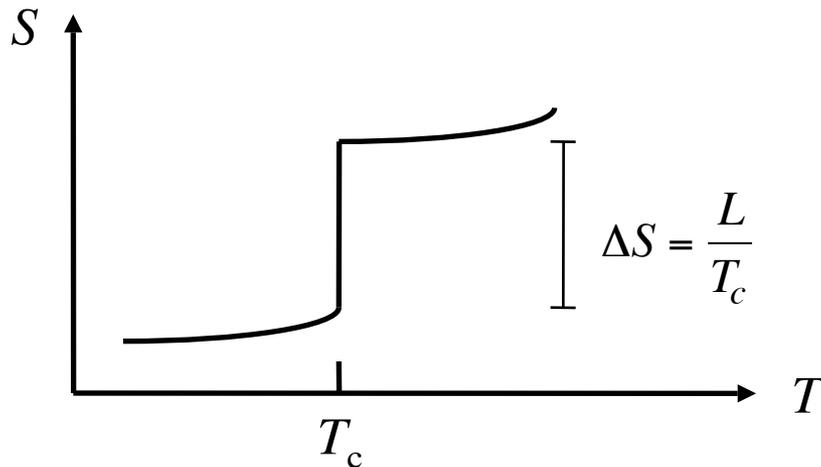
Quark Gluon Plasma
~6x Normal nuclear density

Thermodynamics - phase transitions

Phase transition or a crossover?

Signs of a phase transition:

1st order: **discontinuous in entropy** at T_c → Latent heat, a mixed phase



Higher order: **discontinuous in higher derivatives of $\delta^n S / \delta T^n$** → no mixed phase - system passed smoothly and uniformly into new state (ferromagnet)

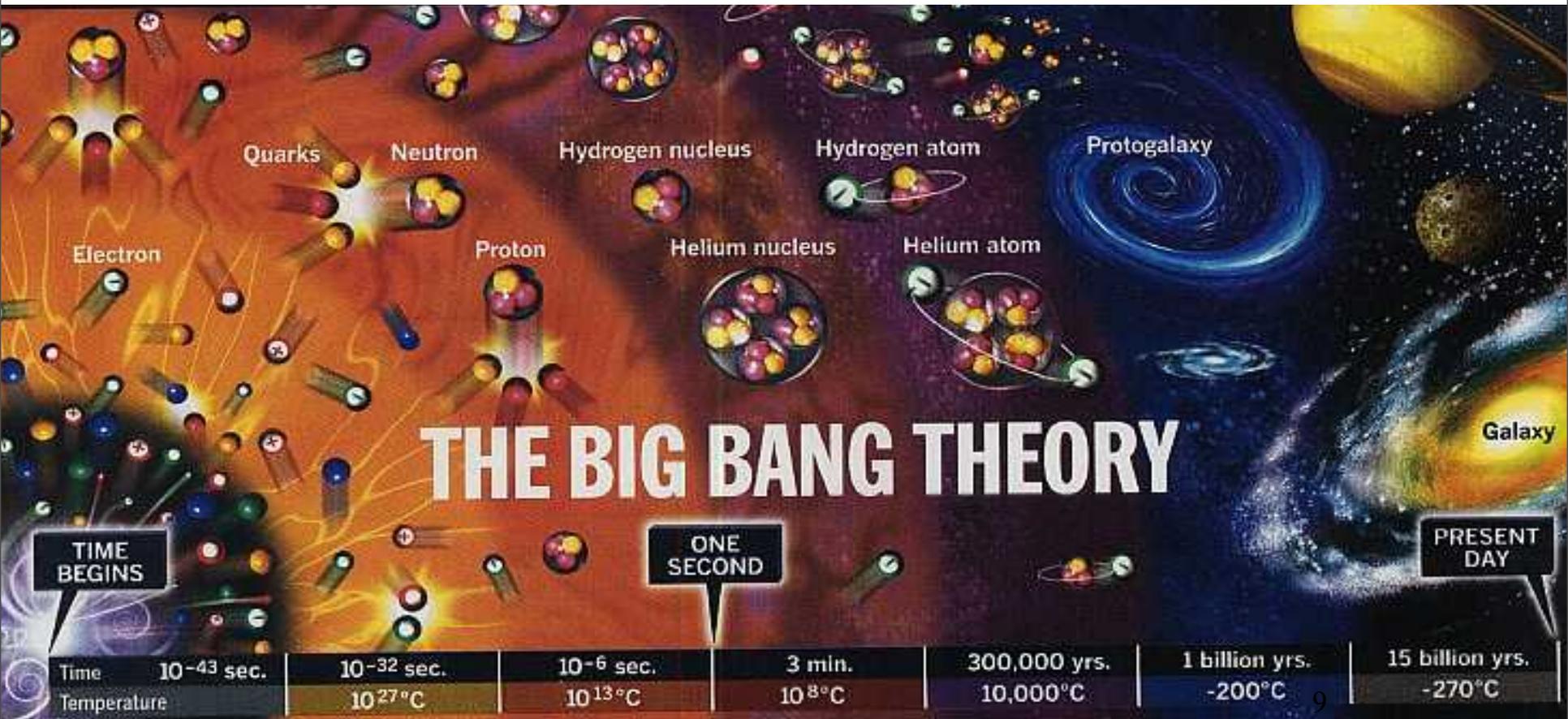
Temperature \Leftrightarrow transverse momentum $T \propto \langle p_T \rangle$

Energy density \Leftrightarrow transverse energy $\epsilon \propto dE_T / dy \cong \langle m_T \rangle dN / dy$

Entropy \Leftrightarrow multiplicity $S \propto dN / dy$

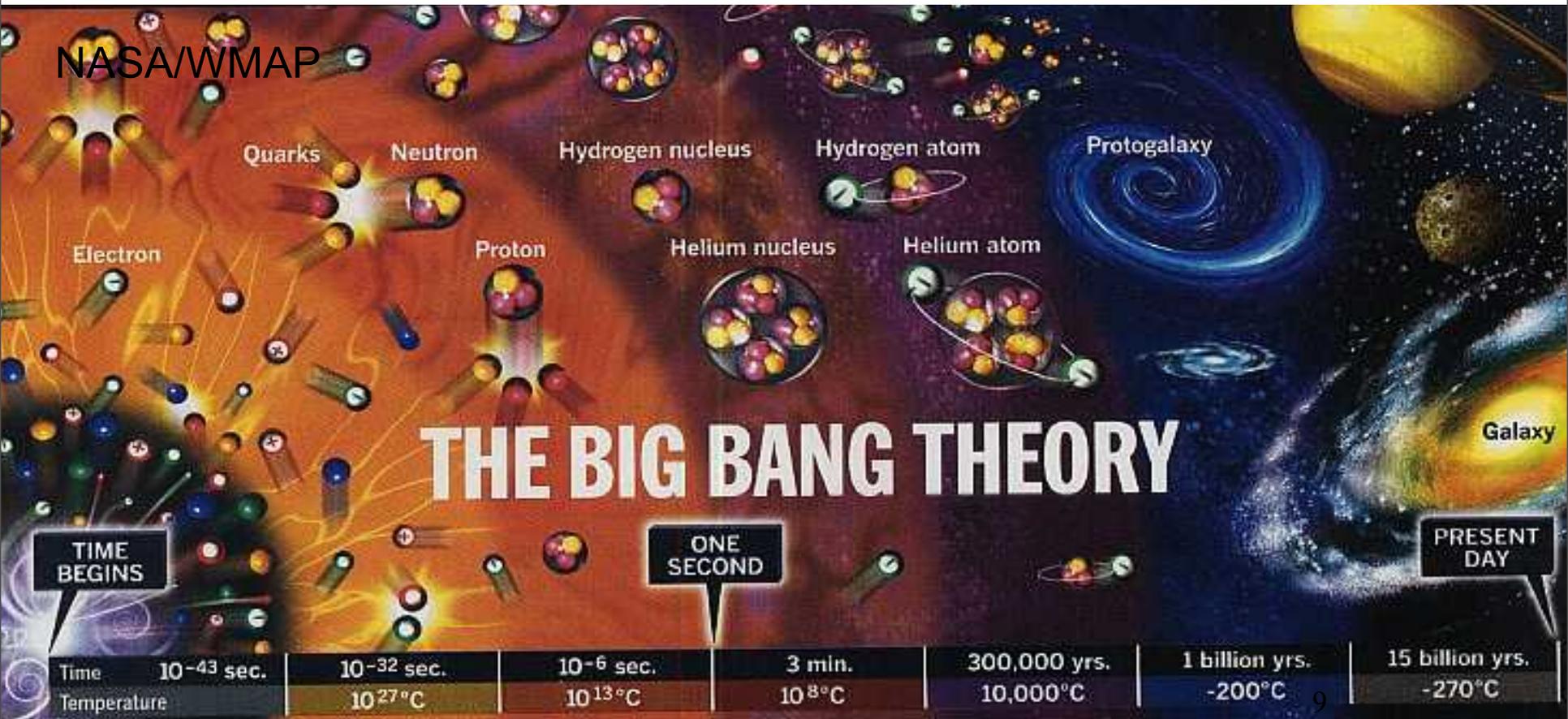
The order of the phase transition

“A **first-order QCD phase transition** that occurred in the early universe would lead to a **surprisingly rich cosmological scenario.**” Ed Witten, Phys. Rev. D (1984)



The order of the phase transition

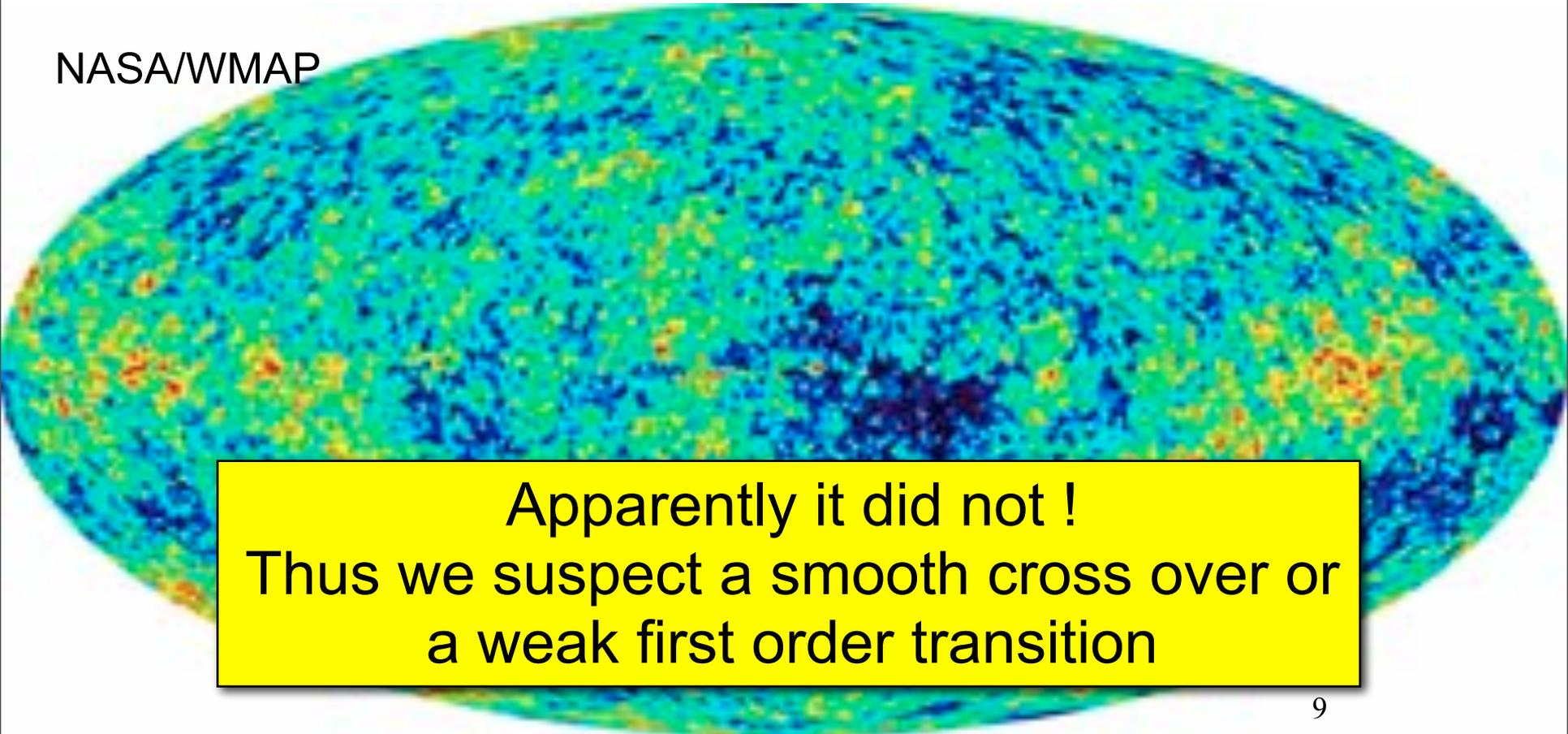
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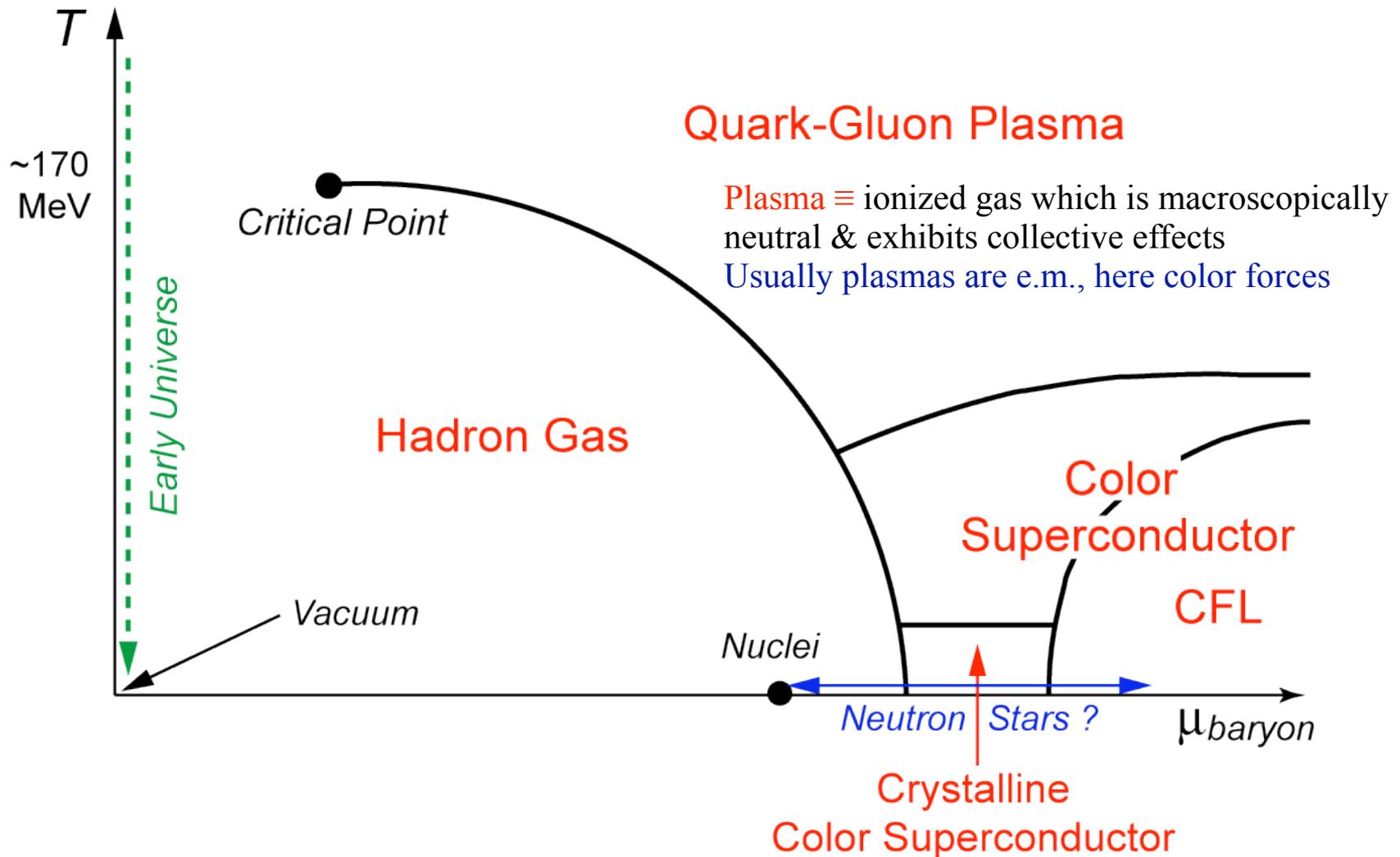
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NASA/WMAP

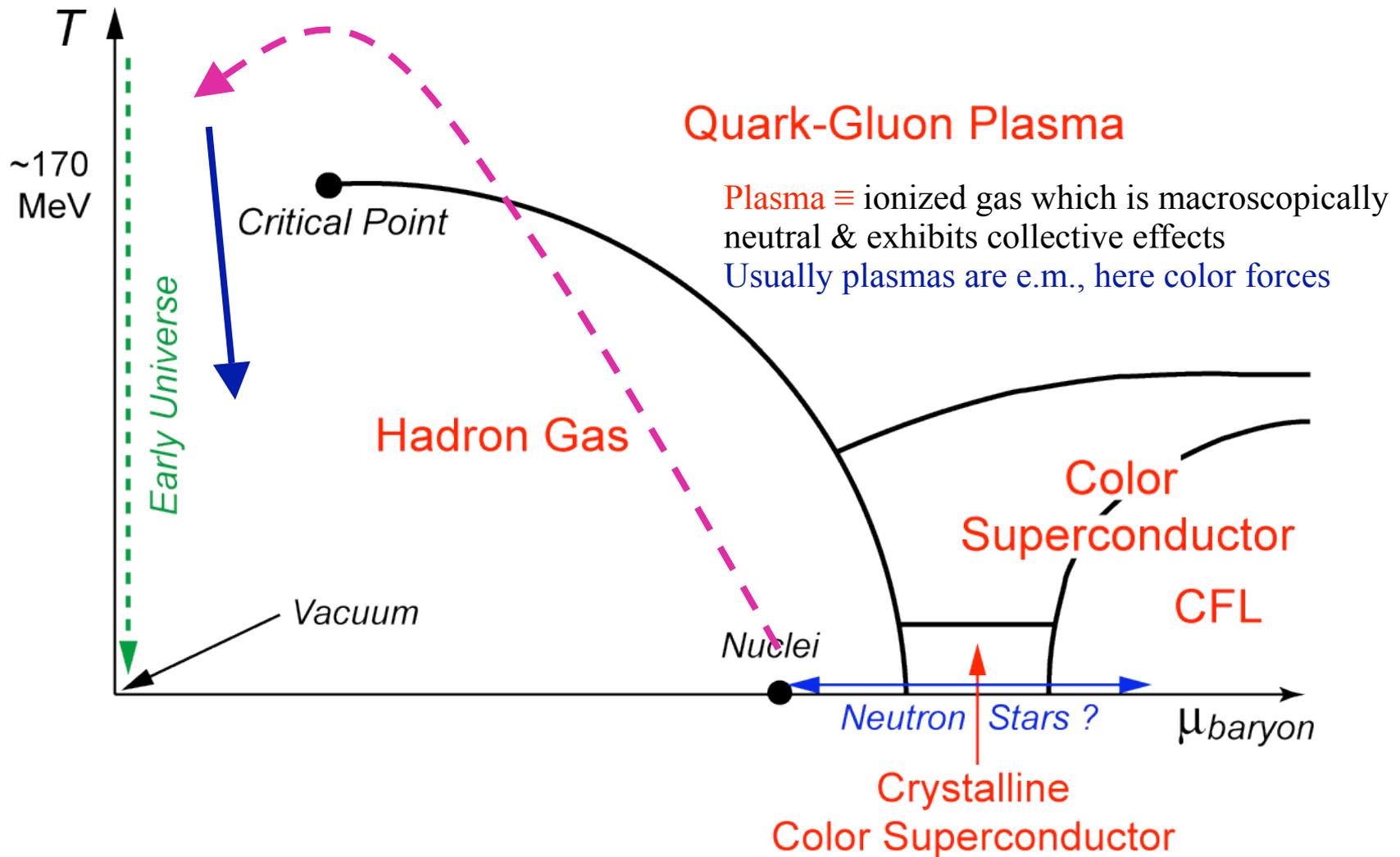


Apparently it did not !
Thus we suspect a smooth cross over or
a weak first order transition

QCD phase diagram of hadronic matter



QCD phase diagram of hadronic matter



RHIC @ Brookhaven National Lab

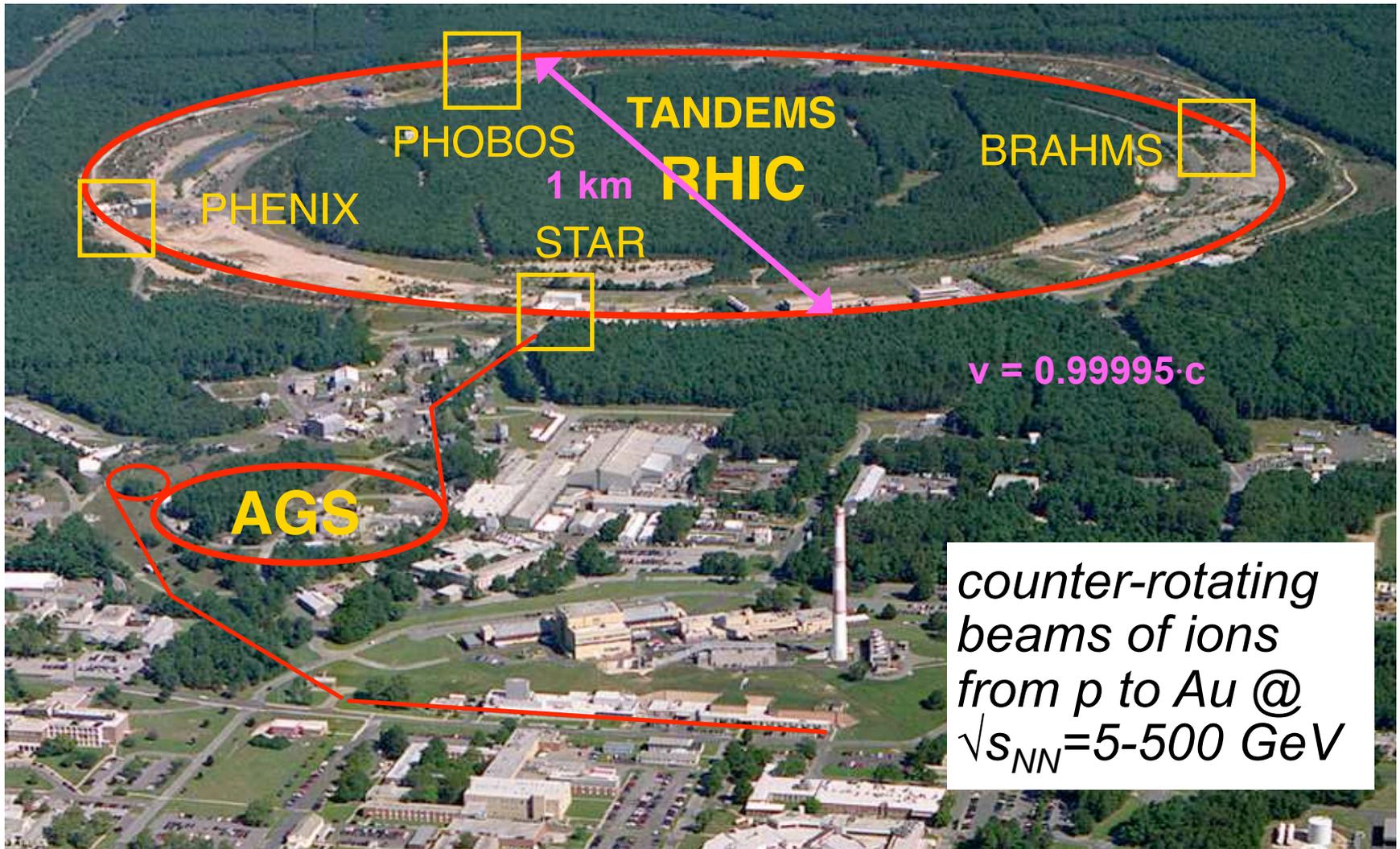


RHIC - **R**elativistic **H**eavy **I**on **C**ollider
3.8 km accelerator that can be seen from space

