

# Using particle jet correlations to probe the medium at RHIC

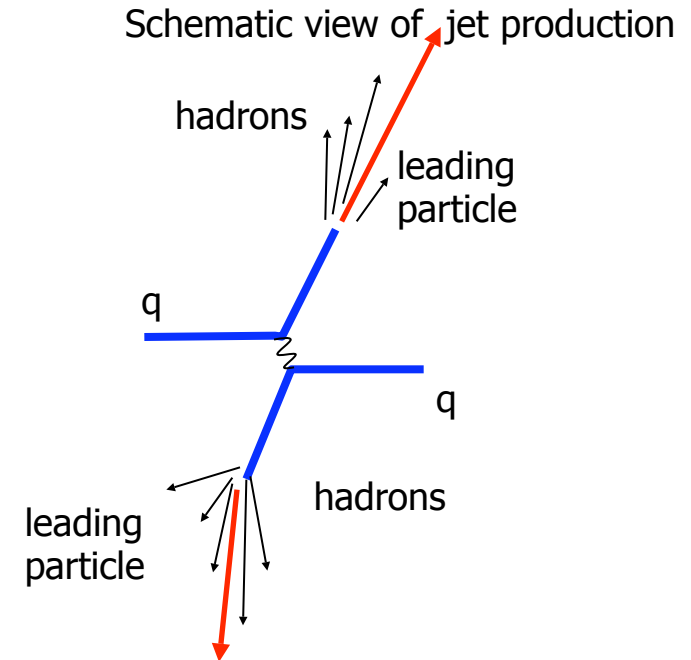
*Helen Caines - Yale University*

DIS 2008 – University College London  
April 7<sup>th</sup>-11<sup>th</sup> 2008



# Using “jets” as probes

- **Early production** in parton-parton scatterings with large  $Q^2$ .
- **Direct interaction with partonic phases** of the reaction

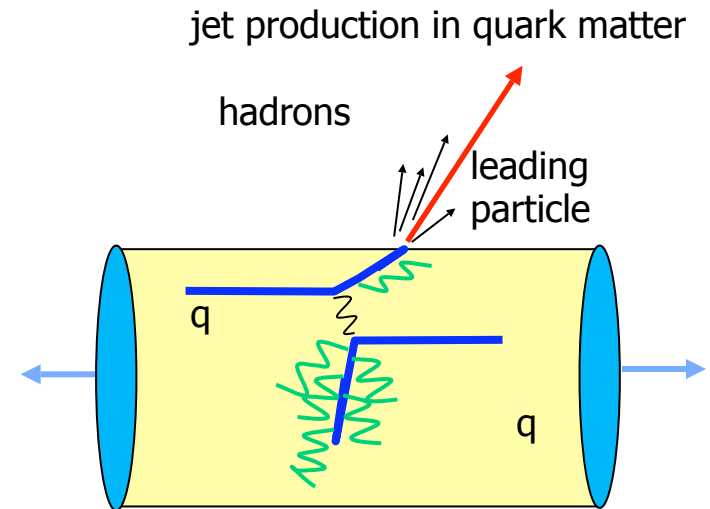


# Using “jets” as probes

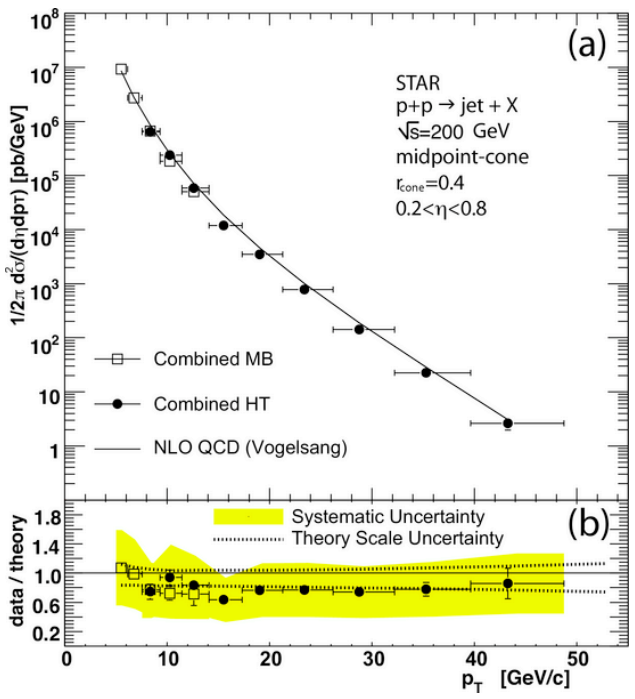
- **Early production** in parton-parton scatterings with large  $Q^2$ .
- **Direct interaction with partonic phases** of the reaction