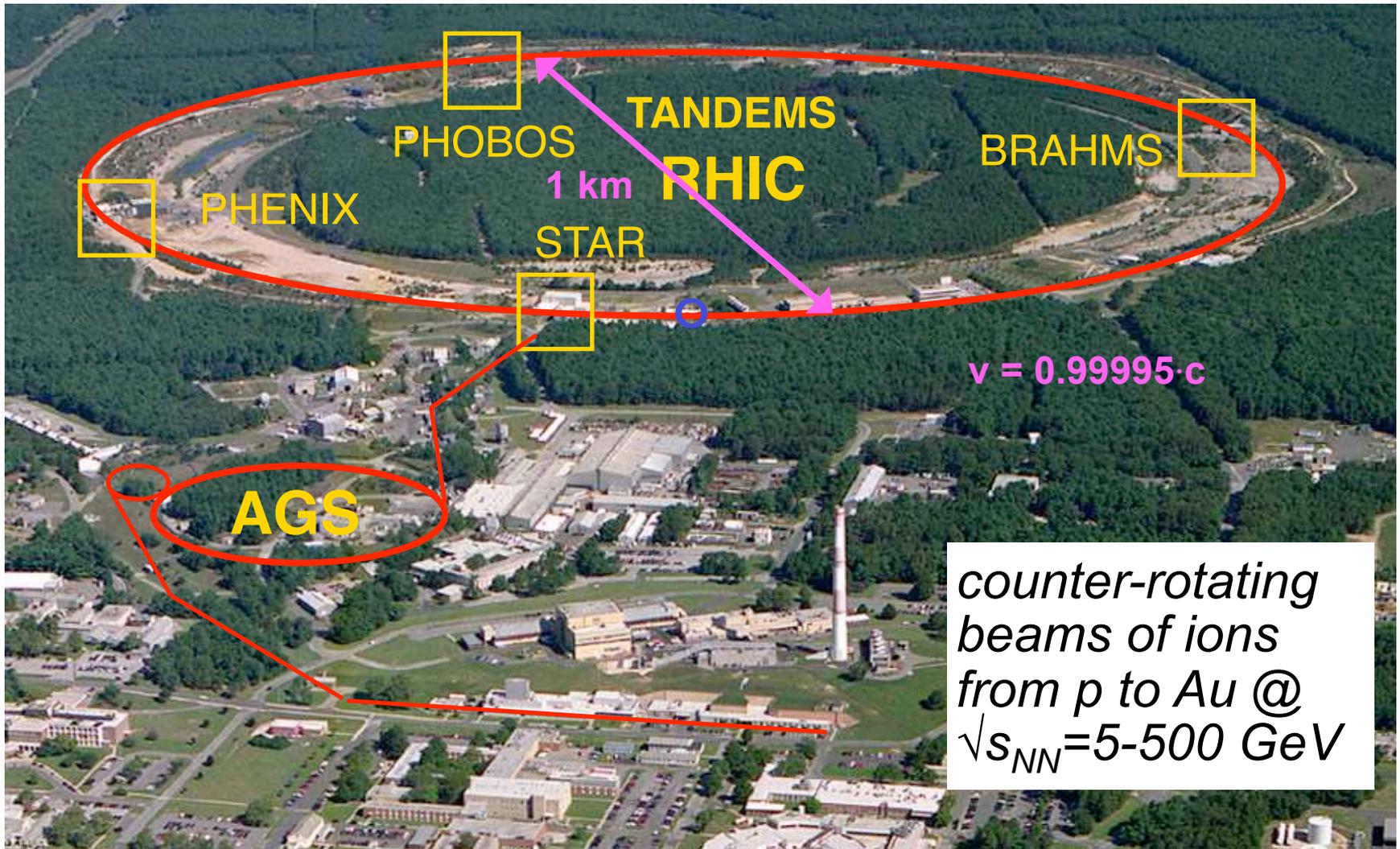
Unusual facts about RHIC

- RHIC's beam travels at **99.995% the speed of light** (186,000 miles per second).
- RHIC's two rings consist of **1740 super-conducting** magnets each cooled by liquid helium to **-269°C**
- RHIC contains **seven tons of helium**
 - enough to fill all the balloons in Macy's Thanksgiving Day Parades for the next 100 years
- The refrigerator to cool the helium needs a **power of 15 MW** (as much as 15000 homes! we shut down over the summer)
- Over 20 years **less than one gram of gold** is used in the beam.
- At top energy: stored beam energy is **200kJ** per ring
 - energy 2000 people get drinking a single drop of beer each

RHIC - the experiments

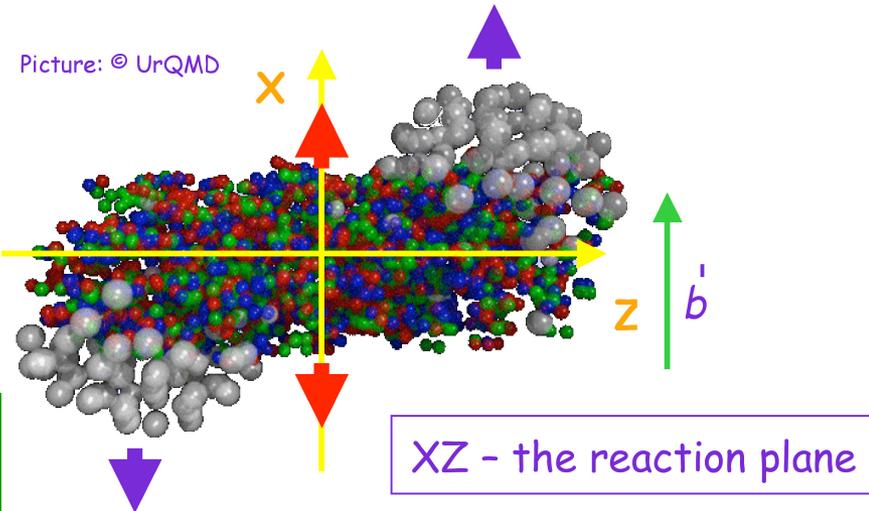
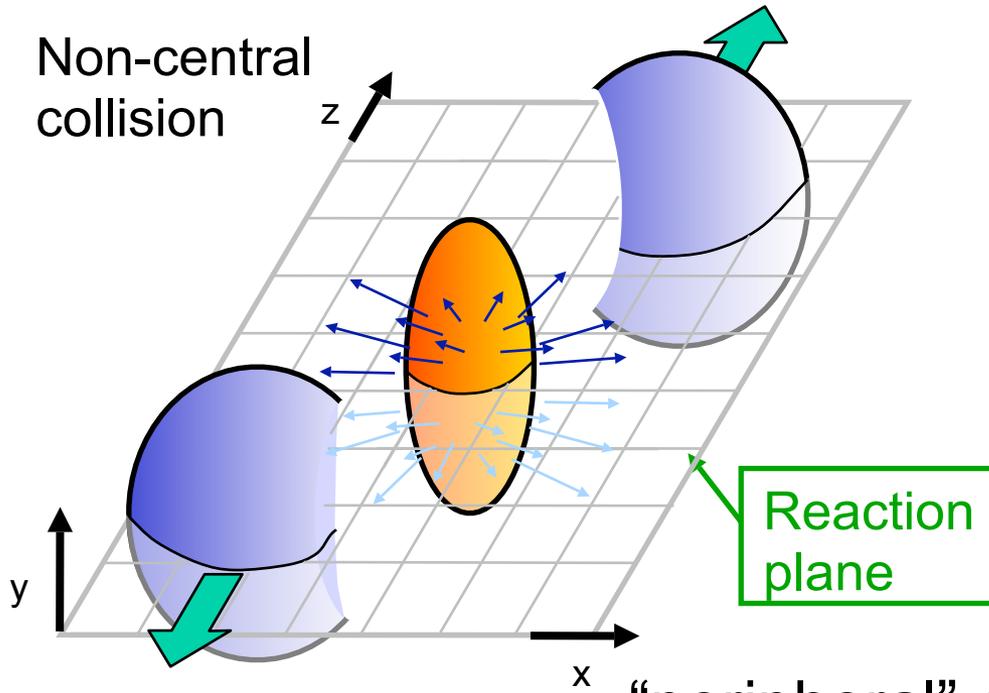


RHIC - the experiments



Geometry of a heavy-ion collision

Non-central collision



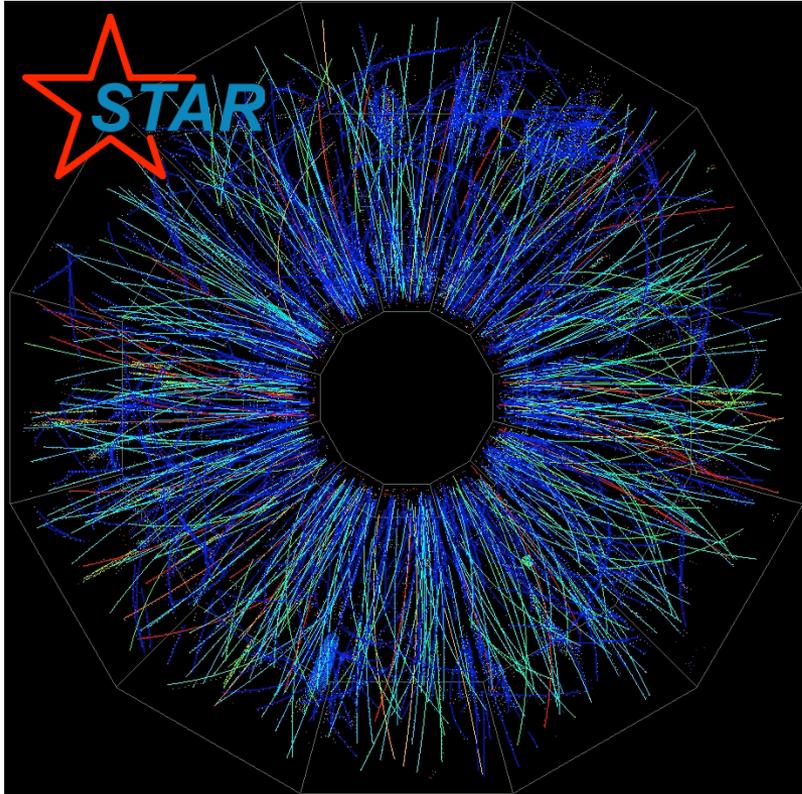
“peripheral” collision ($b \sim b_{\max}$)
 “central” collision ($b \sim 0$)

Number of participants (N_{part}): number of incoming nucleons (participants) in the overlap region

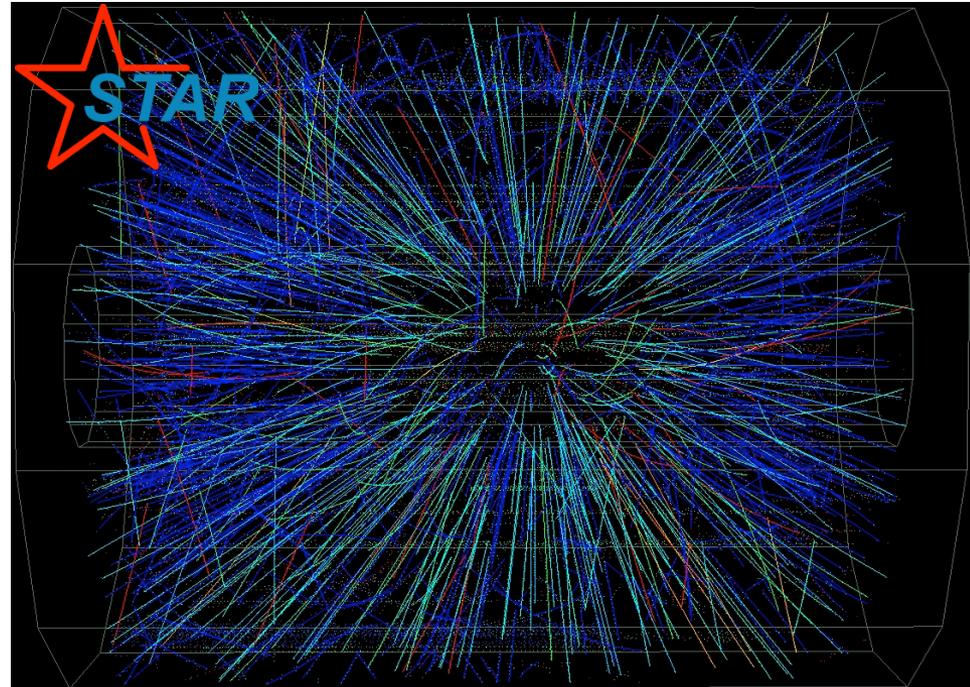
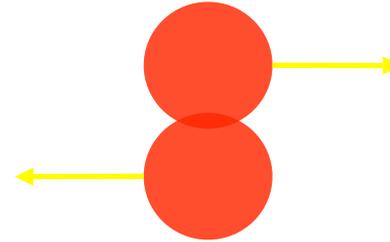
Number of binary collisions (N_{bin}): number of equivalent inelastic nucleon-nucleon collisions

$$N_{\text{bin}} \geq N_{\text{part}}$$

A peripheral Au-Au collision

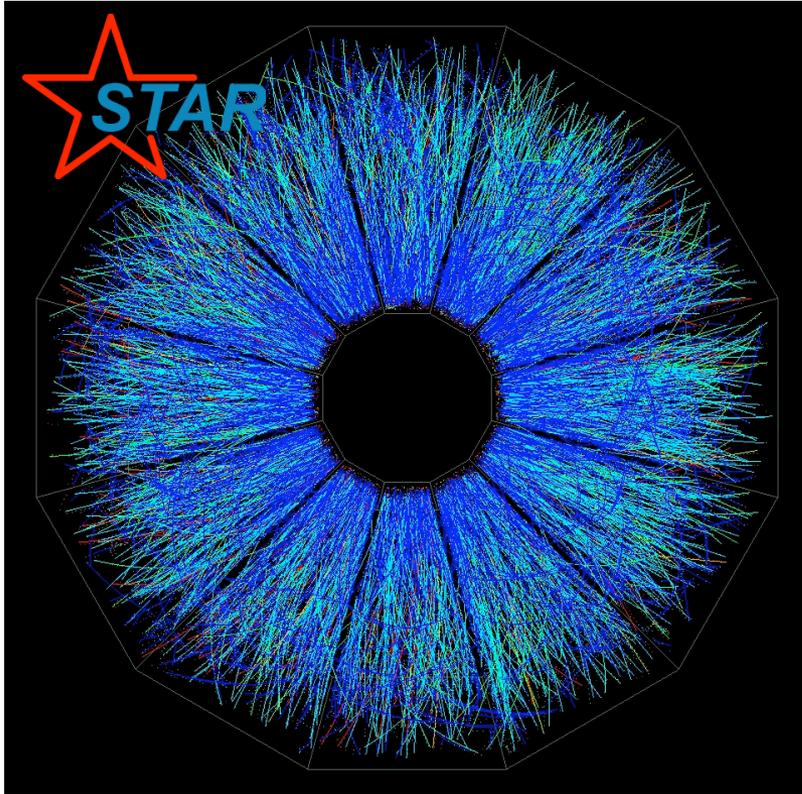


Peripheral Collision



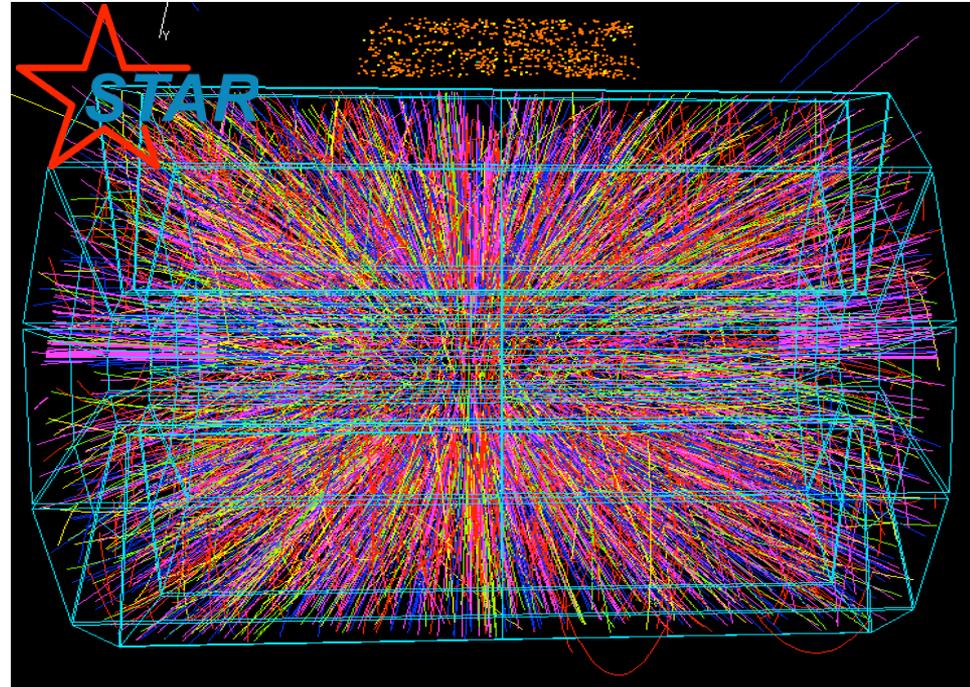
Color \Rightarrow Energy loss in TPC gas

39.4 TeV in central Au-Au collision

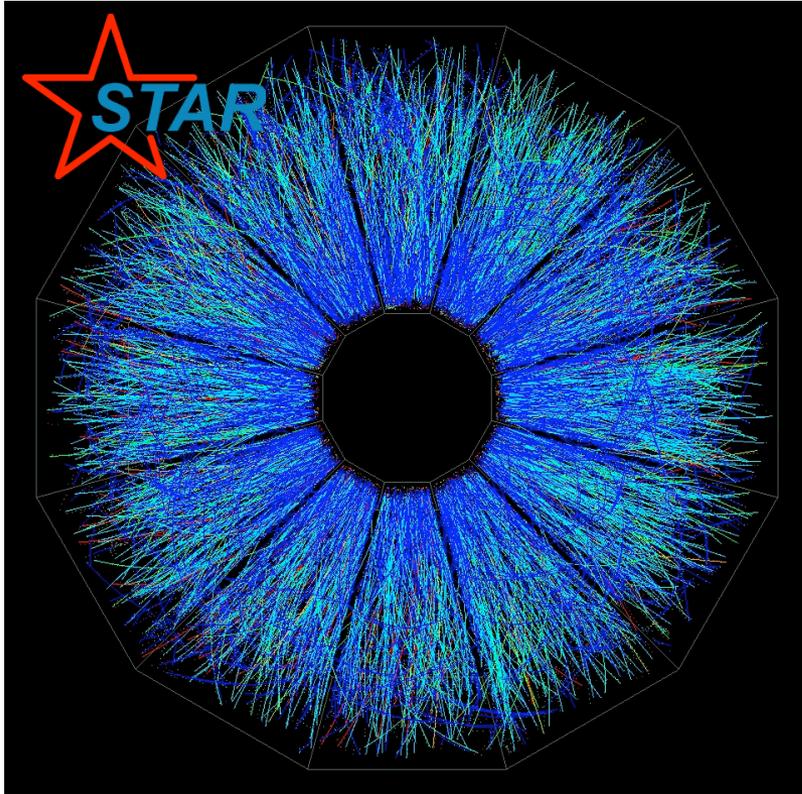


>5000 hadrons and leptons

- Only **charged** particles shown
- Neutrals don't ionise the TPC's gas so are not "seen" by this detector.



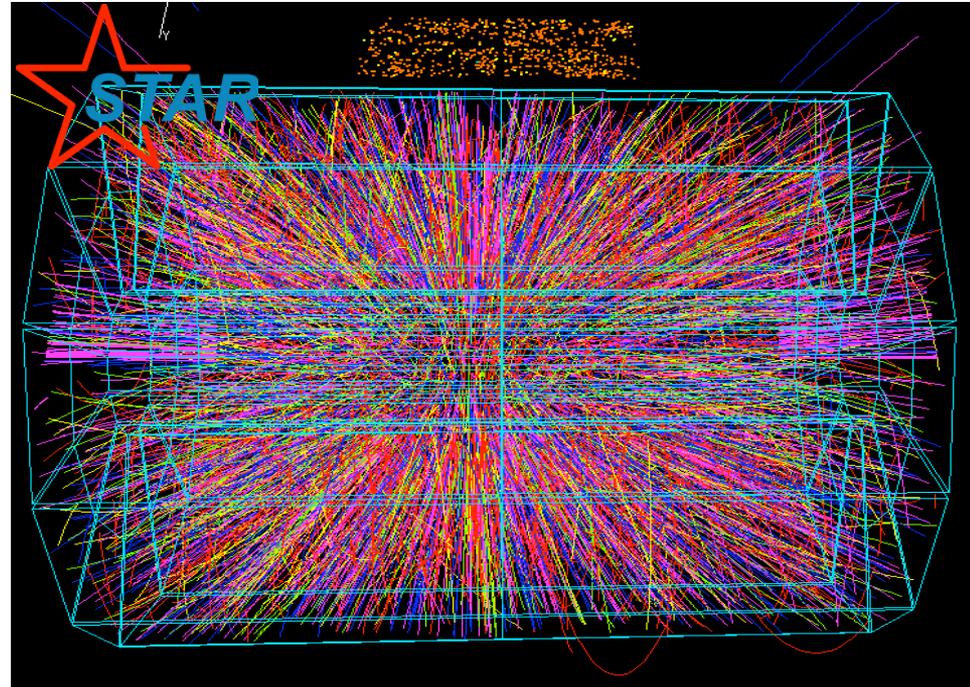
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>5000 hadrons and leptons

26 TeV is removed
from colliding beams.

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The energy is contained in one collision

Central Au+Au Collision:
26 TeV \sim 6 μ Joule



The energy is contained in one collision

Central Au+Au Collision:
26 TeV \sim 6 μ Joule

Sensitivity of human ear:
 10^{-11} erg = 10^{-18} Joule = 10^{-12} μ Joule
A Loud “Bang” if $E \Rightarrow$ Sound



The energy is contained in one collision

Central Au+Au Collision:
26 TeV \sim 6 μ Joule

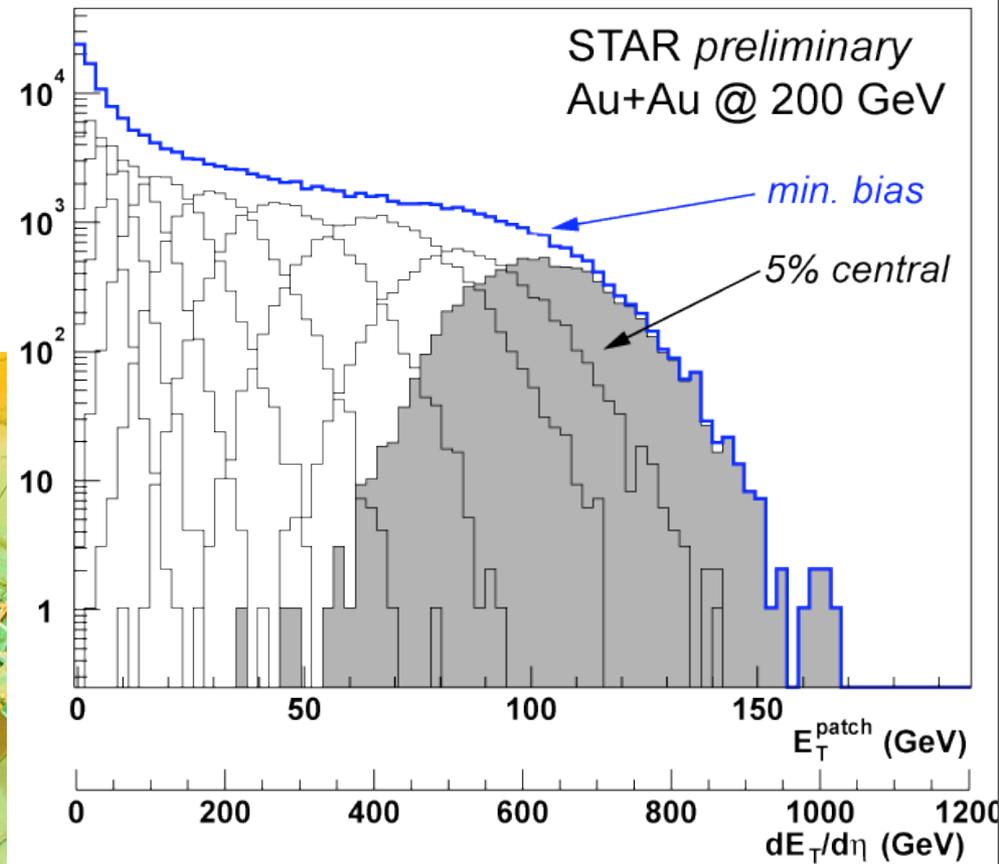
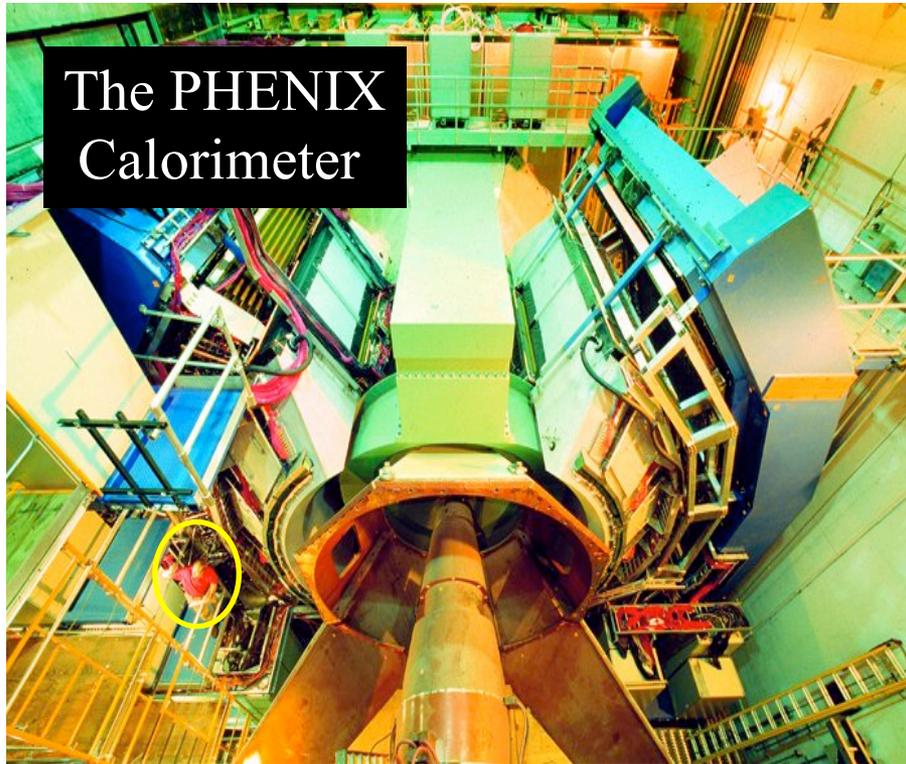
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Most goes into particle creation

Energy density in central Au-Au collisions

- use calorimeters to measure total energy



Energy density in central Au-Au collisions

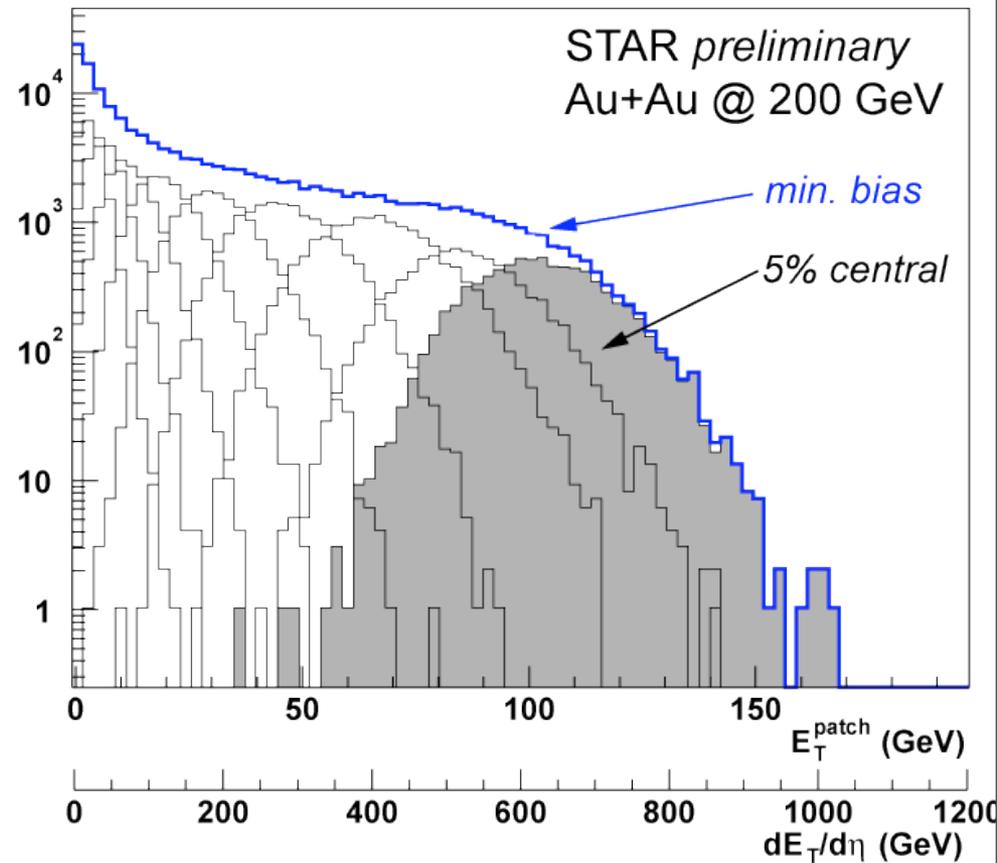
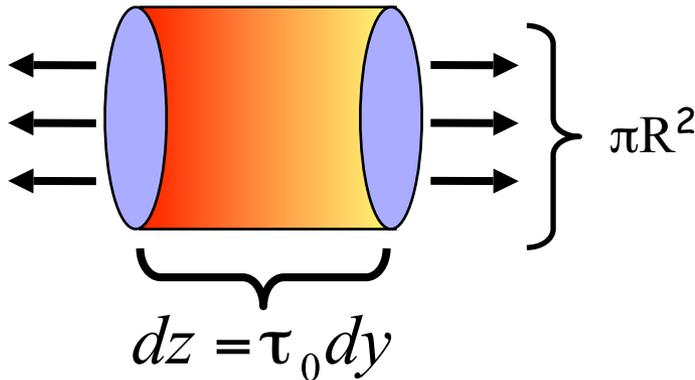
- use calorimeters to measure total energy
- estimate volume of collision

Bjorken-Formula for Energy Density:

$$\epsilon_{Bj} = \frac{\Delta E_T}{\Delta V} = \frac{1}{\pi R^2} \frac{1}{\tau_0} \frac{dE_T}{dy}$$

$R \sim 6.5$ fm

Time it takes to thermalize system ($t_0 \sim 1$ fm/c)



Energy density in central Au-Au collisions

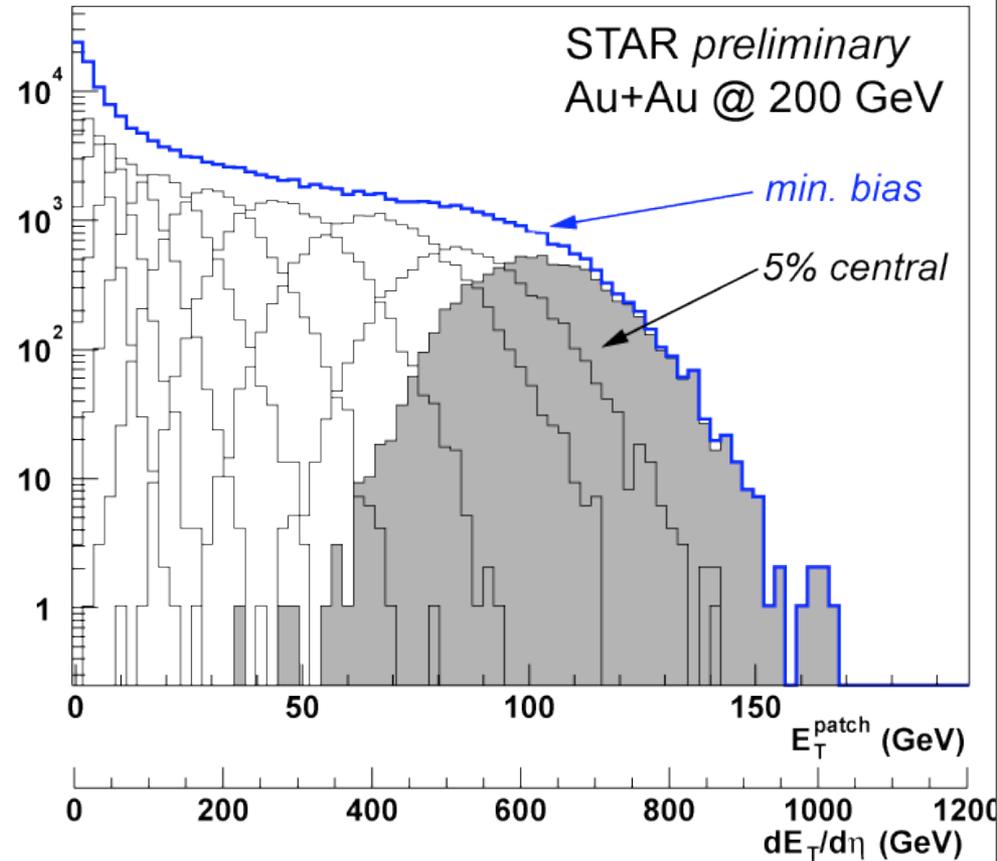
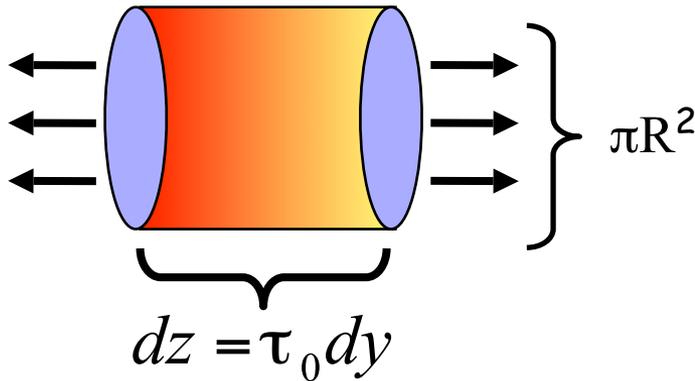
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Time it takes to thermalize system ($t_0 \sim 1$ fm/c)



$$\epsilon_{BJ} \approx 5.0 \text{ GeV/fm}^3$$

~ 30 times normal nuclear density

~ 5 times $> \epsilon_{\text{critical}}$ (lattice QCD)

5 GeV/fm³. Is that a lot?

In a year, the U.S. uses ~100 quadrillion BTUs of energy
(1 BTU = 1 burnt match):

$$100 \times 10^{15} \text{ BTU} \times \frac{1060 \text{ J}}{\text{BTU}} \times \frac{1 \text{ eV}}{1.6 \times 10^{-19} \text{ J}} = 6.6 \times 10^{38} \text{ eV}$$

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At 5 GeV/fm³, this would fit in a volume of:

$$6.6 \times 10^{38} \text{ eV} \div \frac{5 \times 10^9 \text{ eV}}{\text{fm}^3} = 1.3 \times 10^{29} \text{ fm}^3$$

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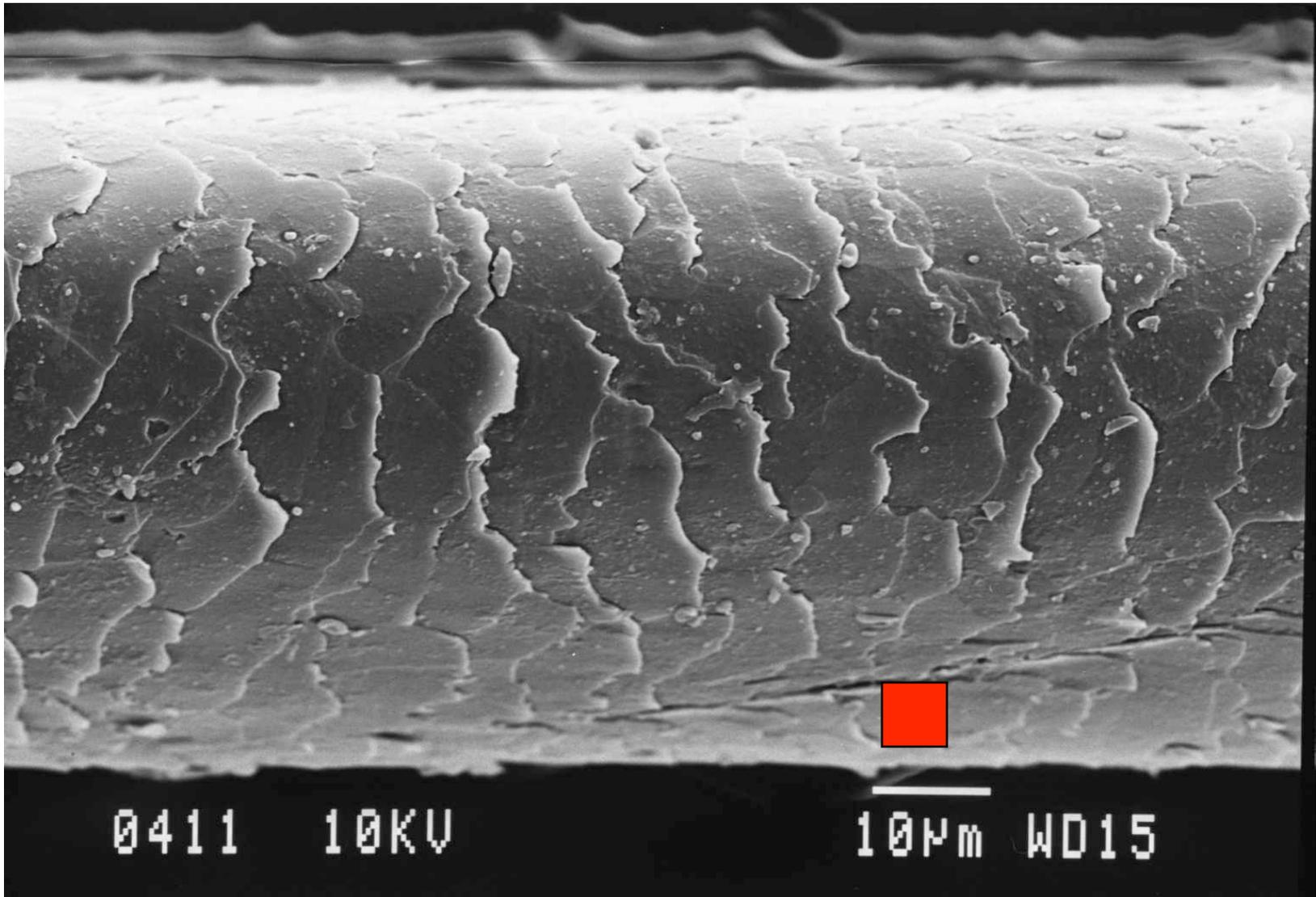
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$$6.6 \times 10^{38} \text{ eV} \div \frac{5 \times 10^9 \text{ eV}}{\text{fm}^3} = 1.3 \times 10^{29} \text{ fm}^3$$

Or, in other words, in a box of the following dimensions:

$$\sqrt[3]{1.3 \times 10^{29} \text{ fm}^3} = 5 \times 10^9 \text{ fm} = 5 \mu\text{m}$$

A human hair



What is the temperature of the medium?

- **Statistical Thermal Models:**
 - Assume a system that is **thermally** (constant T_{ch}) and **chemically** (constant n_i) **equilibrated**
 - System composed of non-interacting hadrons and resonances
 - Obey conservation laws: Baryon Number, Strangeness, Isospin
- Given T_{ch} and μ 's (+ system size), n_i 's can be calculated in a grand canonical ensemble

$$n_i = \frac{g}{2\pi^2} \int_0^\infty \frac{p^2 dp}{e^{(E_i(p) - \mu_i)/T} \pm 1}, \quad E_i = \sqrt{p^2 + m_i^2}$$

Fitting the particle ratios

Number of particles of a given species related to temperature

$$dn_i \sim e^{-(E-\mu_B)/T} d^3p$$

- Assume all particles described by same temperature T and μ_B
- one ratio (e.g., \bar{p} / p) determines μ / T :

$$\frac{\bar{p}}{p} = \frac{e^{-(E-\mu_B)/T}}{e^{-(E-\mu_B)/T}} = e^{-2\mu_B/T}$$

- A second ratio (e.g., K / π) provides $T \rightarrow \mu$

$$\frac{K}{\pi} = \frac{e^{-E_K/T}}{e^{-E_\pi/T}} = e^{-(E_K - E_\pi)/T}$$

- Then all other hadronic ratios (and yields) defined

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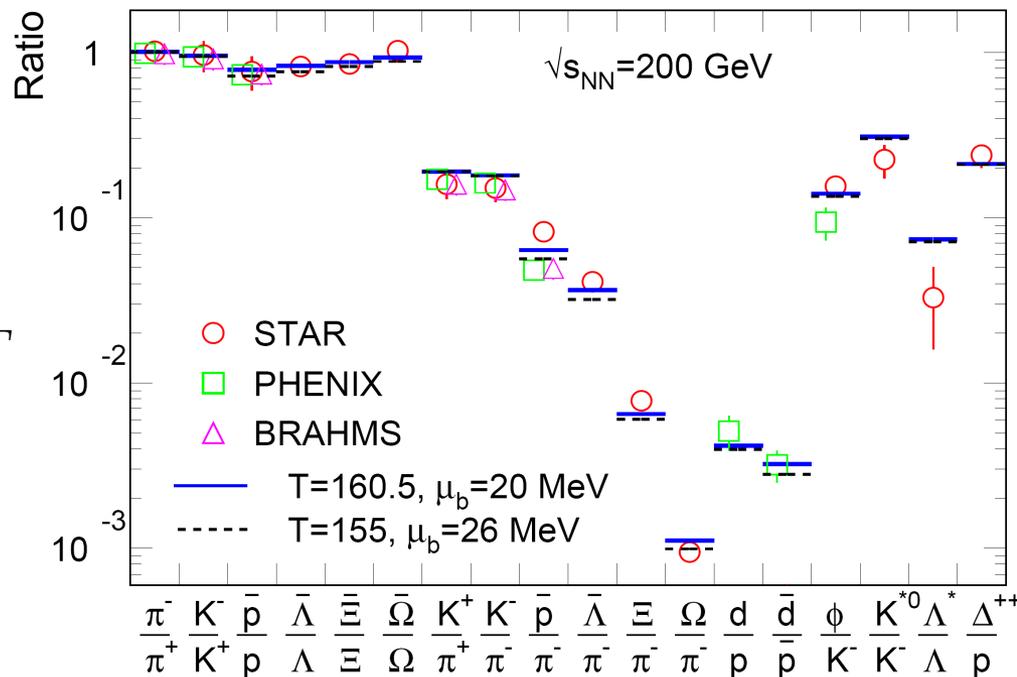
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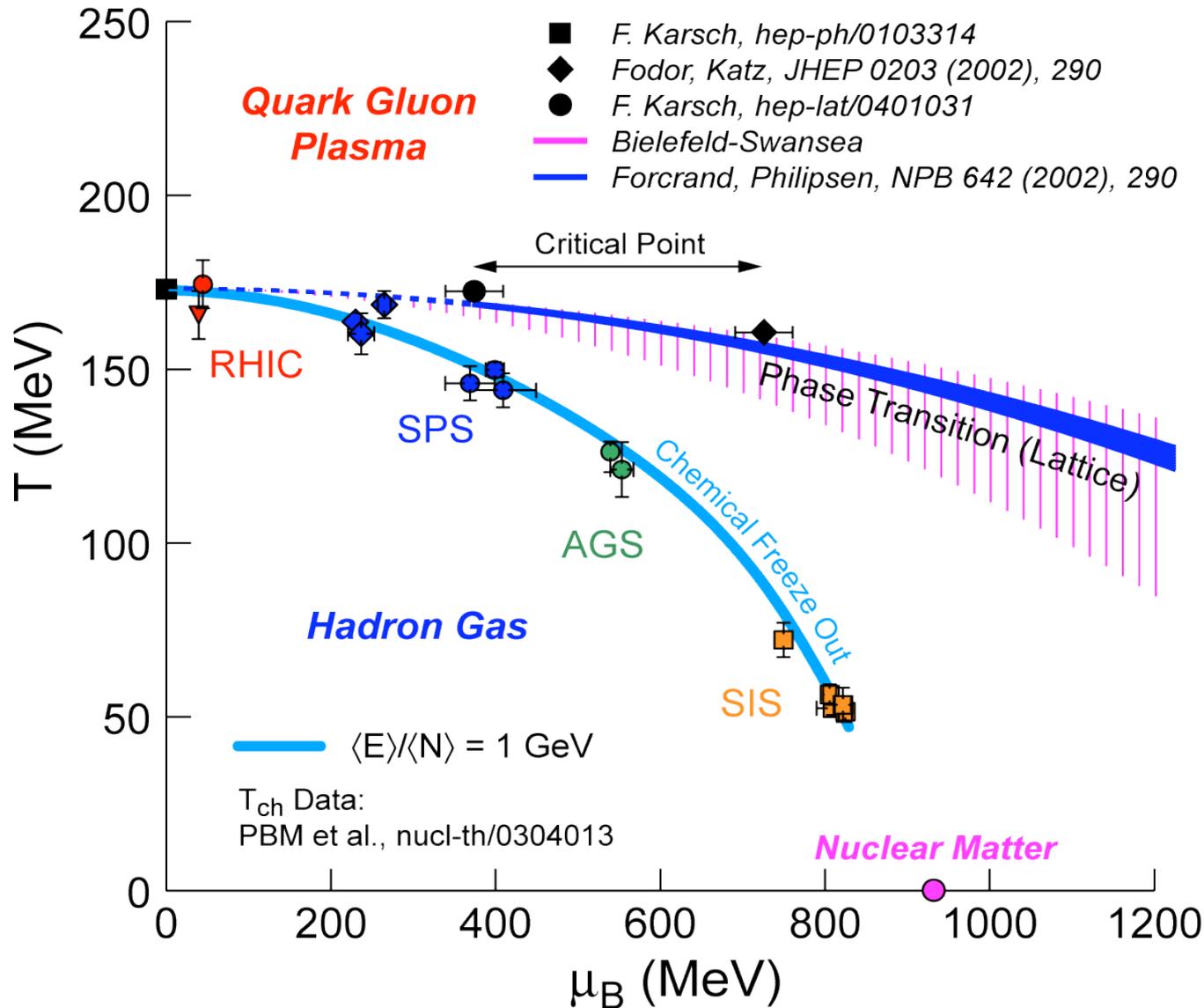
A. Adronic *et al.*, NPA772:167