Therefore use “jets” as probes at RHIC

- ◆ attenuation or absorption of jets:  
“jet quenching”
- ◆ suppression of high  $p_T$  hadrons
- ◆ modification of angular correlation
- ◆ changes of particle composition



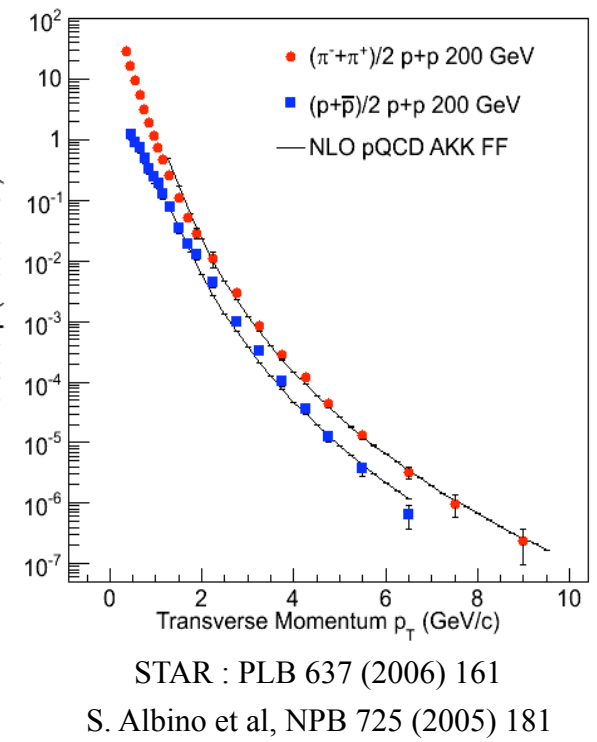
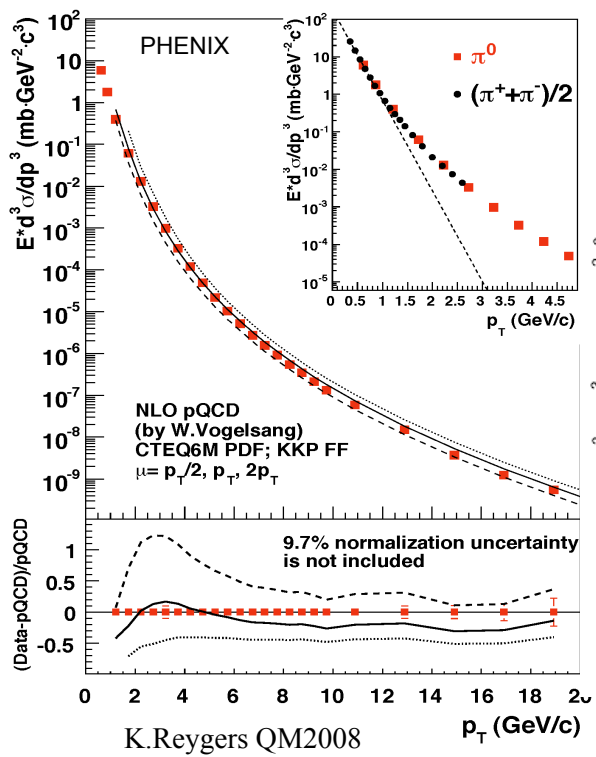
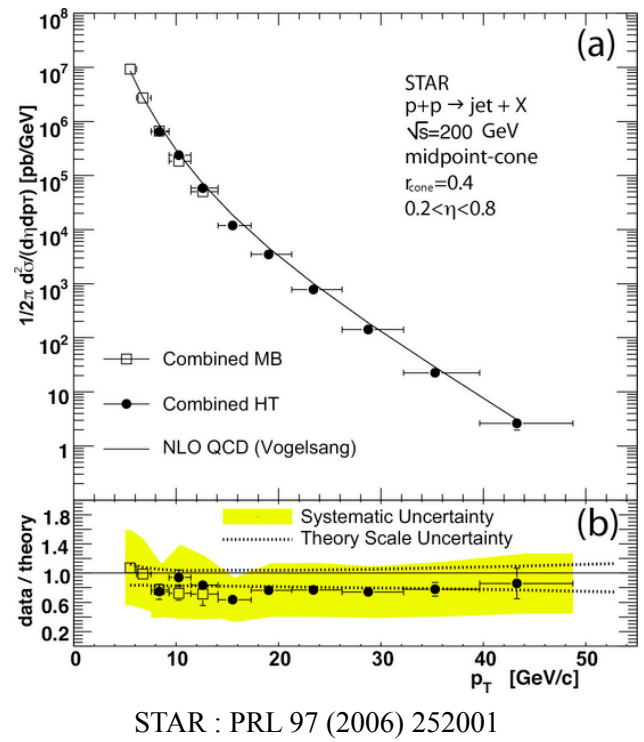
# Jets – a calibrated probe?