$T \sim 160 \text{ MeV}, \mu_b \sim 20 \text{ MeV}$

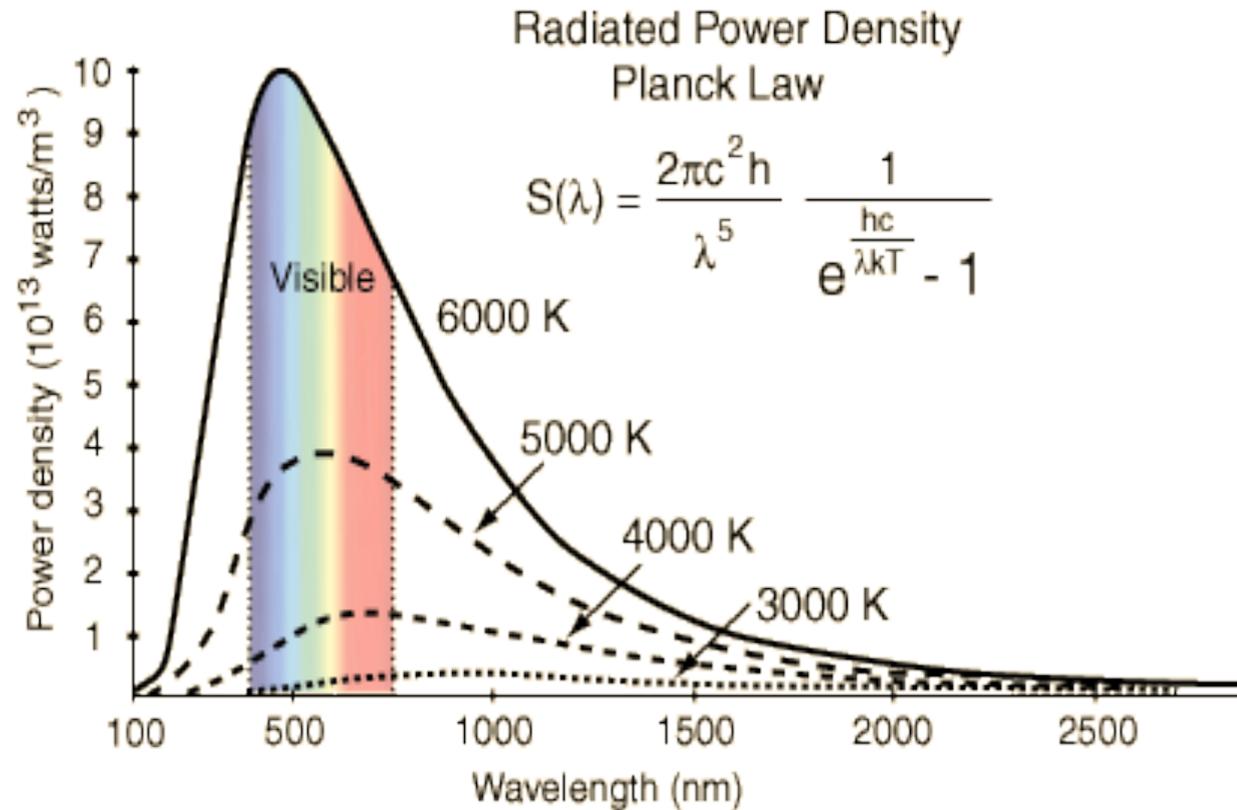
Temperature needed to make QGP

Where RHIC sits on the phase diagram



Blackbody radiation

Planck distribution describes **intensity** as a **function of the wavelength** of the emitted radiation

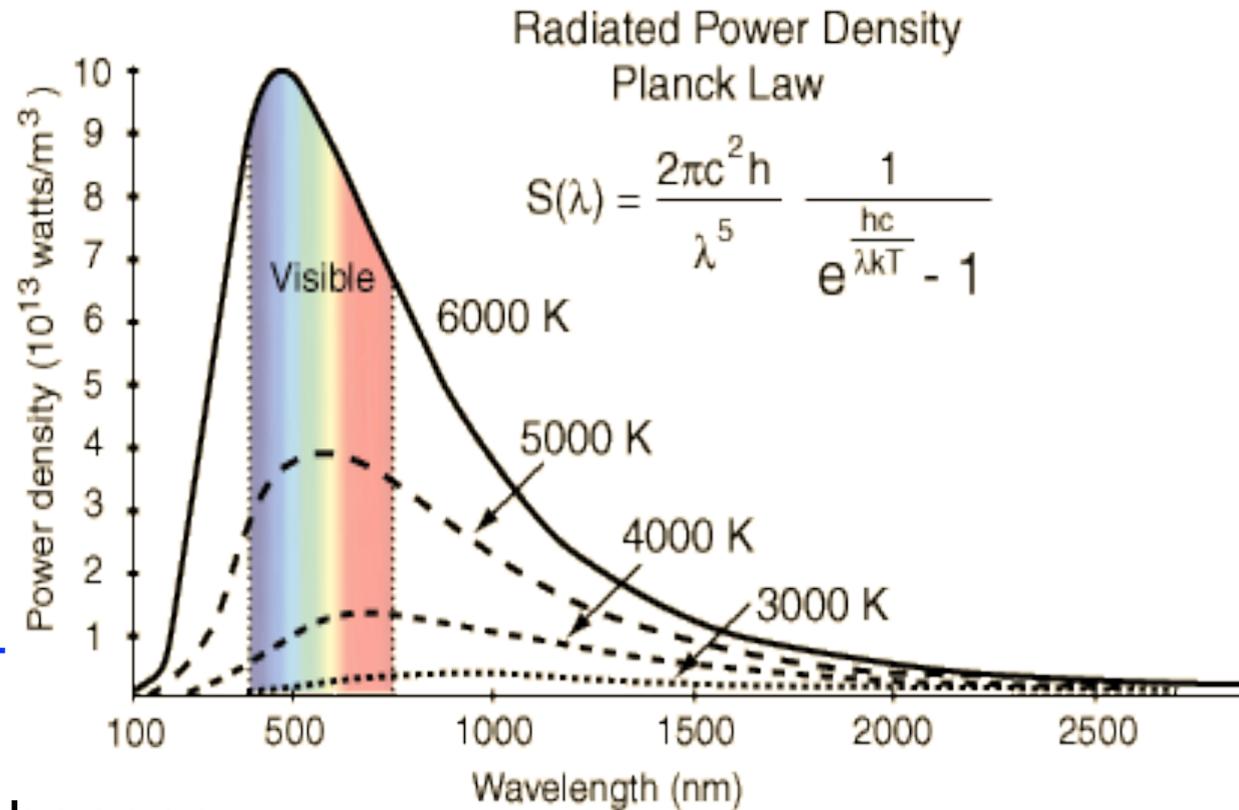


Blackbody radiation

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“Blackbody” radiation is the spectrum of radiation emitted by an object at temperature T

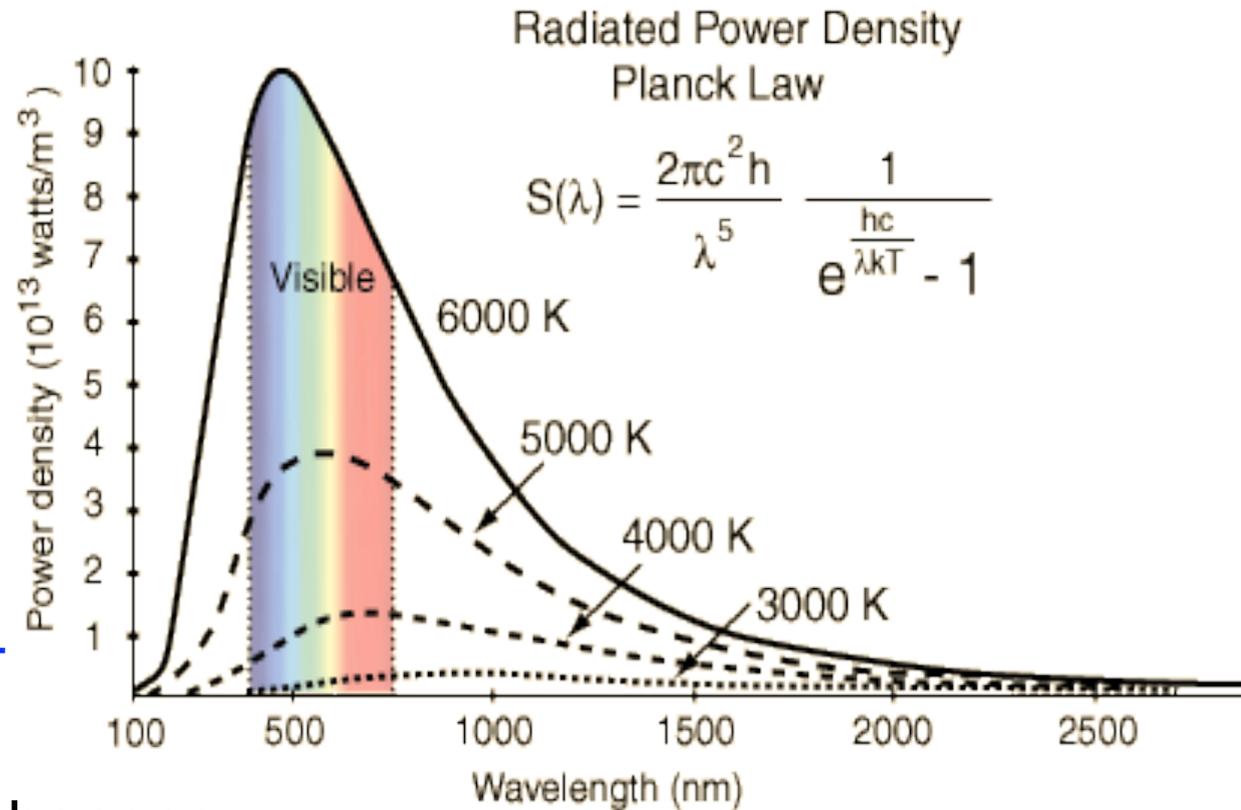
As T increases curve changes



Blackbody radiation

Planck distribution describes intensity as a function of the wavelength of the emitted radiation

“Blackbody” radiation is the spectrum of radiation emitted by an object at temperature T



As T increases curve changes

$$1/\text{Wavelength} \propto \text{Frequency} \propto E \propto p$$

Determining the temperature

From transverse momentum distribution deduce temperature

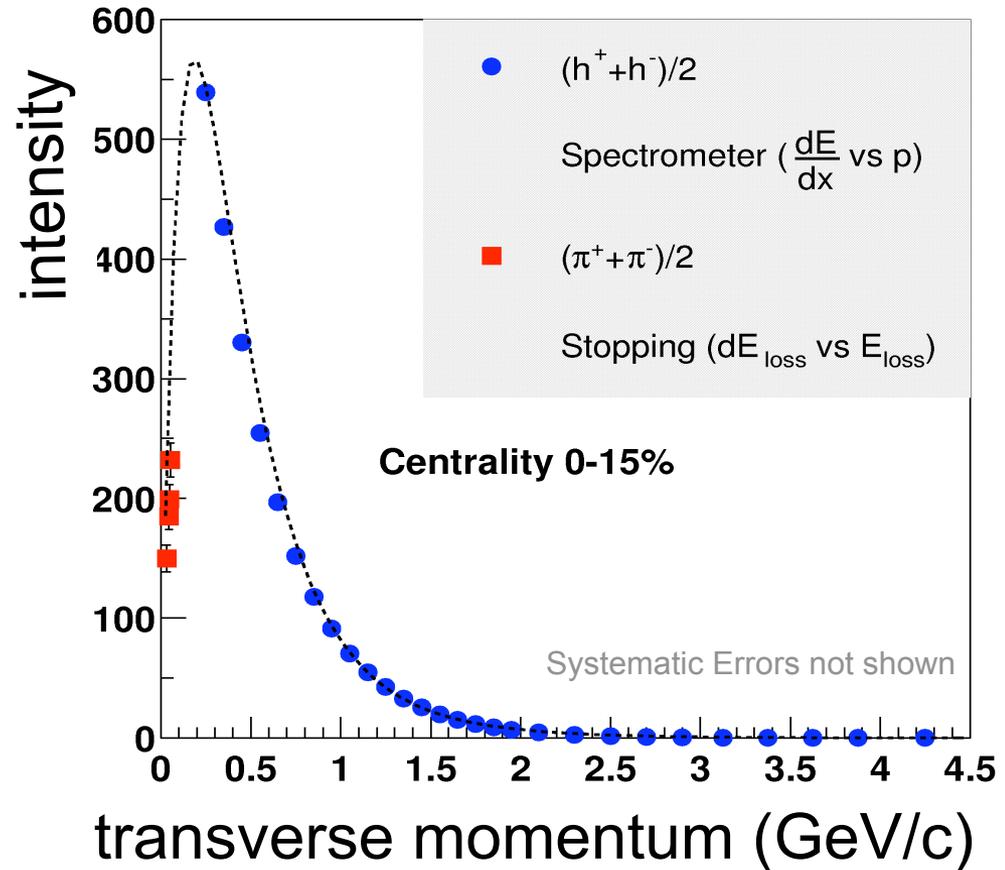
~120 MeV

$$E = \frac{3}{2}kT$$

$$T = \frac{2E}{3k}$$

$$= \frac{2 \times 120 \times 10^6}{3 \times 1.4 \times 10^{-23}} \times 1.6 \times 10^{-19}$$

$$\sim 9 \times 10^{11} K$$



Determining the temperature

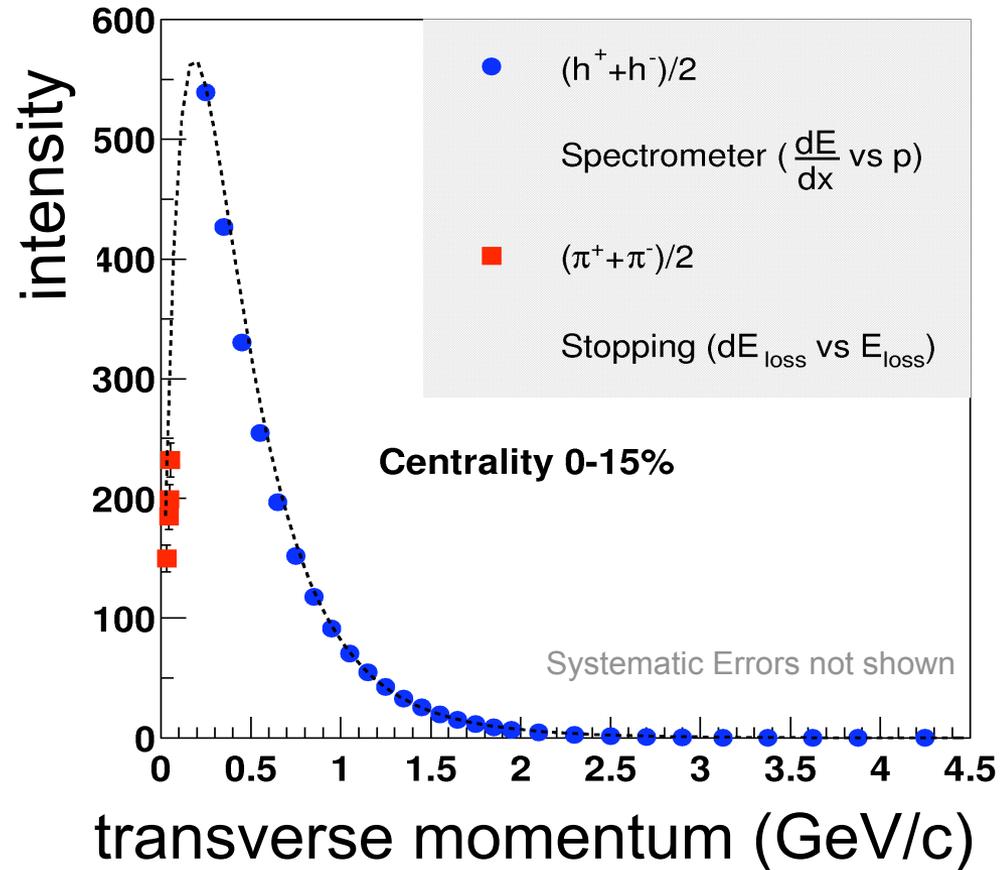
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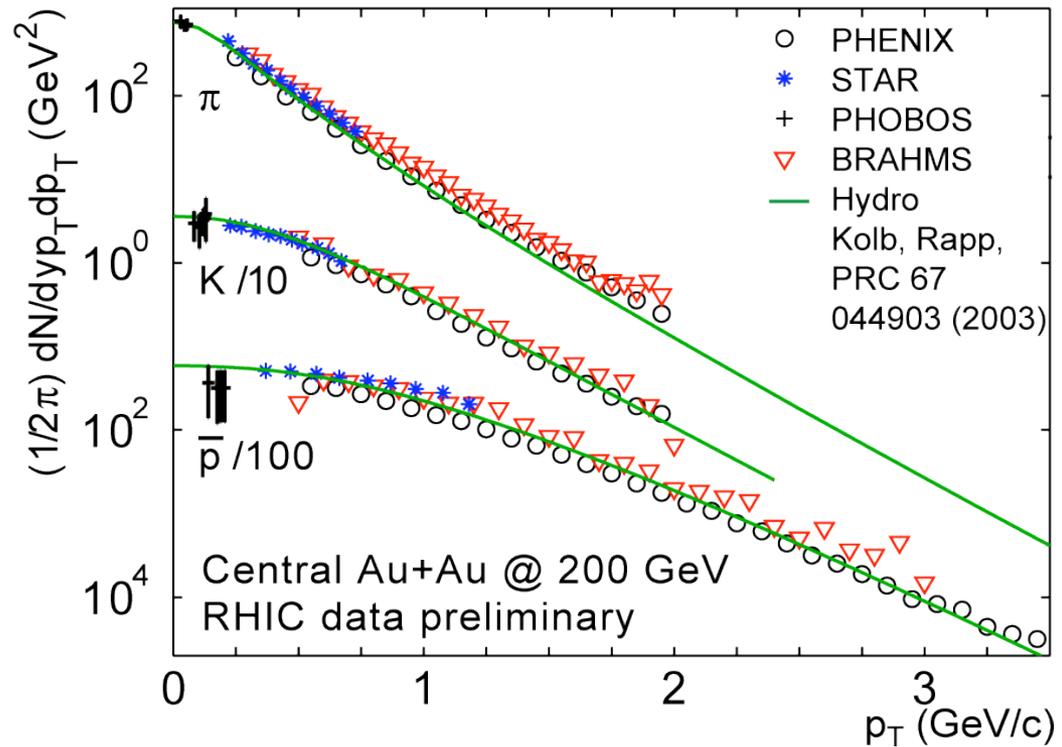
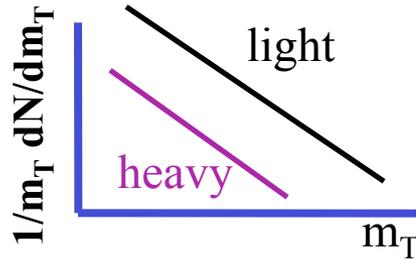
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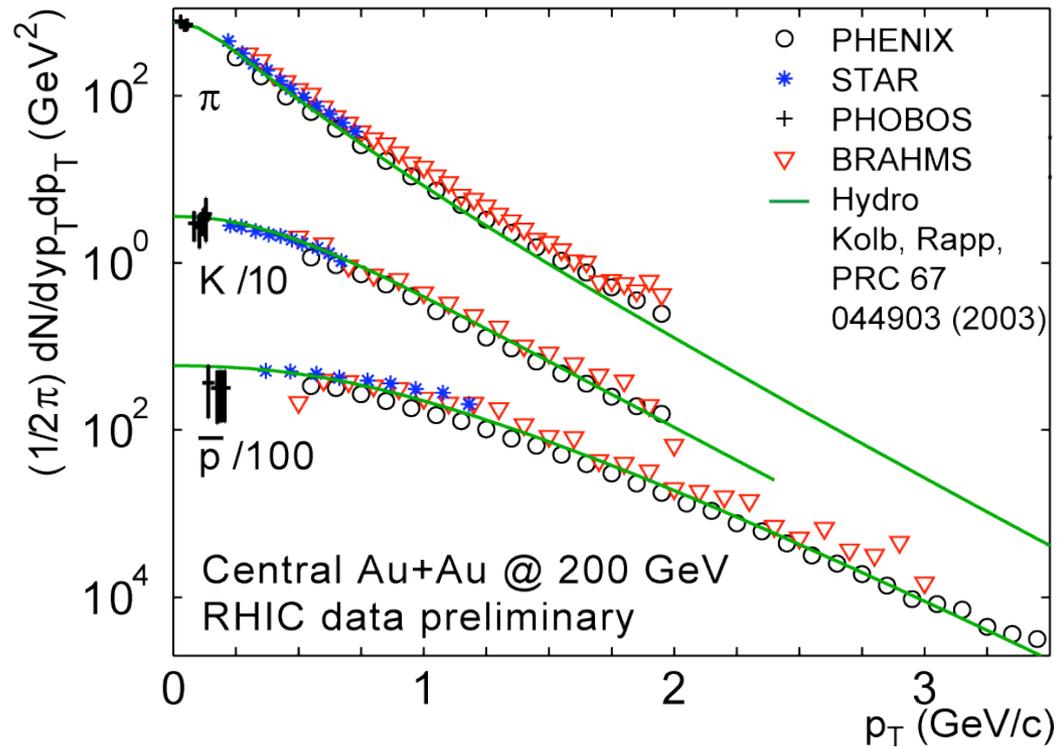
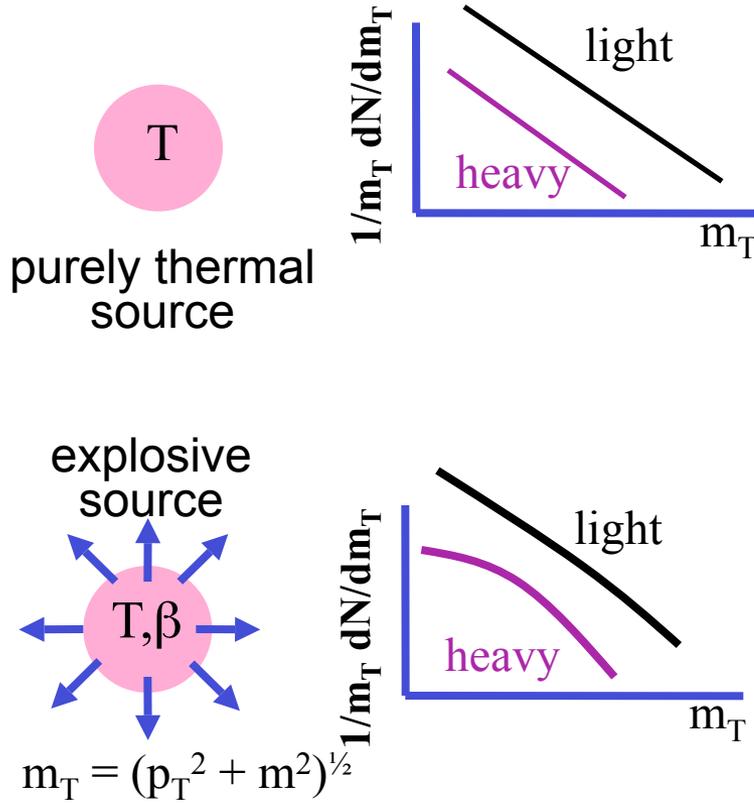
Again temperature needed to create QGP

Strong collective radial expansion

T
purely thermal
source

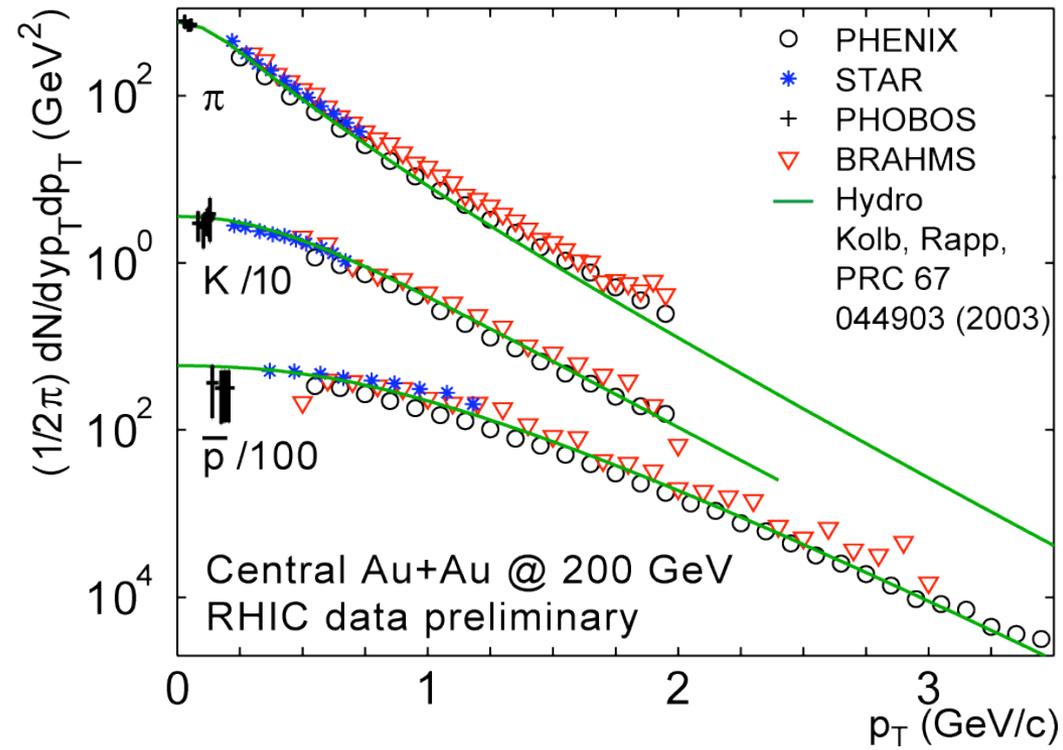
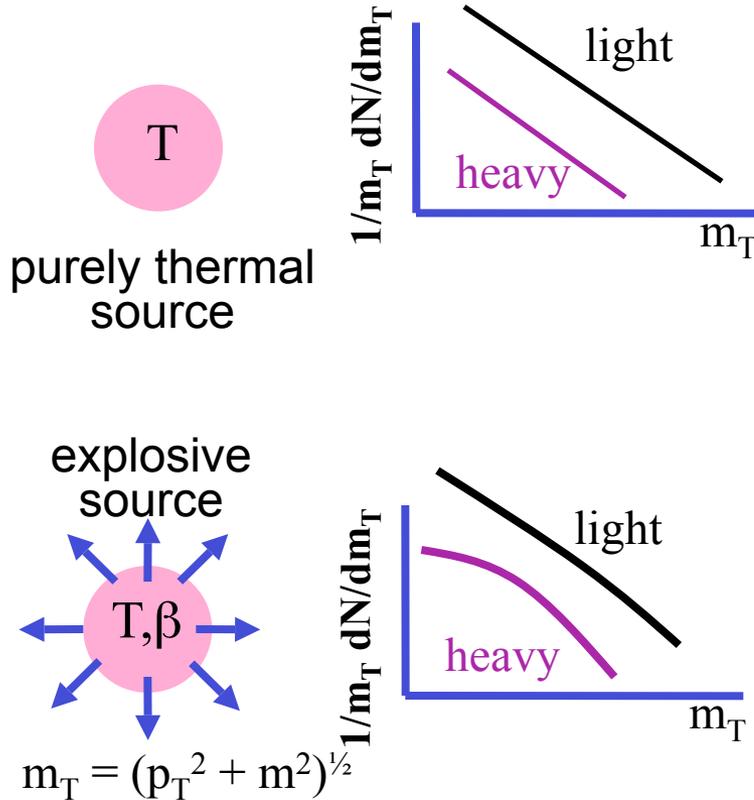


Strong collective radial expansion



- Different spectral shapes for particles of differing mass
 → strong **collective radial flow**

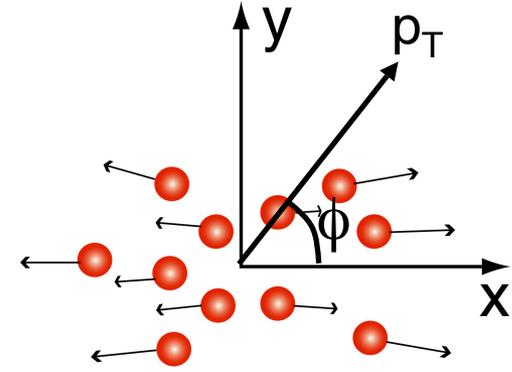
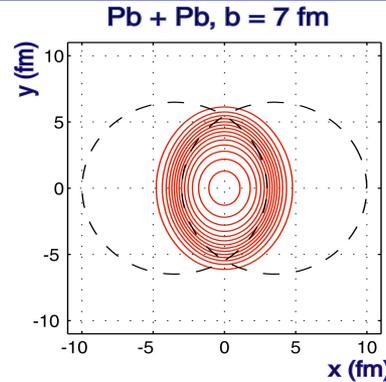
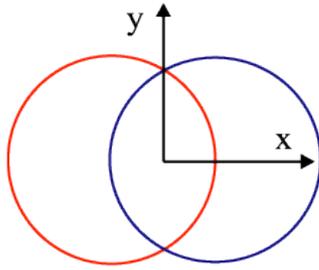
Strong collective radial expansion



- Different spectral shapes for particles of differing mass
→ strong collective radial flow

Good agreement with hydrodynamic prediction for soft EOS (QGP+HG)

Anisotropic/Elliptic flow



Almond shape overlap
region in **coordinate space**



**Interactions/
Rescattering**



**Anisotropy in
momentum
space**

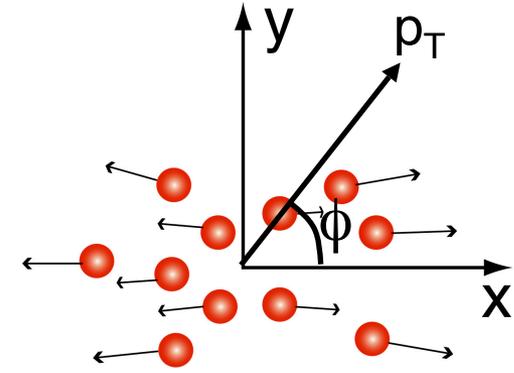
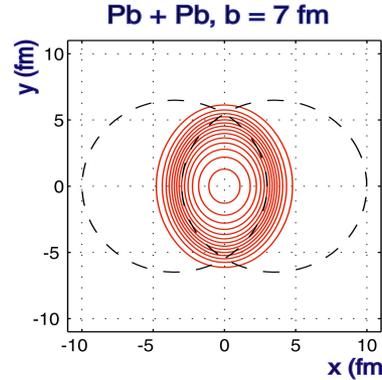
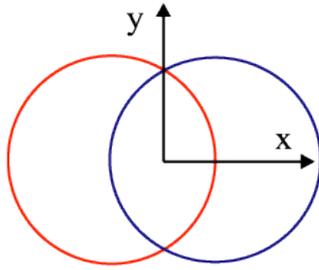
$$dN/d\phi \sim 1 + 2 v_2(p_T) \cos(2\phi) + \dots$$

$$\phi = \text{atan}(p_y/p_x)$$

$$v_2 = \langle \cos 2\phi \rangle$$

v_2 : 2nd harmonic Fourier coefficient in $dN/d\phi$ with respect to the reaction plane

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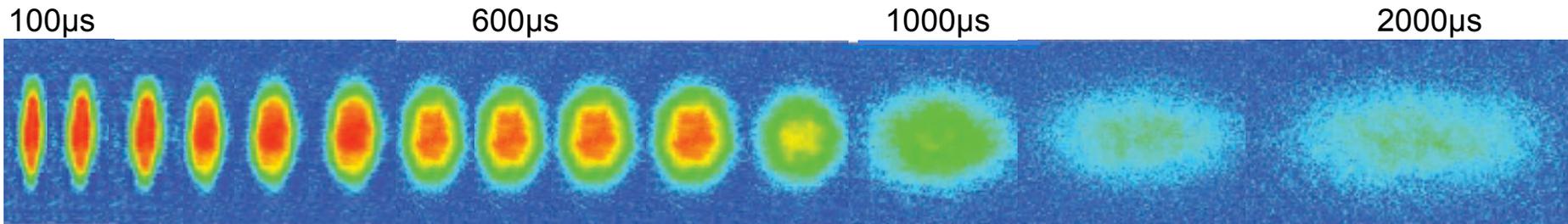
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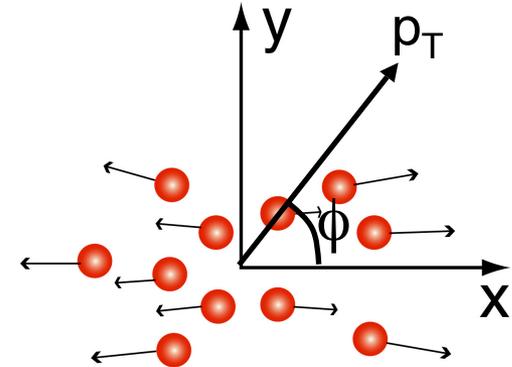
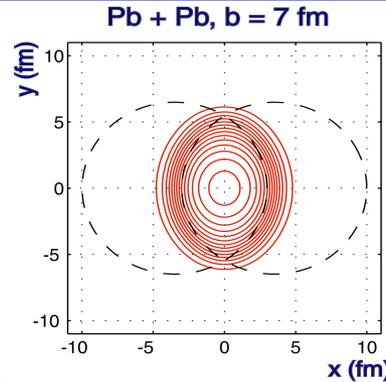
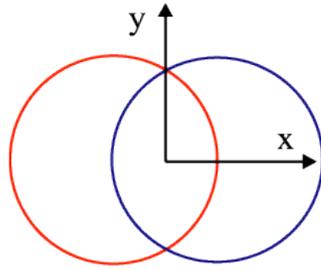
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Time

–M. Gehr, S. Granade, S. Hemmer, K. O'Hara, J. Thomas - **Science** 298 2179 (2002)

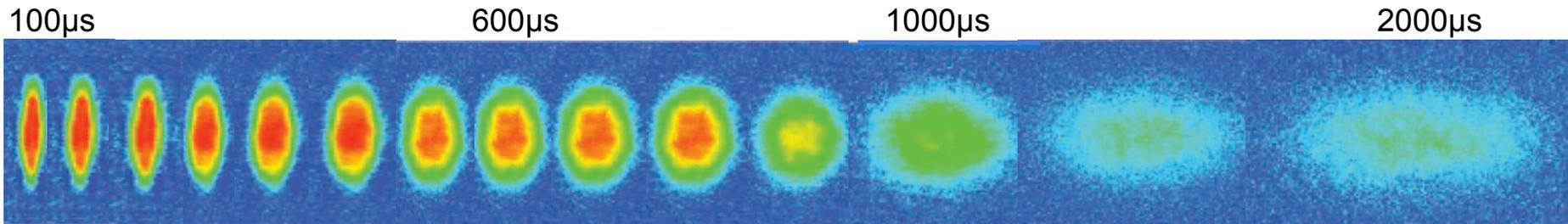
Anisotropic/Elliptic flow



Elliptic flow observable sensitive to early evolution of system

Mechanism is self-quenching

Large v_2 is an indication of **early** thermalization

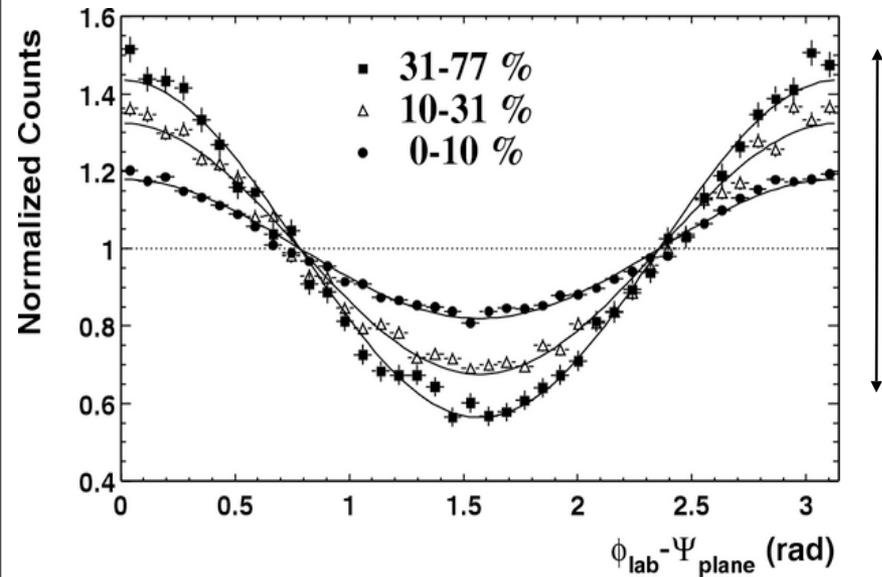


→ Time

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Elliptic flow

Distribution of particles with respect to event plane, $\phi-\psi$, $p_t > 2$ GeV; STAR PRL 90 (2003) 032301

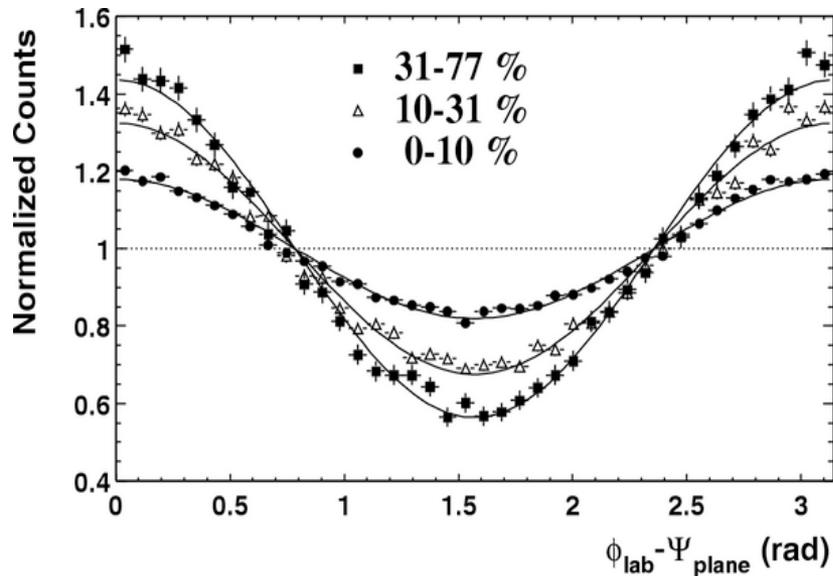


• Very strong elliptic flow → early equilibration

Factor 3:1 peak to valley

Elliptic flow

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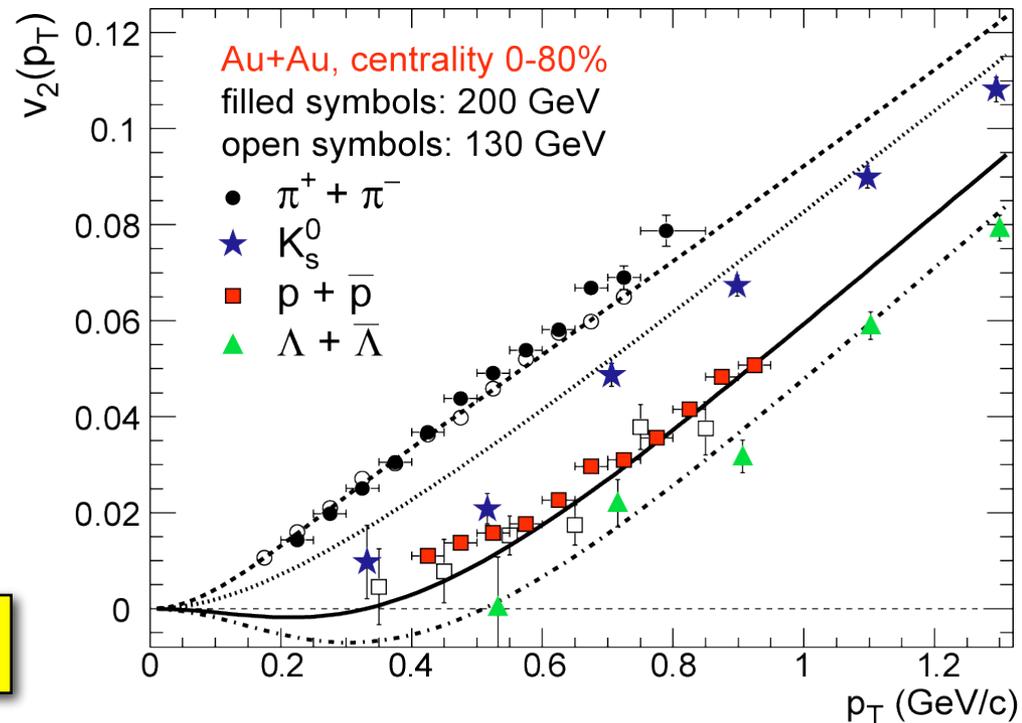


• Very strong elliptic flow \rightarrow early equilibration

Factor 3:1 peak to valley

• Pure hydrodynamical models including QGP phase describe elliptic and radial flow for many species

QGP \rightarrow almost perfect fluid



The constituents “flow”

- Elliptic flow is additive.
- If partons are flowing the *complicated* observed flow pattern in $v_2(p_T)$ for hadrons

$$\frac{d^2N}{dp_T d\phi} \propto 1 + 2 v_2(p_T) \cos(2\phi)$$

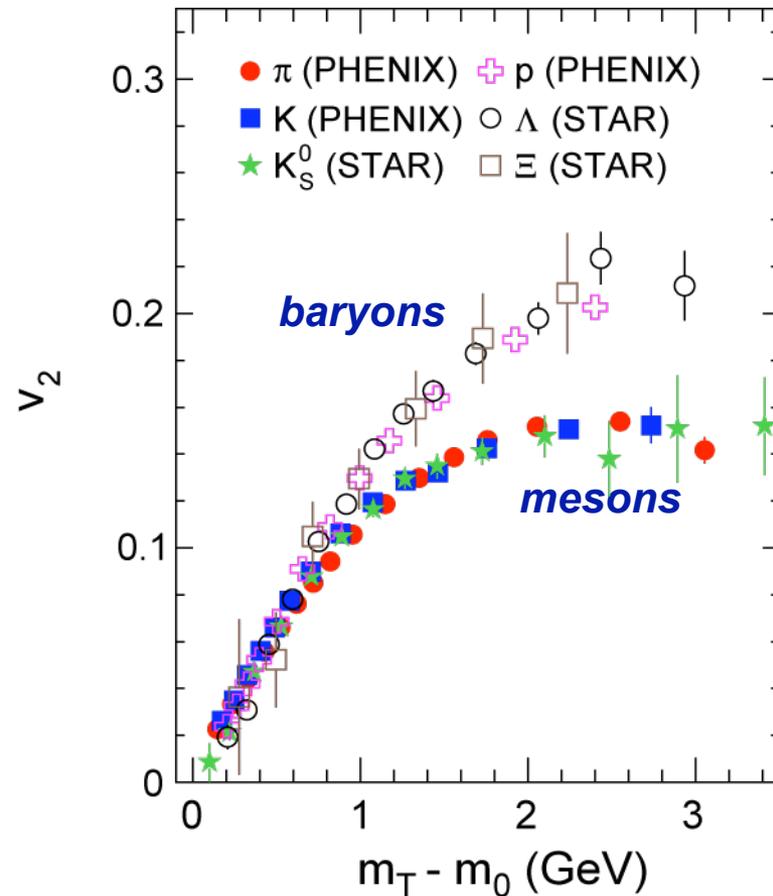
should become *simple* at the quark level

$$p_T \rightarrow p_T / n$$

$$v_2 \rightarrow v_2 / n ,$$

$n = (2, 3)$ for (meson, baryon)

$$m_T = \sqrt{p_T^2 + m_0^2}$$



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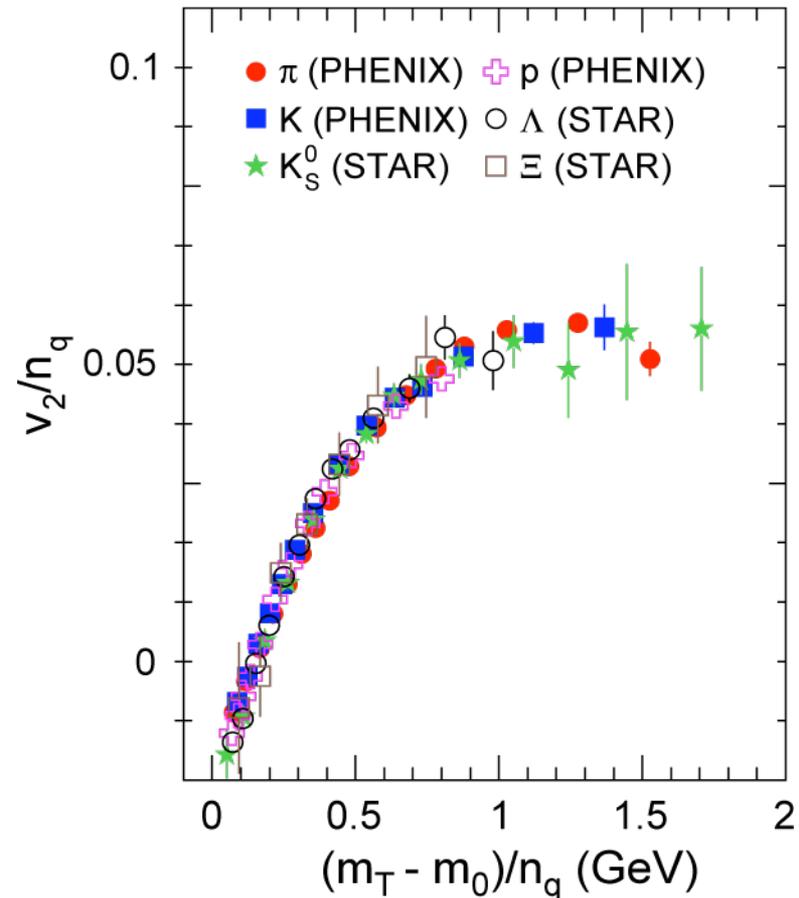
$$v_2 \rightarrow v_2 / n ,$$

$n = (2, 3)$ for (meson, baryon)

Works for $p, \pi, K_S^0, \Lambda, \Xi$..

$$v_2^s \sim v_2^{u,d} \sim 7\%$$

$$m_T = \sqrt{p_T^2 + m_0^2}$$



Constituents of QGP are partons

Summary of what we learned so far

- Energy density in the collision region is way above that where hadrons can exist
- The initial temperature of collision region is way above that where hadrons can exist
- The medium has quark and gluon degrees of freedom in initial stages

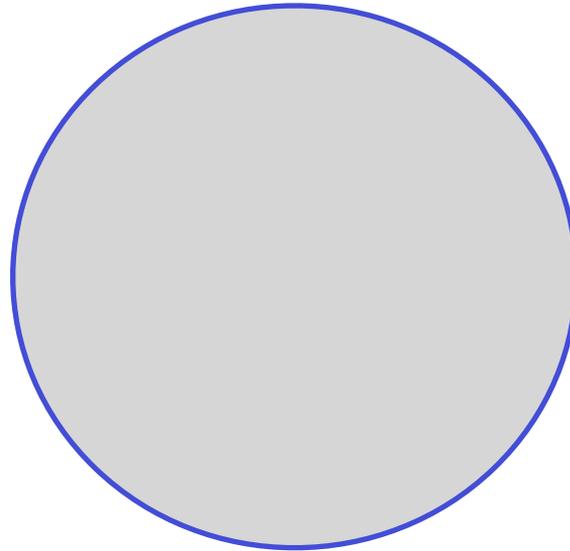
We have created a new state of matter at RHIC
- the QGP

- The QGP is flowing like an almost “perfect” liquid

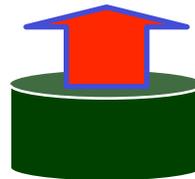
How to learn more about QGP?