STAR : PRL 97 (2006) 252001

Jet production in p+p understood in pQCD framework

# Jets – a calibrated probe?

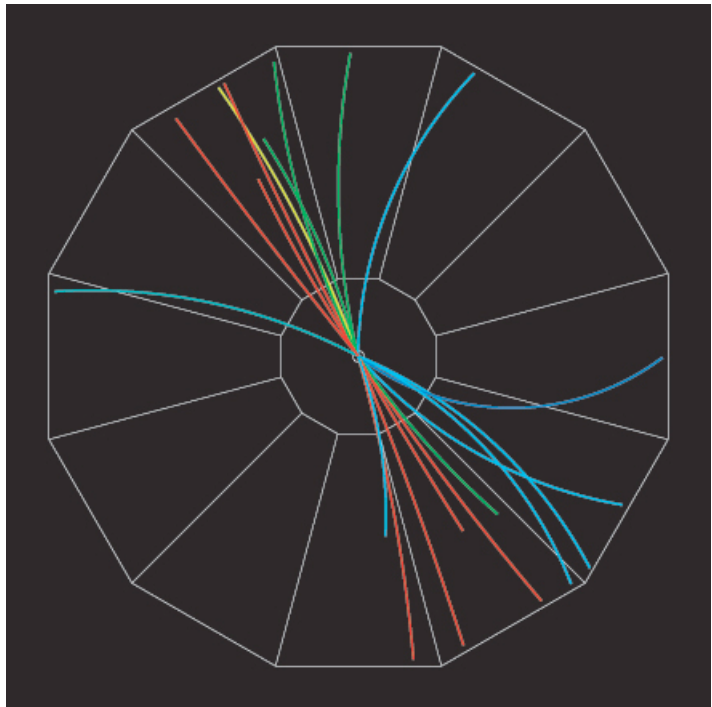


Jet production in p+p understood in pQCD framework  
Particle production in p+p also well modeled.

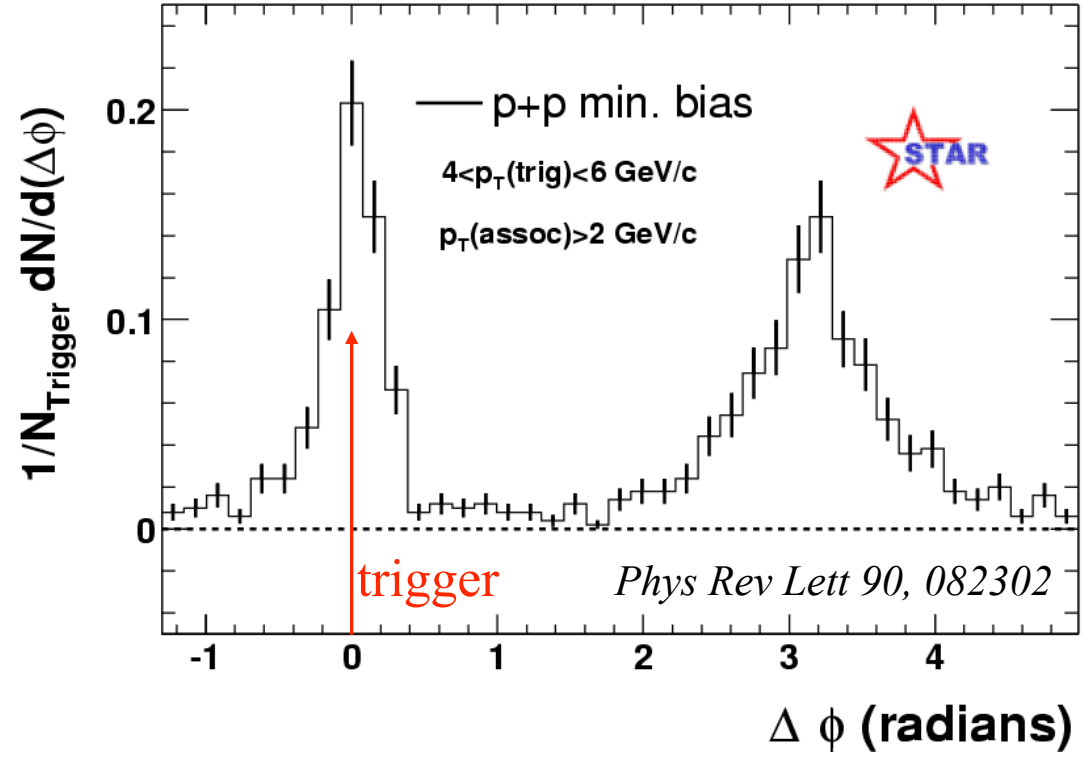
Seems we have a reasonably calibrated probe

# Jets in Au+Au collisions!

p+p → dijet



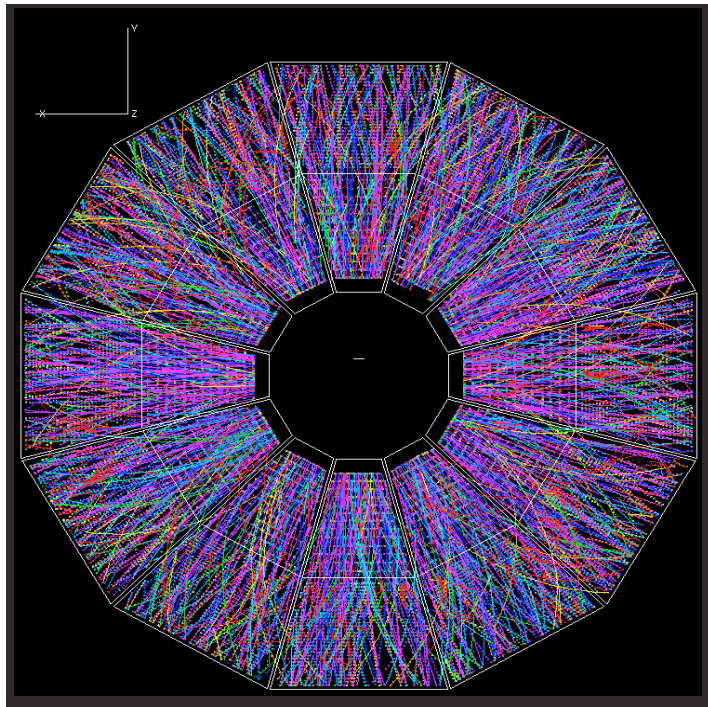
min. bias p+p collisions



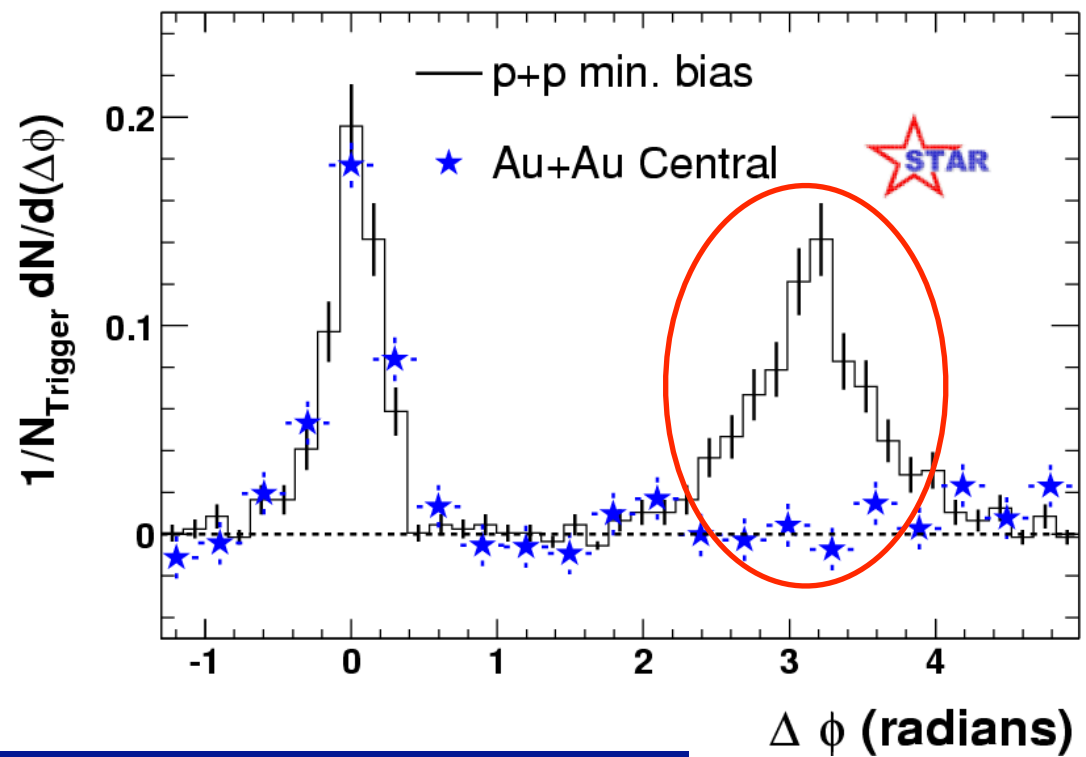
- Trigger: highest p<sub>T</sub> track
- Δφ distribution:

# Jets in Au+Au collisions!

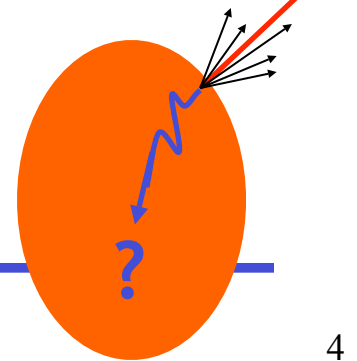
p+p → dijet



central Au+Au collisions

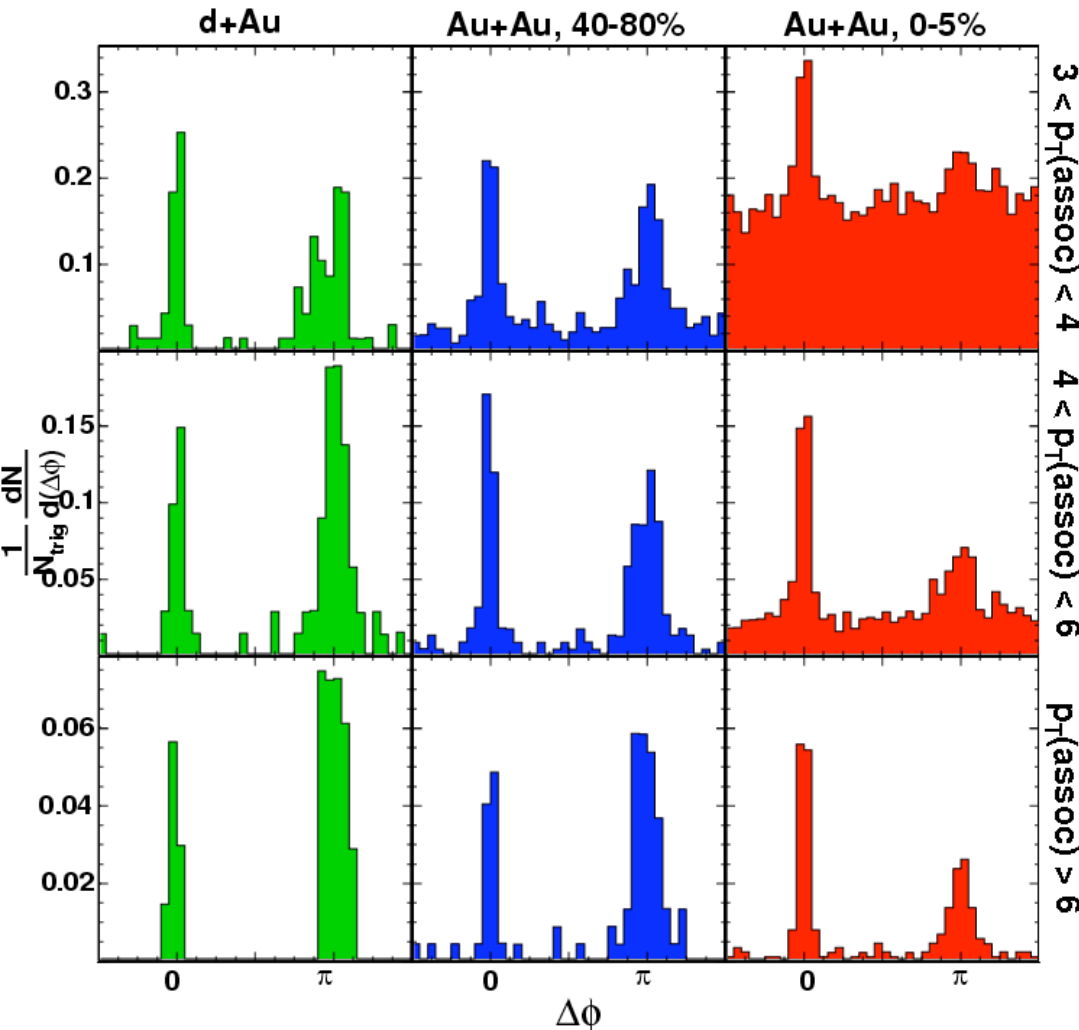


$\Delta\phi \approx 0$ : central Au+Au similar to p+p  