Matter we want to study

Calibrated
LASER



Calibrated
Light Meter

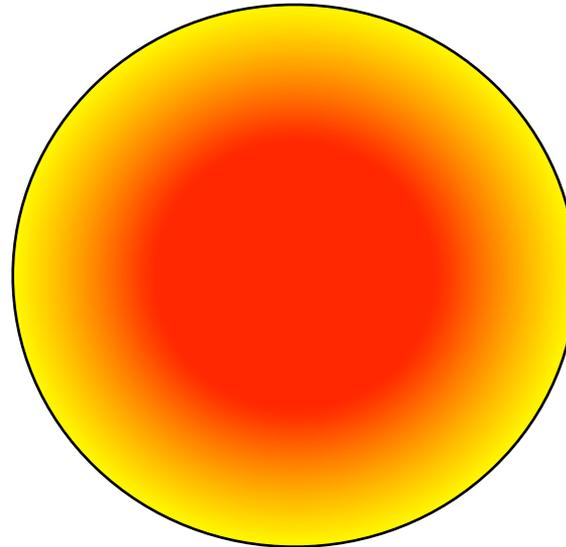


Calibrated
Heat Source

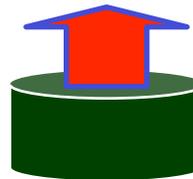
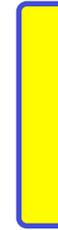
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Detectors



Energy released
in A+A collision
(27 TeV for Au+Au at RHIC)

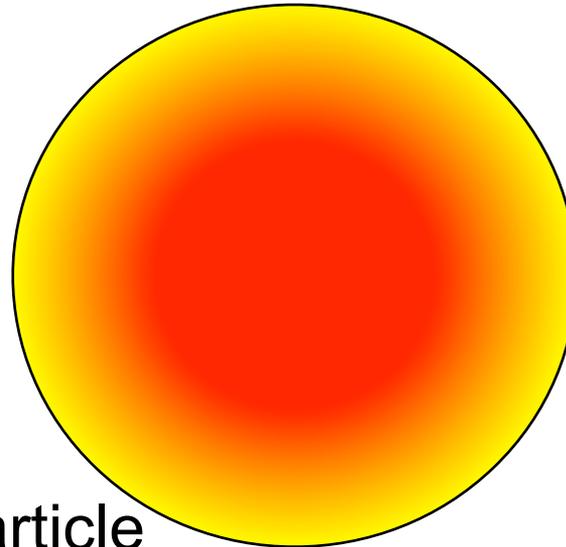
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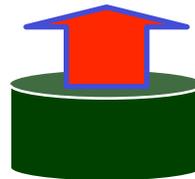
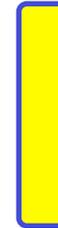
Hard Probes



- Photons
- Partons (q, g)
- High momentum particle



Detectors



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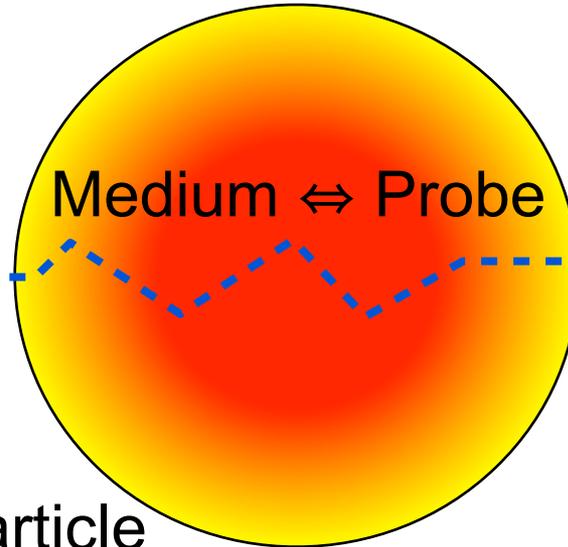
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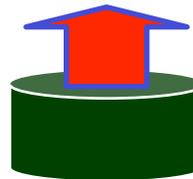
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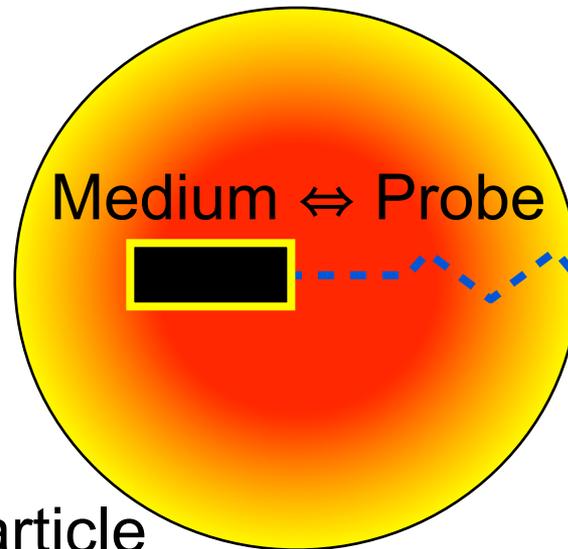
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Matter we want to study

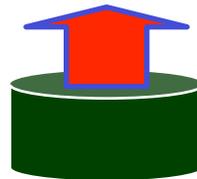
Hard Probes

Self-generated probes

- Photons
- Partons (q, g)
- High momentum particle



Detectors

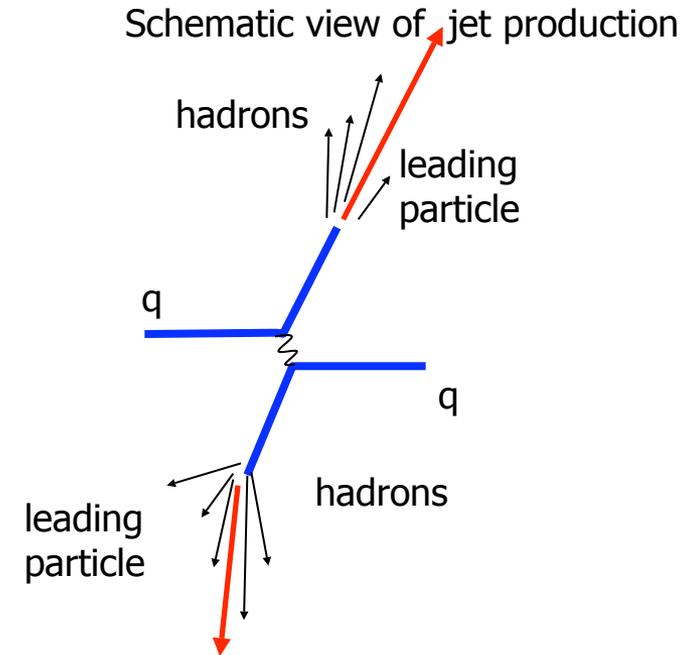


Energy released
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Using high momentum particles as probes

Early production in parton-parton scatterings with large Q^2 .

Direct interaction with partonic phases of the reaction



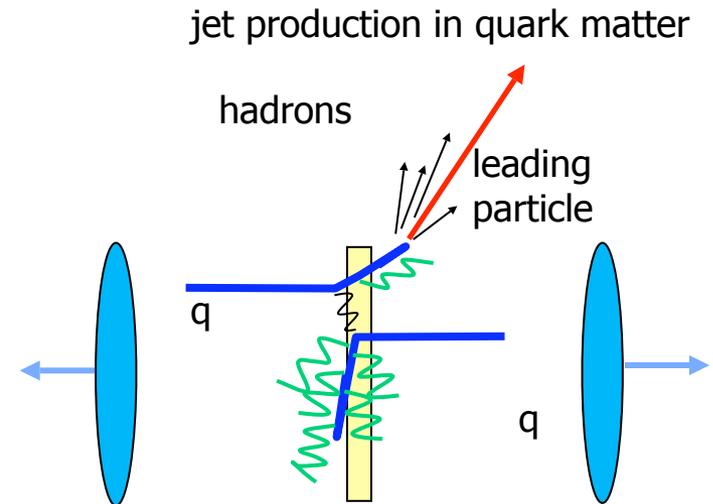
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Therefore use these high momentum products as probes at RHIC

- attenuation or absorption of high p_T hadrons



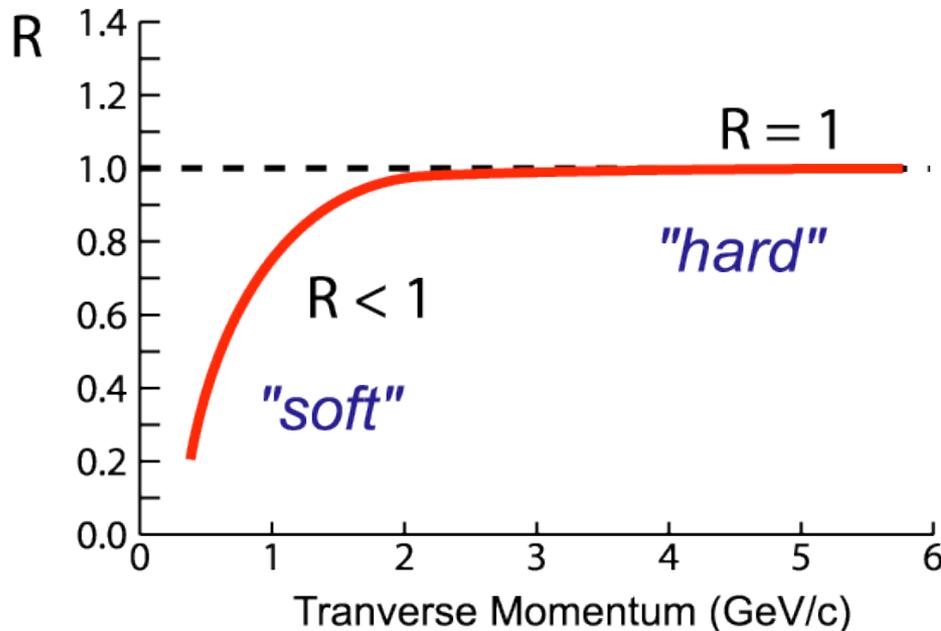
Looking for attenuation/absorption

Compare to p-p at same collision energy

Nuclear
Modification
Factor:

$$R_{AA}(p_T) = \frac{Yield(A+A)}{Yield(p+p) \times \langle N_{coll} \rangle}$$

Average number
of p-p collision
in A-A collision



No "Effect":

- $R < 1$ at small momenta - production from thermal bath
- $R = 1$ at higher momenta where hard processes dominate

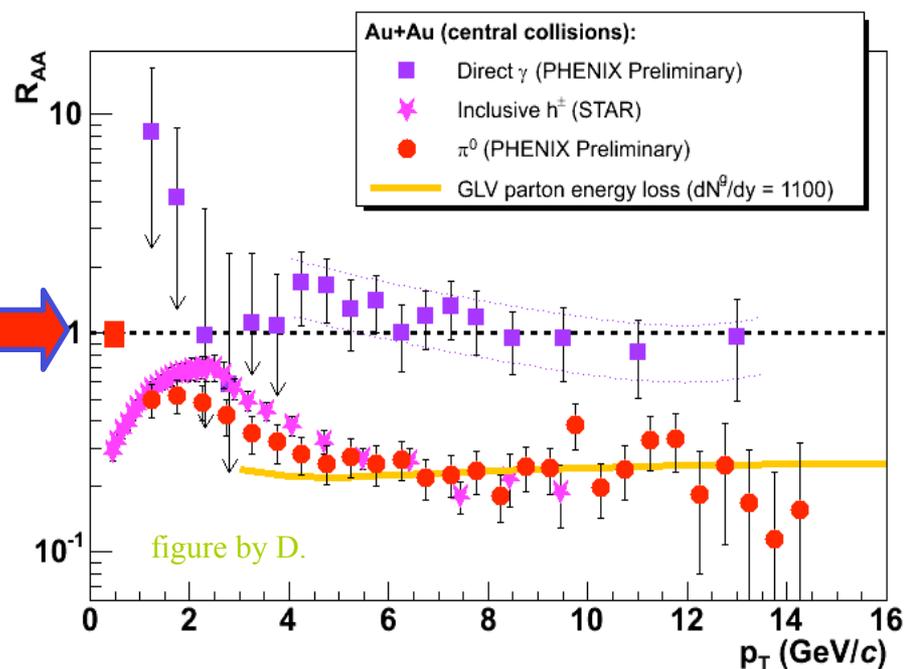
$R < 1$ at high p_T if QGP affecting parton's propagation

High- p_T suppression

Observations at RHIC:

1. Photons are **not** suppressed

- Good! γ don't interact with medium
- N_{coll} scaling works



High- p_T suppression

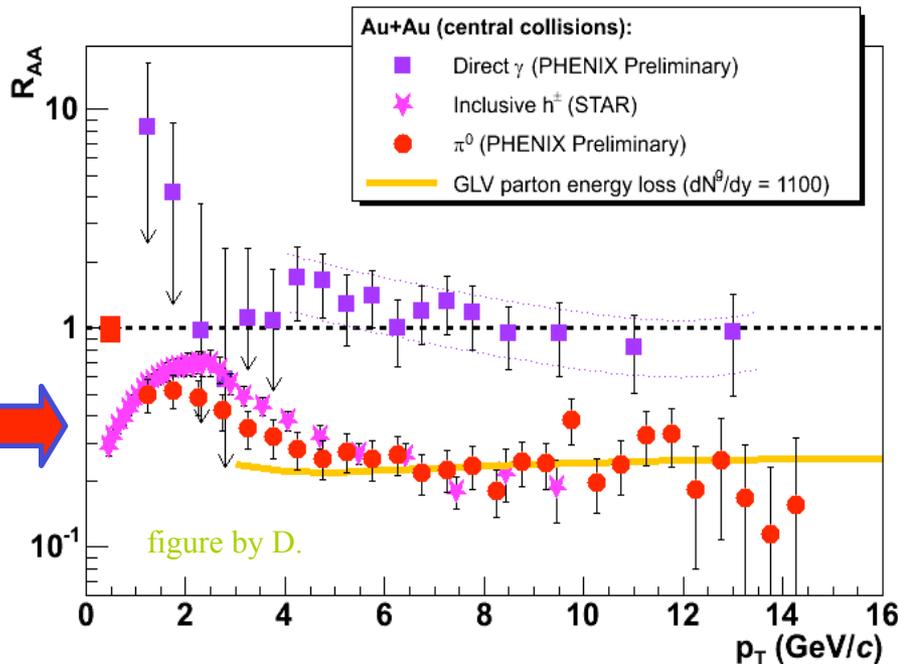
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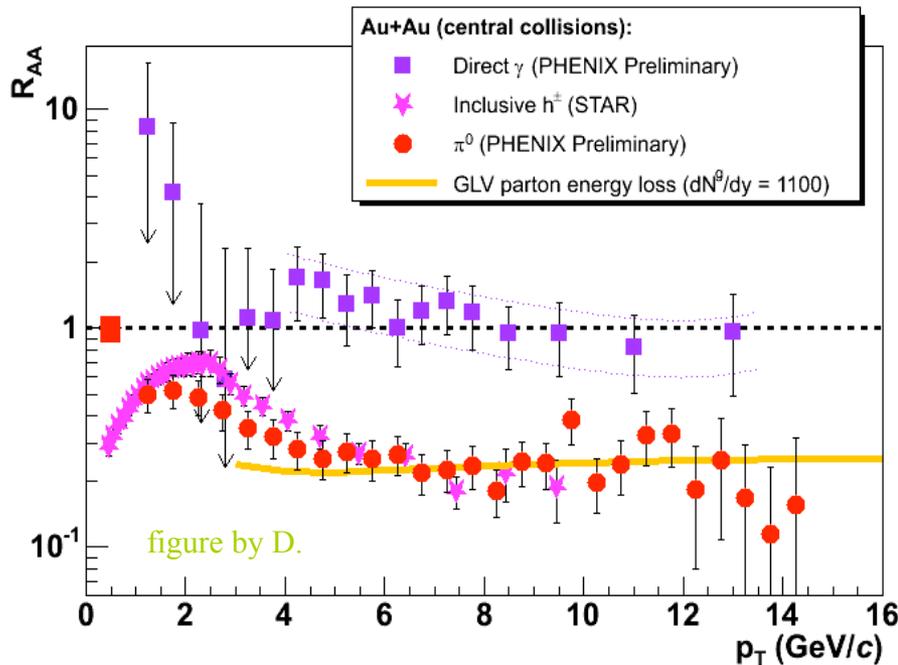
2. Hadrons are **suppressed** in central collisions

- Huge: factor 5



High- p_T suppression

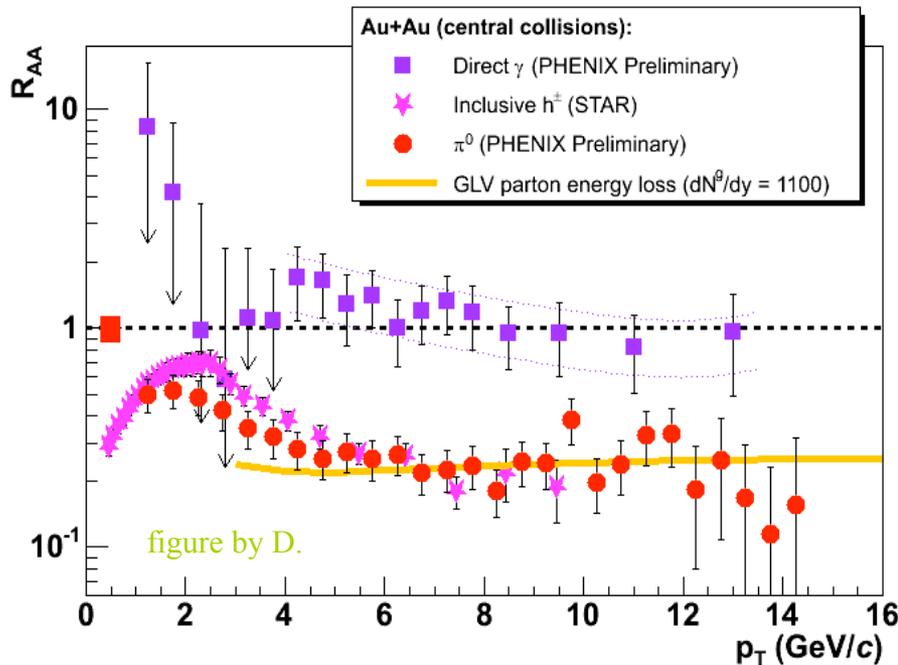
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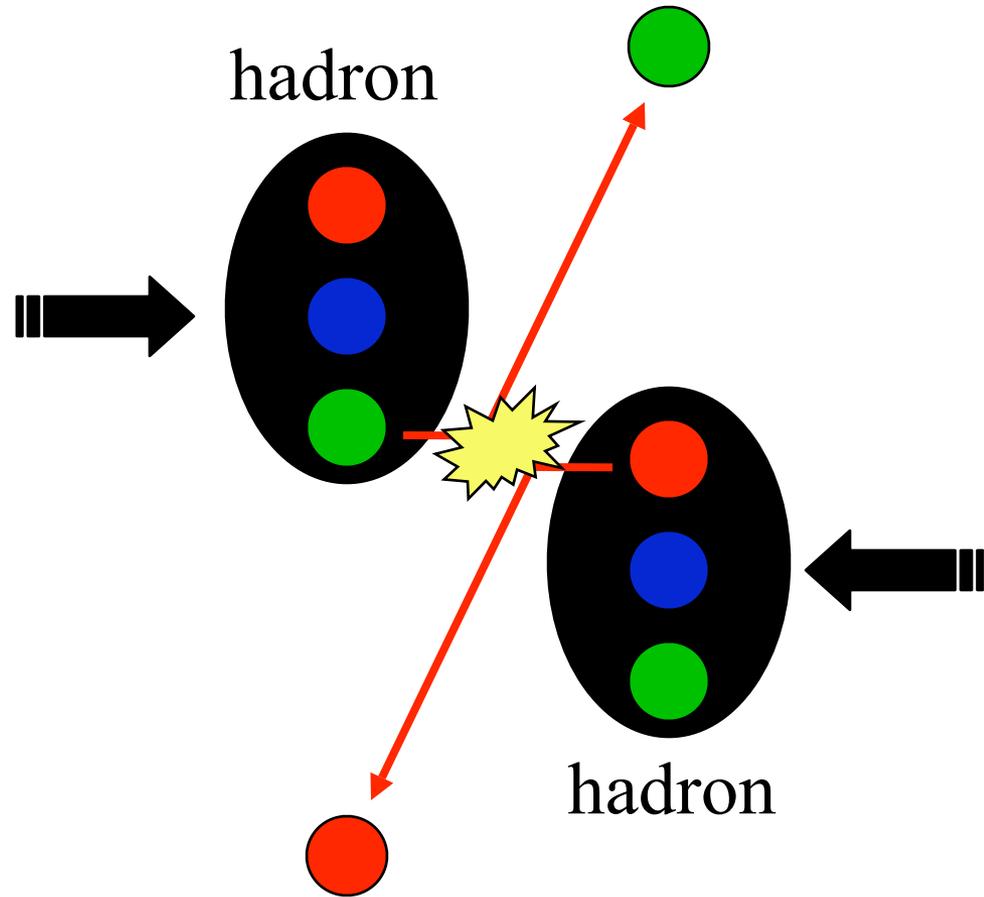
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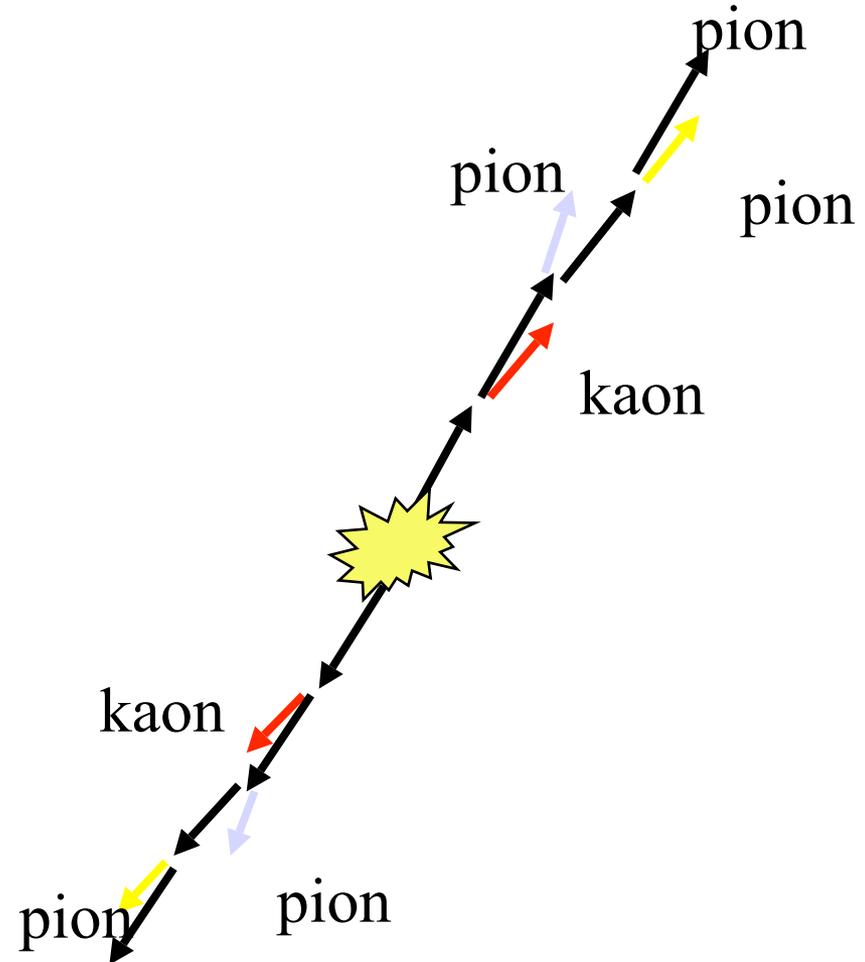
sQGP - strongly coupled - colored objects suffer large energy loss

Definition of a “jet”



Definition of a “jet”

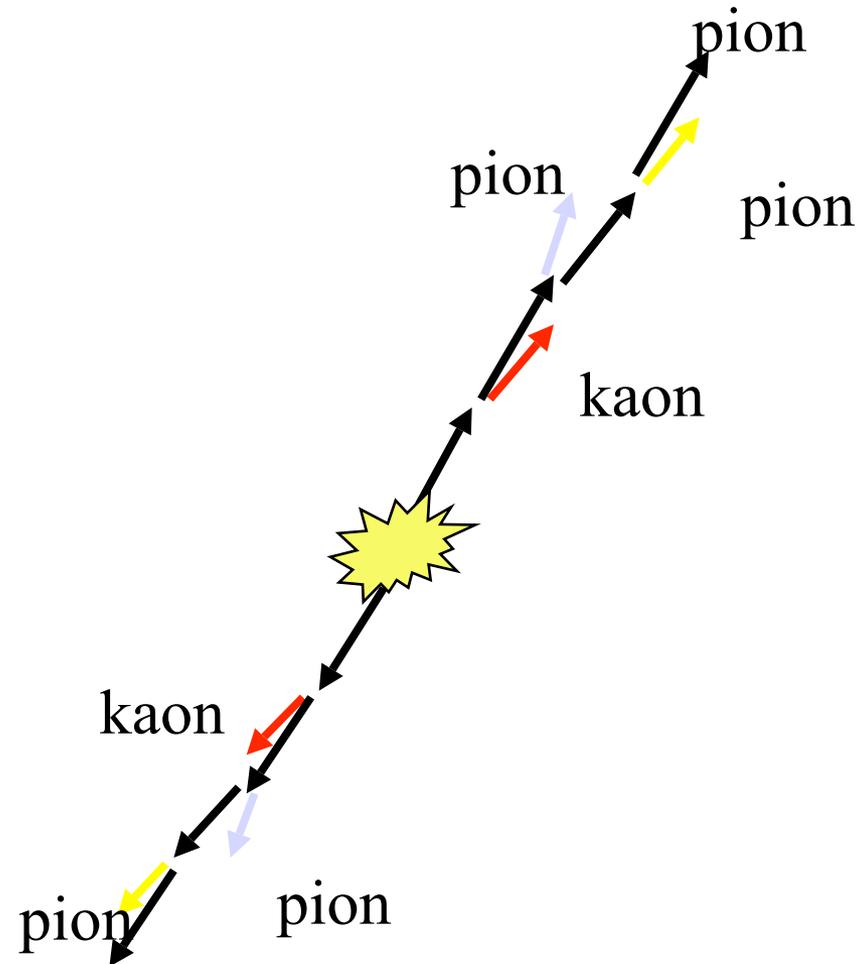
- Use Jets – the possible for “knock-on” collisions of partons
- The fragmented “bits” appear as “normal” subatomic particles
pions, kaons, etc
- Seen in high-energy physics experiments since mid-1970’s



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Jets commonly come in pairs



Using jets to study the QGP properties

A case study: opacity of fog



- “is this thing on?”

Using jets to study the QGP properties

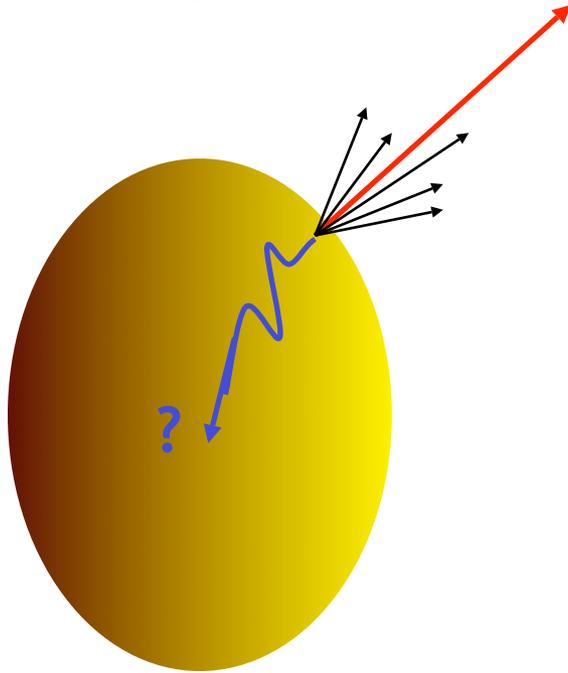
A case study: opacity of fog



- “is this thing on?”
- First beam - least know the source is on.
- Second beam intensity tells you a lot about matter passed through

Using jets to study the QGP properties

A case study: opacity of fog



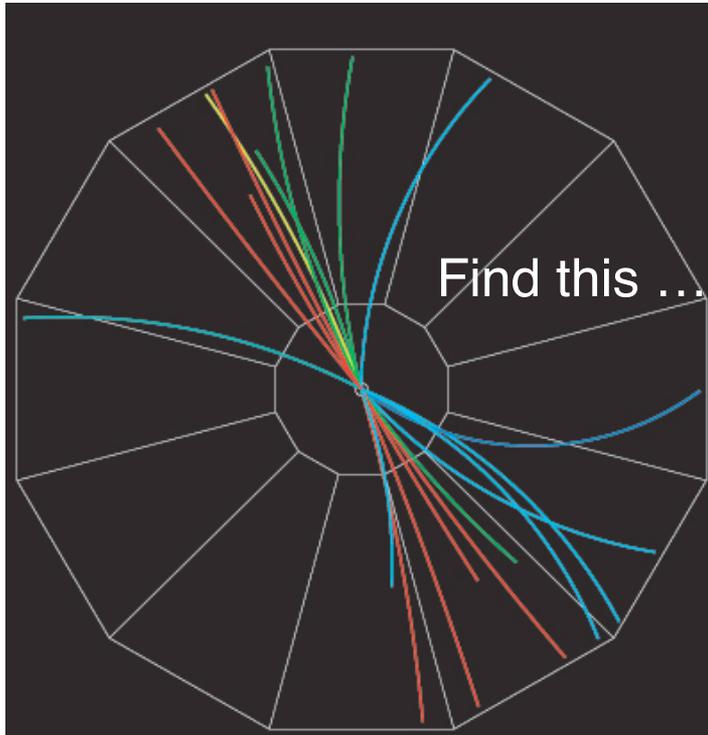
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Predictions

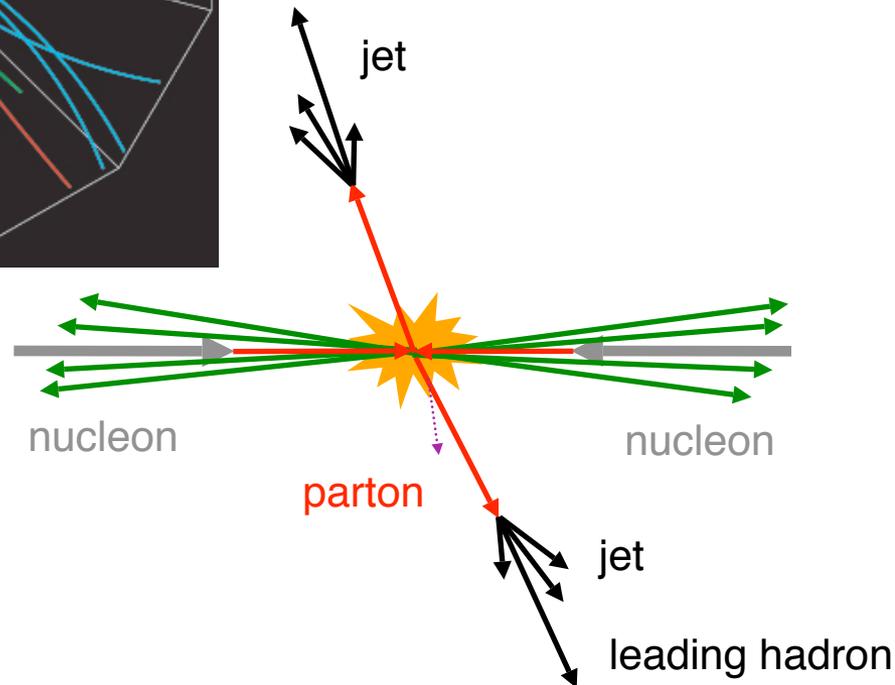
QGP: “backwards” jet will be absorbed by medium

Hadron gas: “backwards” jet be less affected by medium

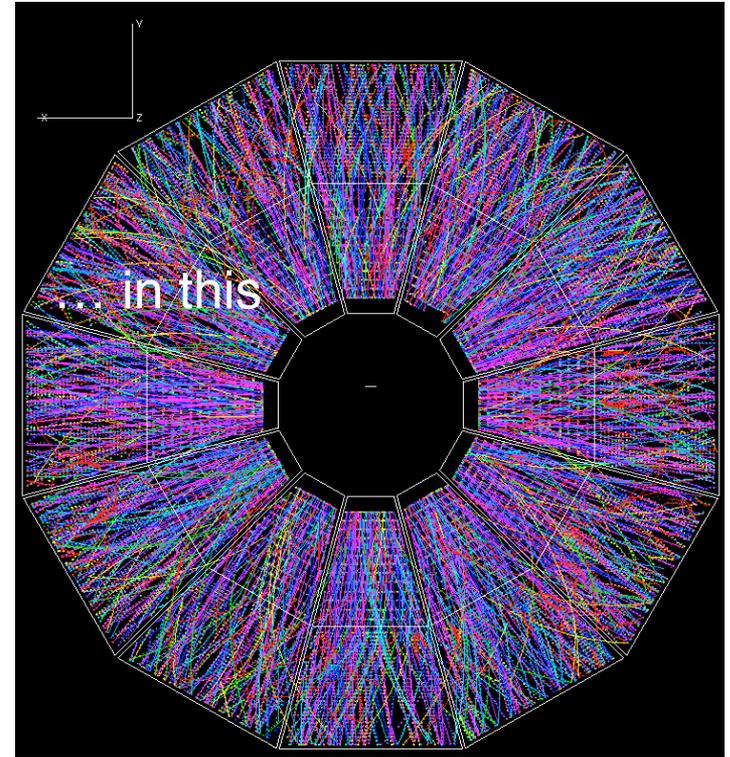
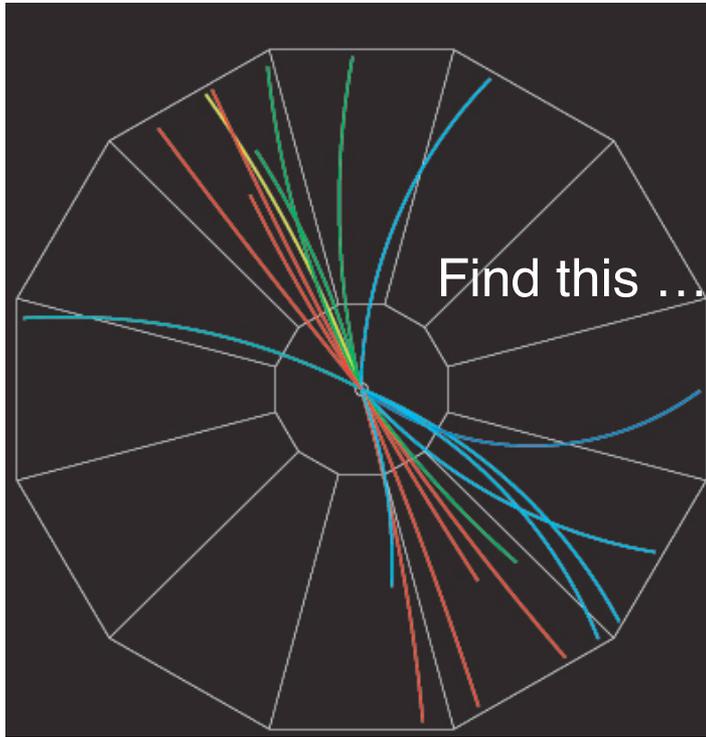
Finding a jet in a Au-Au event



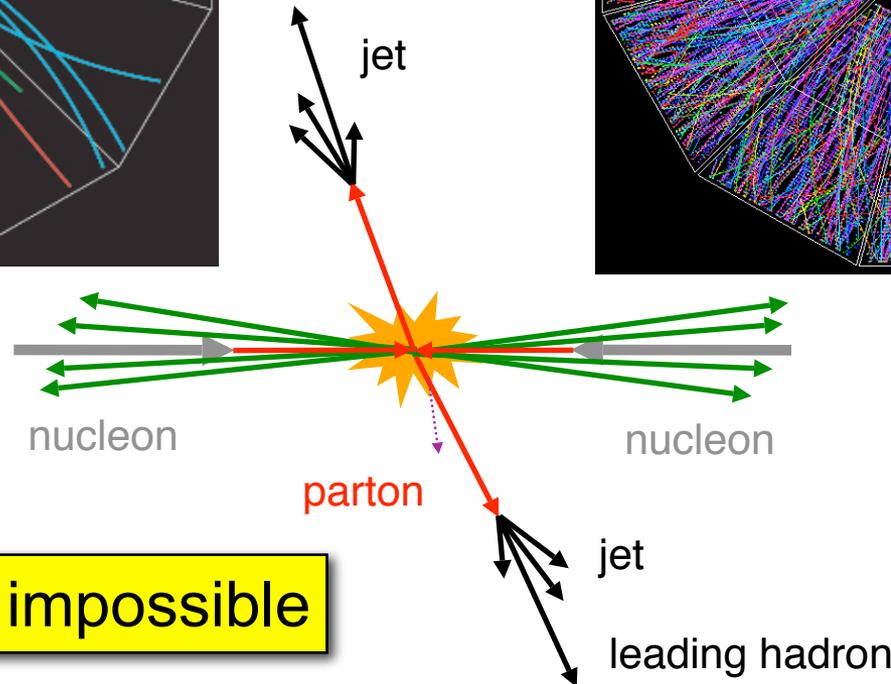
$p+p \rightarrow \text{jet}+\text{jet}$
(STAR@RHIC)



Finding a jet in a Au-Au event



$p+p \rightarrow \text{jet}+\text{jet}$
(STAR@RHIC)



$\text{Au}+\text{Au} \rightarrow ???$
(STAR@RHIC)

Seems almost impossible

How to find a jet? - an algorithm

- How to locate the running of the bulls in Pamplona, Spain:



How to find a jet? - an algorithm

- How to locate the running of the bulls in Pamplona, Spain:
 - start by finding one fast moving bull



How to find a jet? - an algorithm

- How to locate the running of the bulls in Pamplona, Spain:
 - start by finding one fast moving bull
 - look others moving in roughly the same direction - where there's one bull there's usually another



How to find a jet? - an algorithm

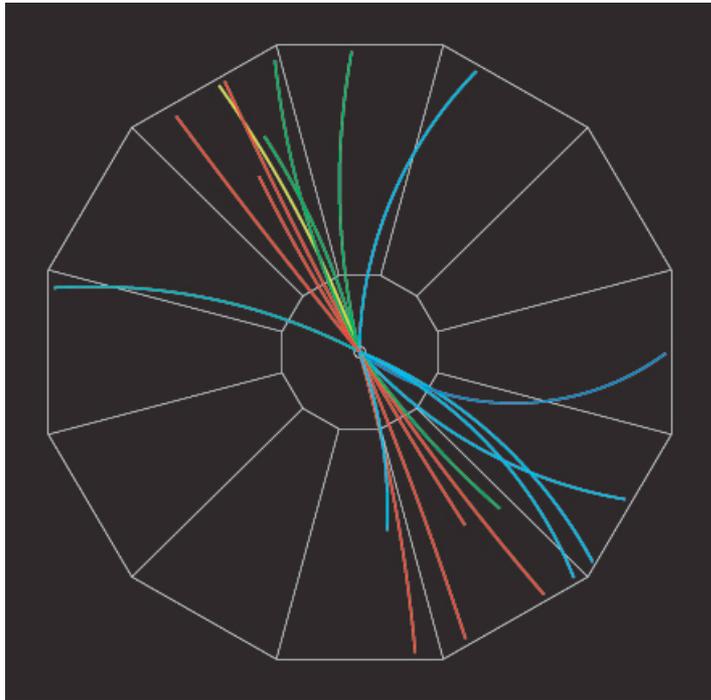
- How to locate the running of the bulls in Pamplona, Spain:
 - start by finding one fast moving bull
 - look others moving in roughly the same direction - where there's one bull there's usually another
 - if the bull density is high, you often find many people moving in opposite direction



Jet finding is now simple: just replace “bull” by “particle”

Jets in Au-Au collisions!

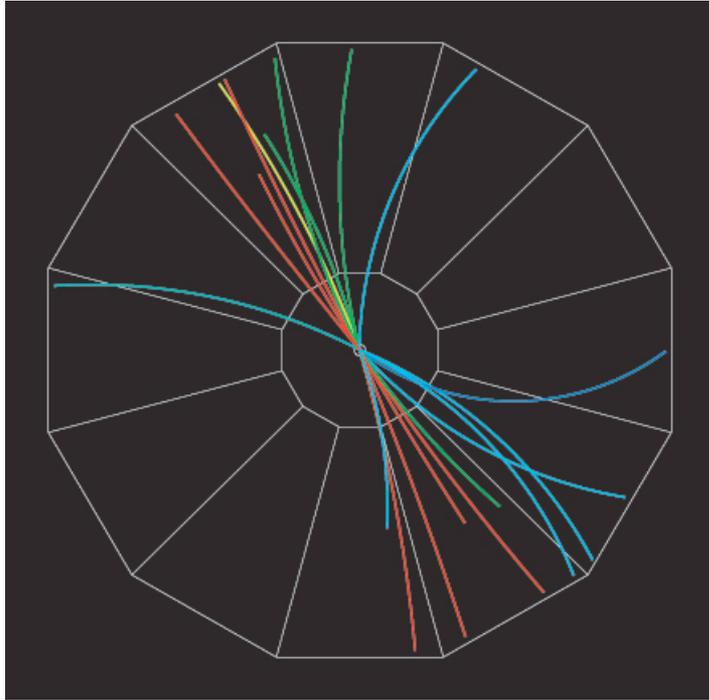
$p+p \rightarrow \text{dijet}$



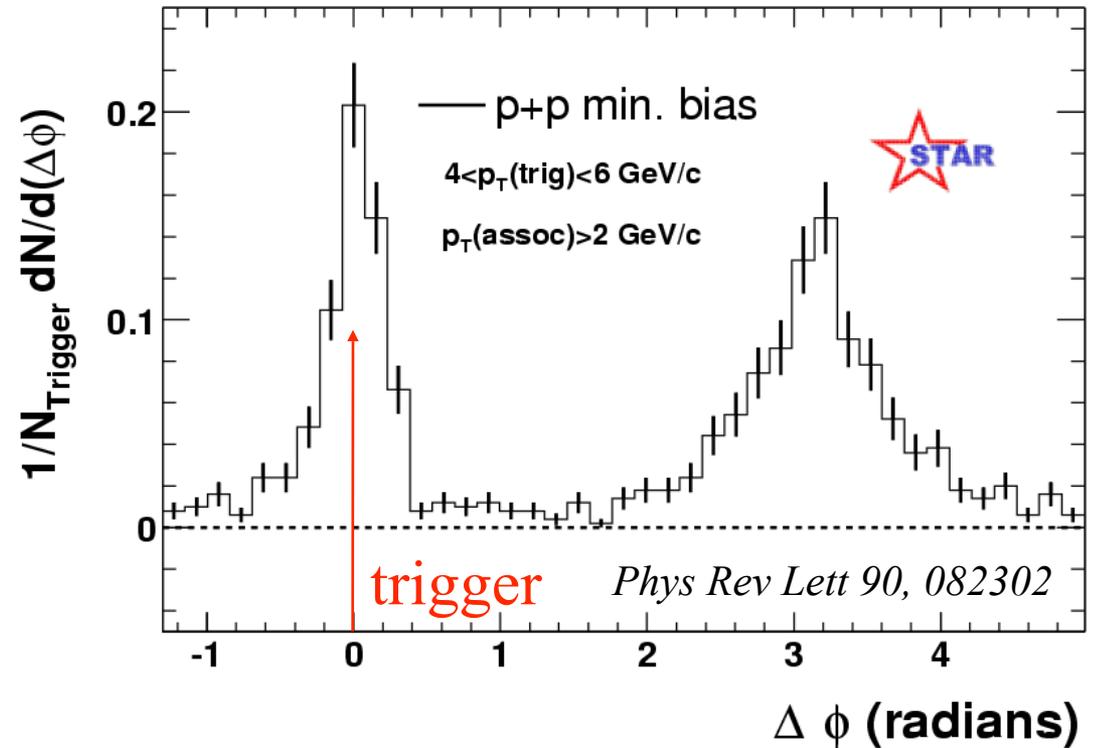
- Trigger: highest p_T track
- $\Delta\phi$ distribution:

Jets in Au-Au collisions!

$p+p \rightarrow \text{dijet}$

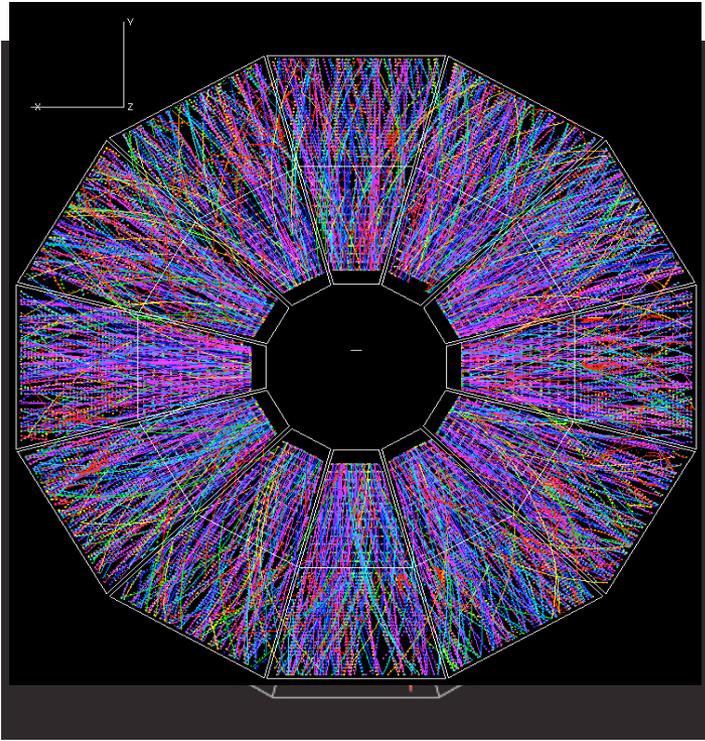


min. bias $p+p$ collisions

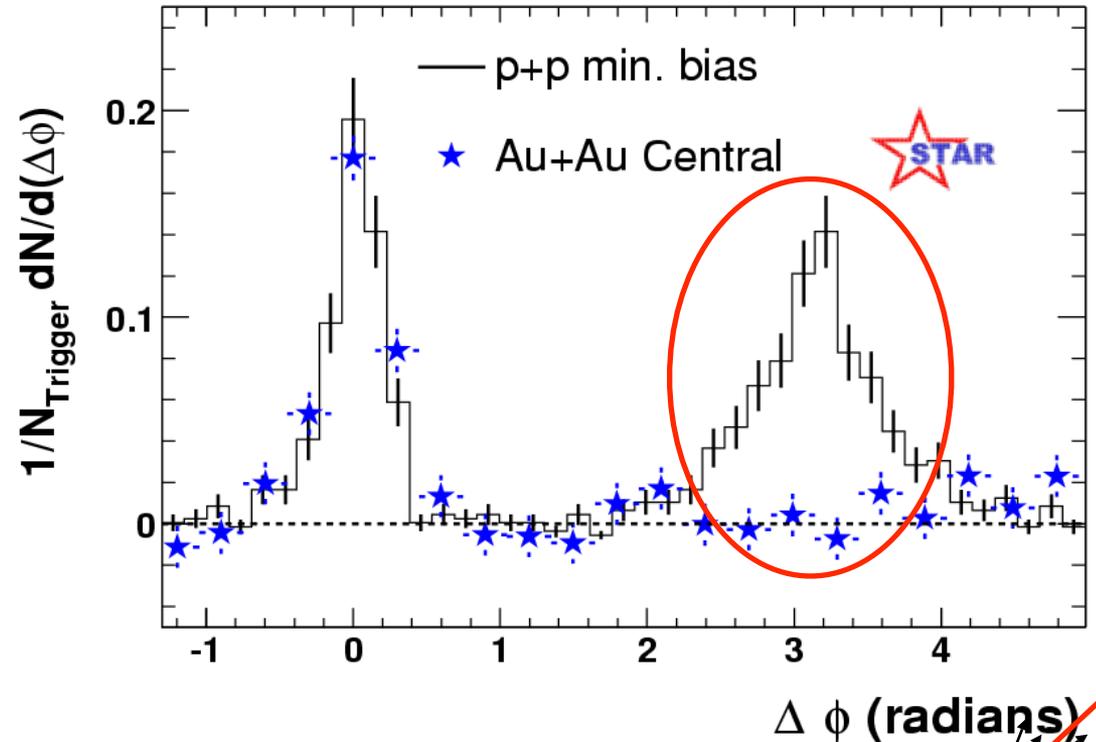


Jets in Au-Au collisions!

$p+p \rightarrow \text{dijet}$

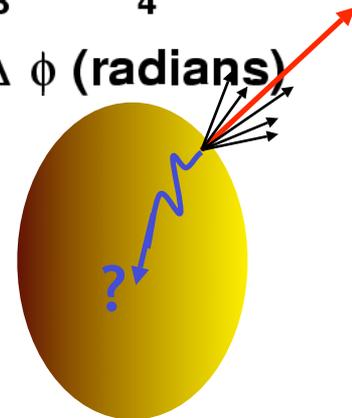


central Au+Au collisions



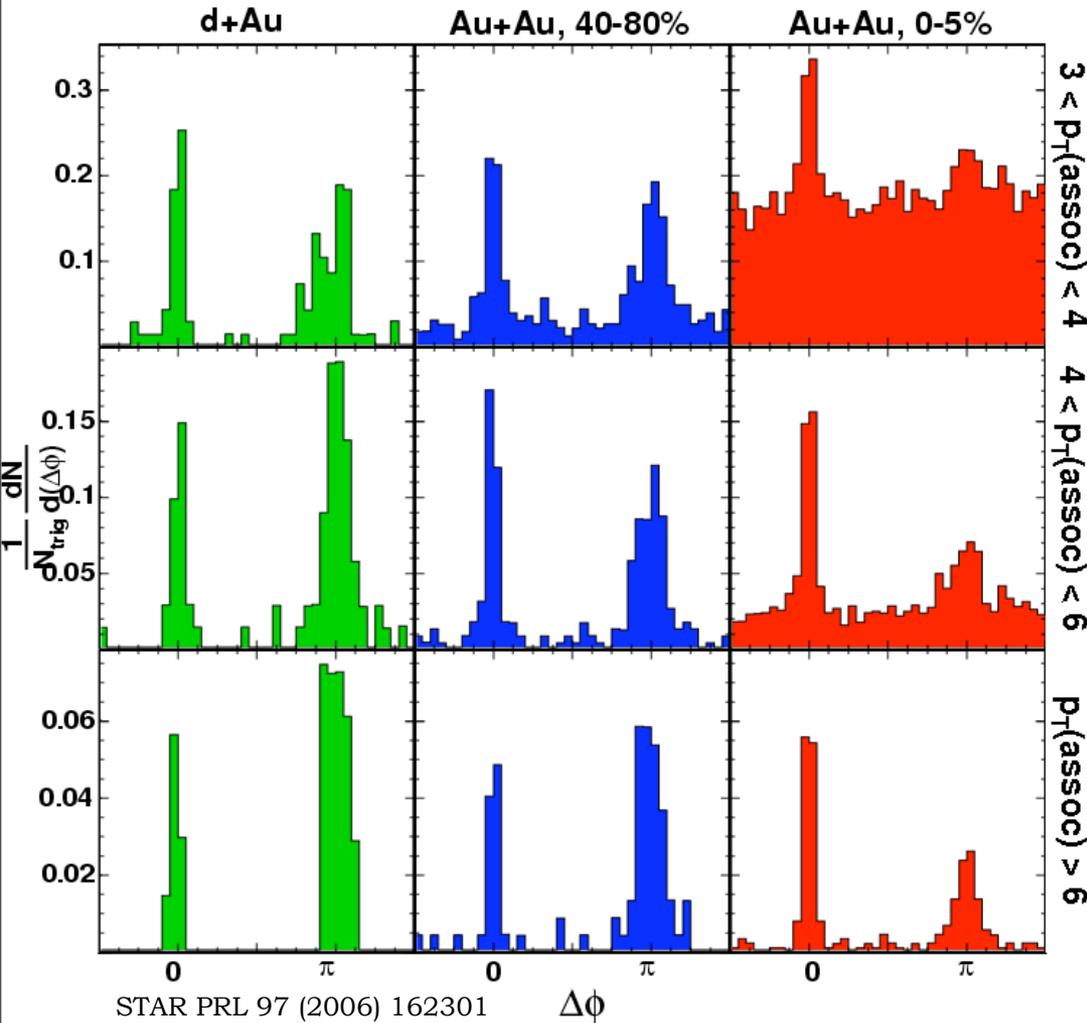
$\Delta\phi \approx 0$: central Au+Au similar to p+p

$\Delta\phi \approx \pi$: strong suppression of back-to-back correlations in central Au+Au



Observation of “Punch through”

$$8 < p_T^{\text{trig}} < 15 \text{ GeV}/c$$

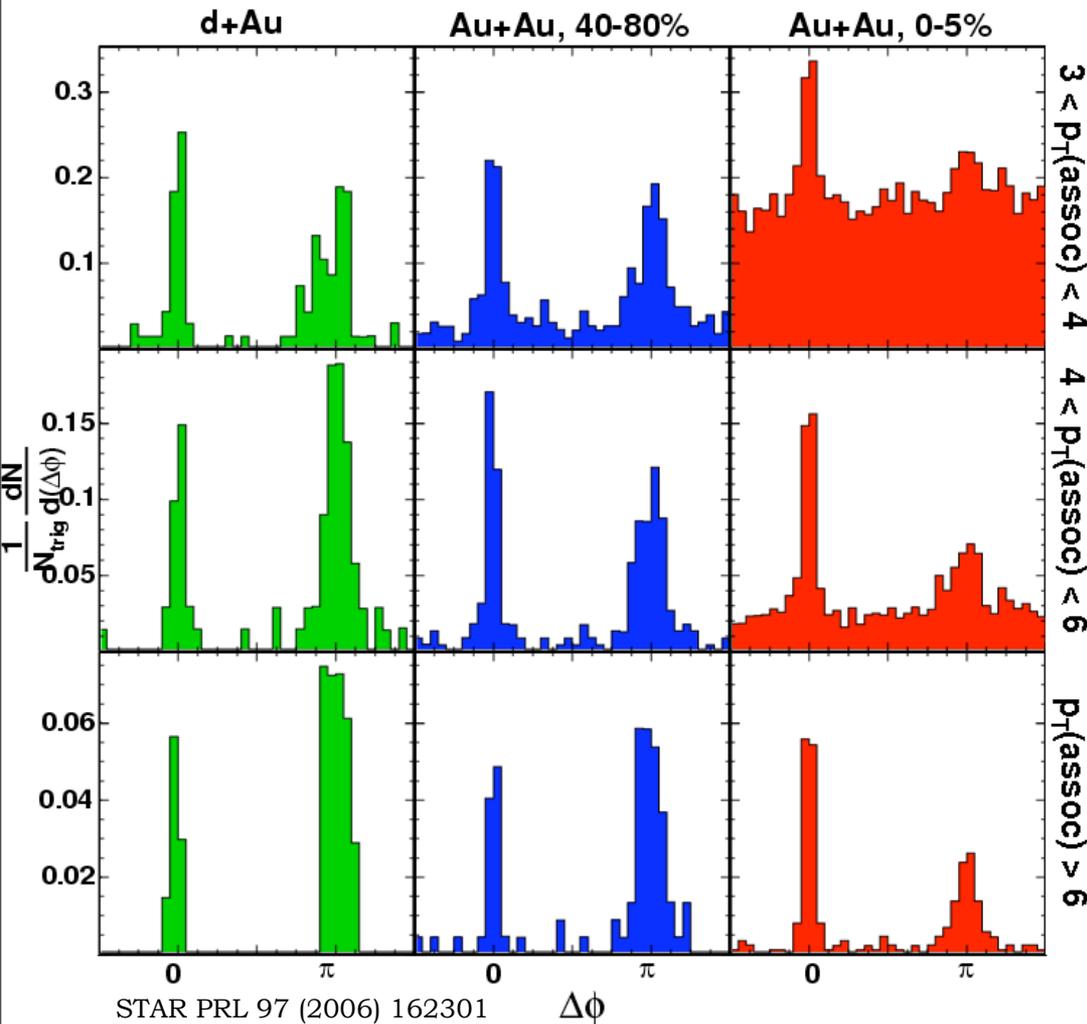


If use high- p_T triggers:

- Away-side peak re-emerges
- Smaller in Au-Au than d-Au
- Virtually no background

Observation of “Punch through”

$$8 < p_T^{\text{trig}} < 15 \text{ GeV}/c$$



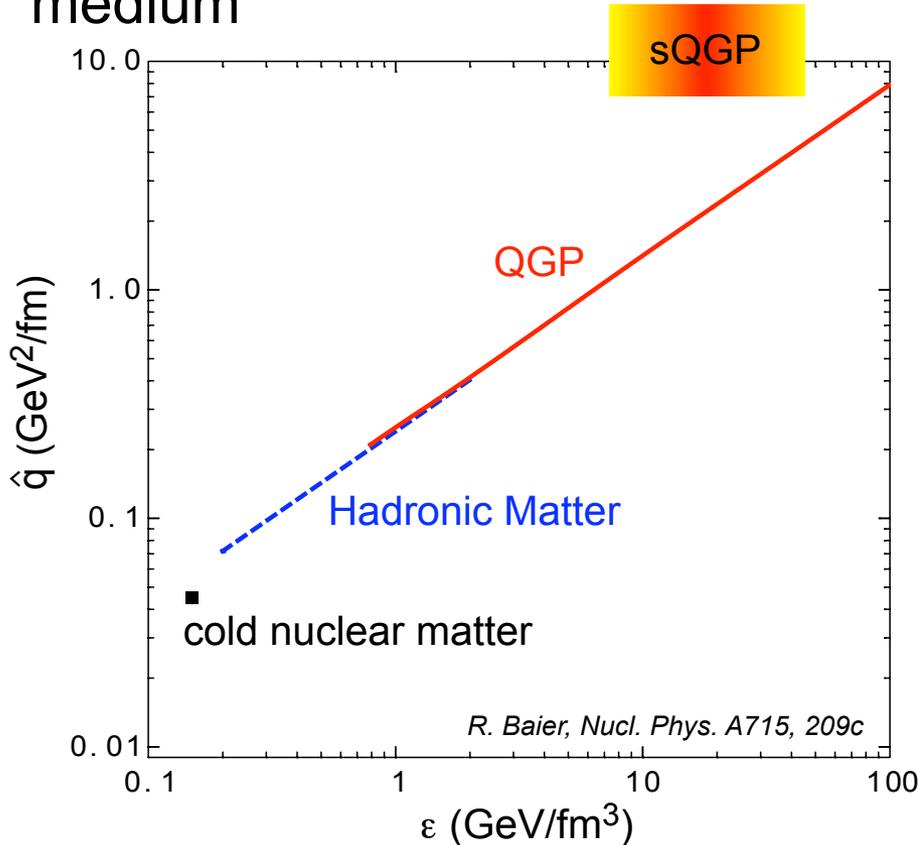
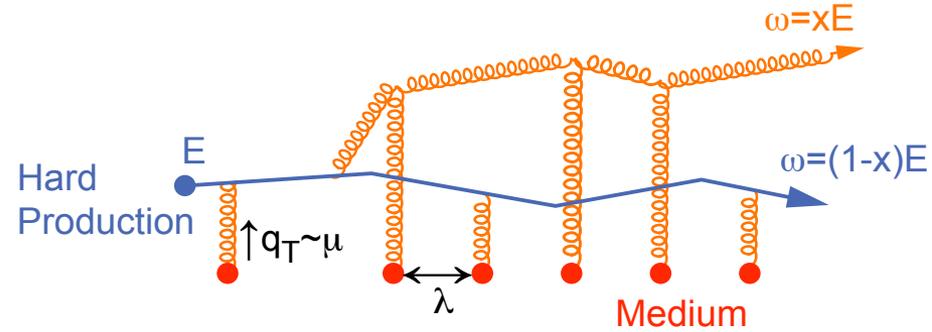
If use high- p_T triggers:

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High energy jets
“punch through”
the medium.

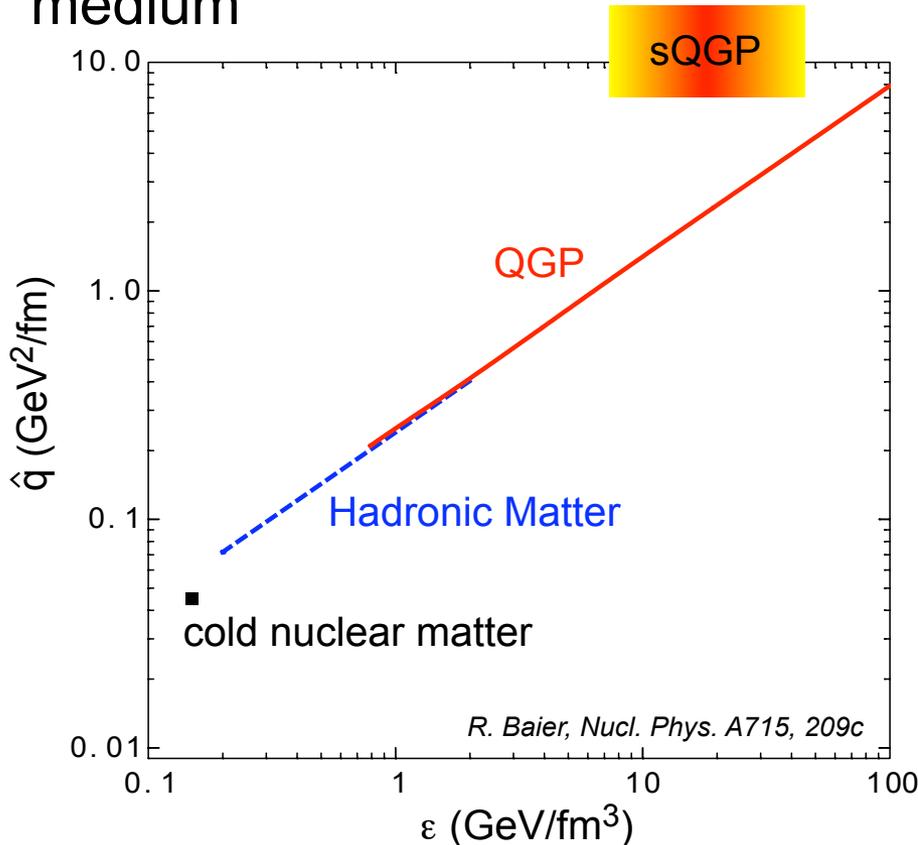
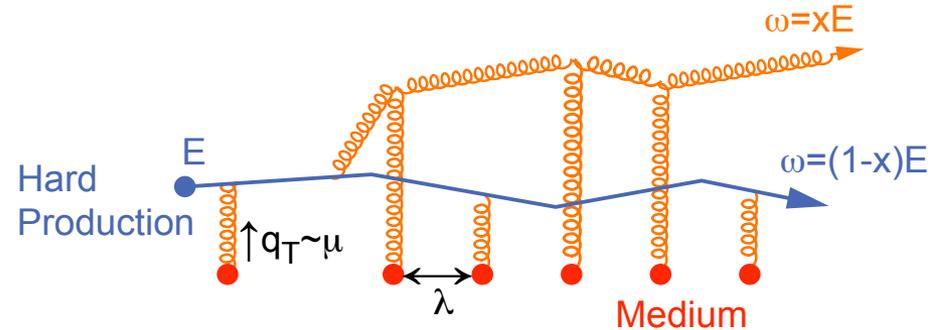
Interpretation

Gluon radiation: Multiple final-state gluon radiation off the produced hard parton induced by the traversed dense colored medium



Interpretation

Gluon radiation: Multiple final-state gluon radiation off the produced hard parton induced by the traversed dense colored medium



- Mean parton energy loss \propto medium properties:

- $\Delta E_{\text{loss}} \sim \rho_{\text{gluon}}$ (gluon density)
- $\Delta E_{\text{loss}} \sim \Delta L^2$ (medium length)
- $\Rightarrow \sim \Delta L$ with expansion

- Characterization of medium

- transport coefficient \hat{q}
- is $\langle p_T^2 \rangle$ transferred from the medium to a hard gluon per unit path length

$$\hat{q} \sim 5-10 \text{ GeV}^2/\text{fm}$$

Summary

The matter we create at RHIC is a **sQGP** it is

fantastically hot

and has an

incredible energy density.

It

exists for only an instant

yet shows

many signs of being in equilibrium.

It flows like a

nearly “perfect” fluid

and appears to have

quark and gluon degrees of freedom

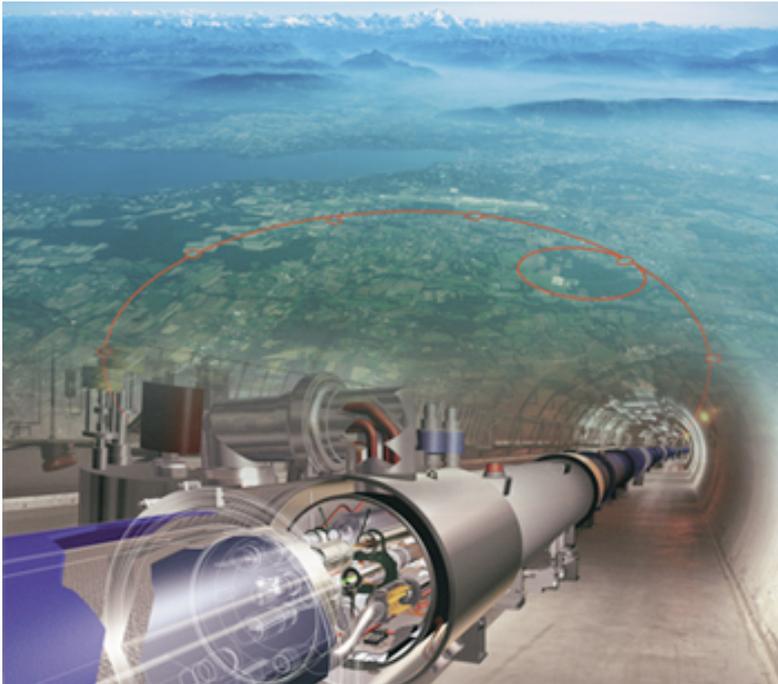
which causes

significant energy loss to partons passing through

The next energy frontier

The Large Hadron Collider (LHC) at CERN will be commissioned in 2008 with over an order of magnitude higher energy than at RHIC.

Instead of 40 TeV, 1000 TeV !



3 experiments with
dedicated heavy-ion
experiments

ALICE
ATLAS
CMS

sQGP: hotter, bigger, longer lived
more detailed measurements