 $\Delta\phi \approx \pi$ : strong suppression of back-to-back correlations in central Au+Au



# Observation of “Punch through”

$8 < p_T^{\text{trig}} < 15 \text{ GeV}/c$



STAR PRL 97 (2006) 162301

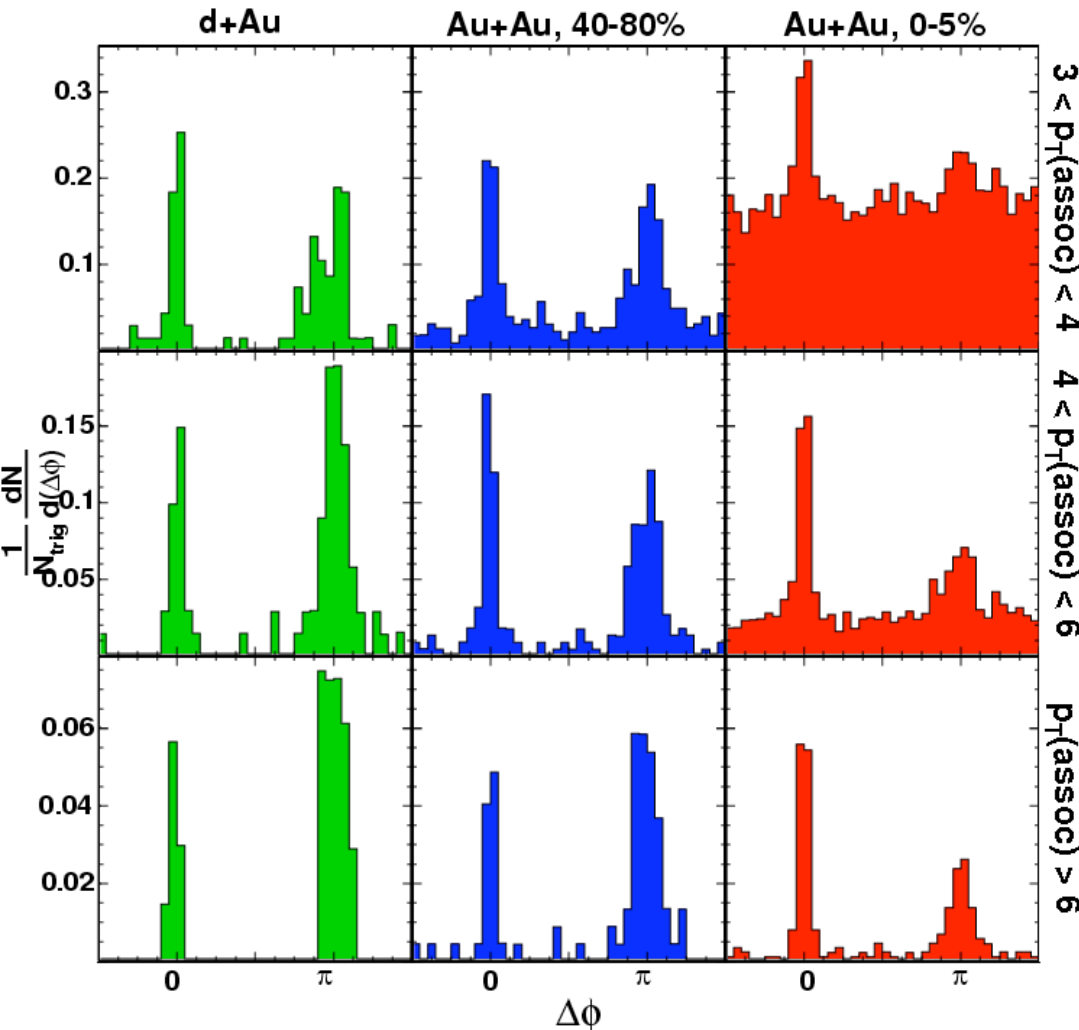
If use high- $p_T$  triggers:

- Away-side peak re-emerges
- Smaller in Au-Au than d-Au
- Virtually no background



# Observation of “Punch through”

$8 < p_T^{\text{trig}} < 15 \text{ GeV}/c$



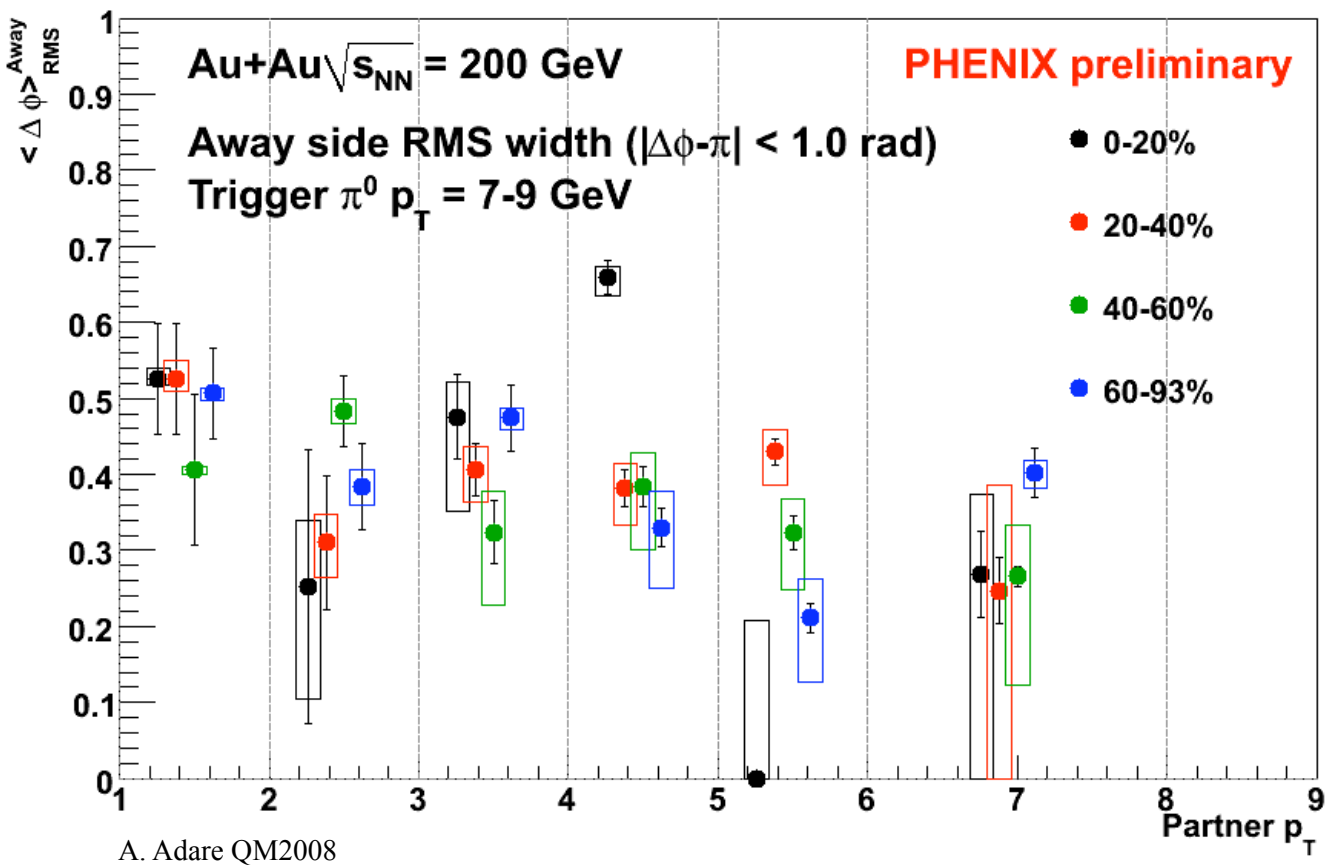
STAR PRL 97 (2006) 162301

If use high- $p_T$  triggers:

- Away-side peak re-emerges
- Smaller in Au-Au than d-Au
- Virtually no background

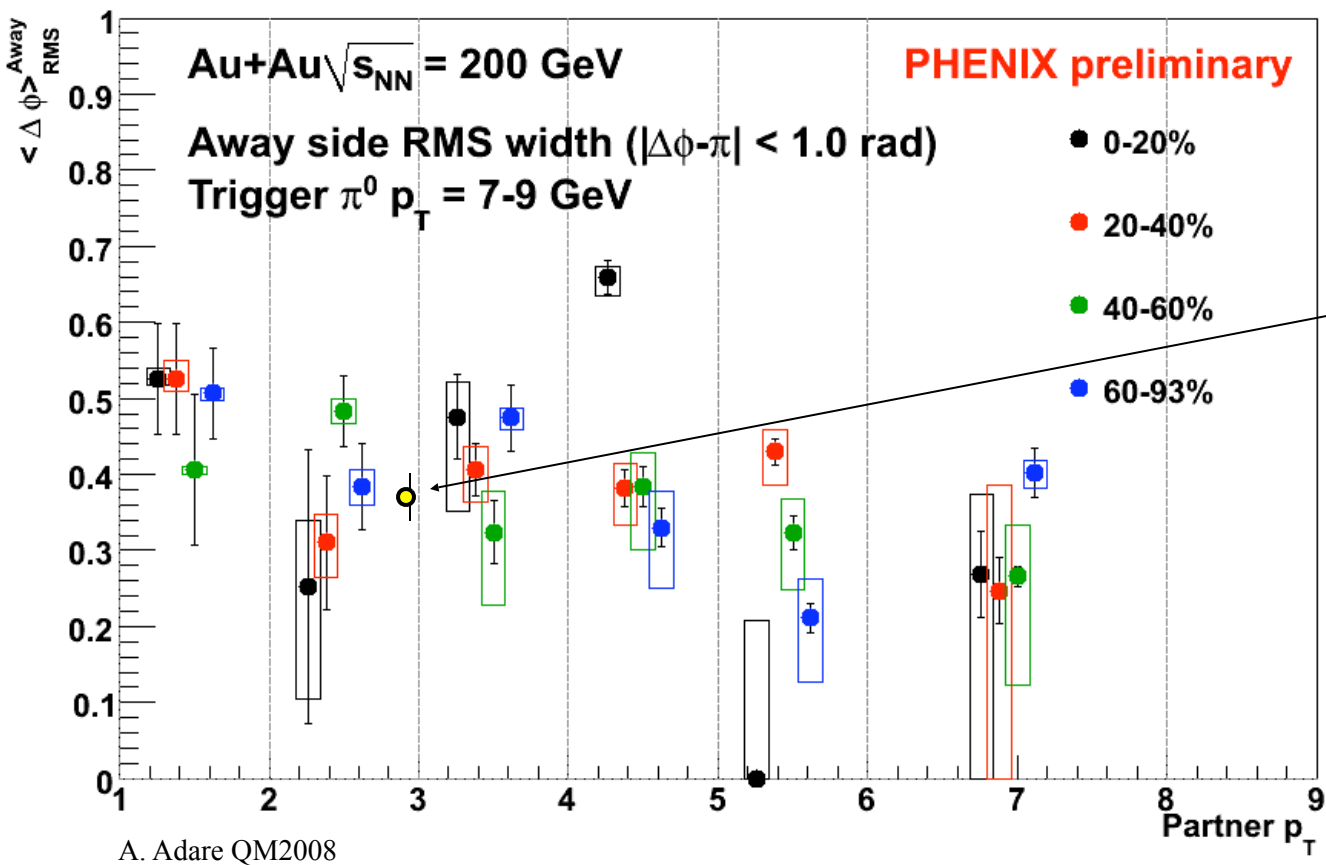
High energy jets  
“punch through”  
the medium.

# High $p_T$ triggered away side RMS width



RMS Width - centrality independent

# High $p_T$ triggered away side RMS width

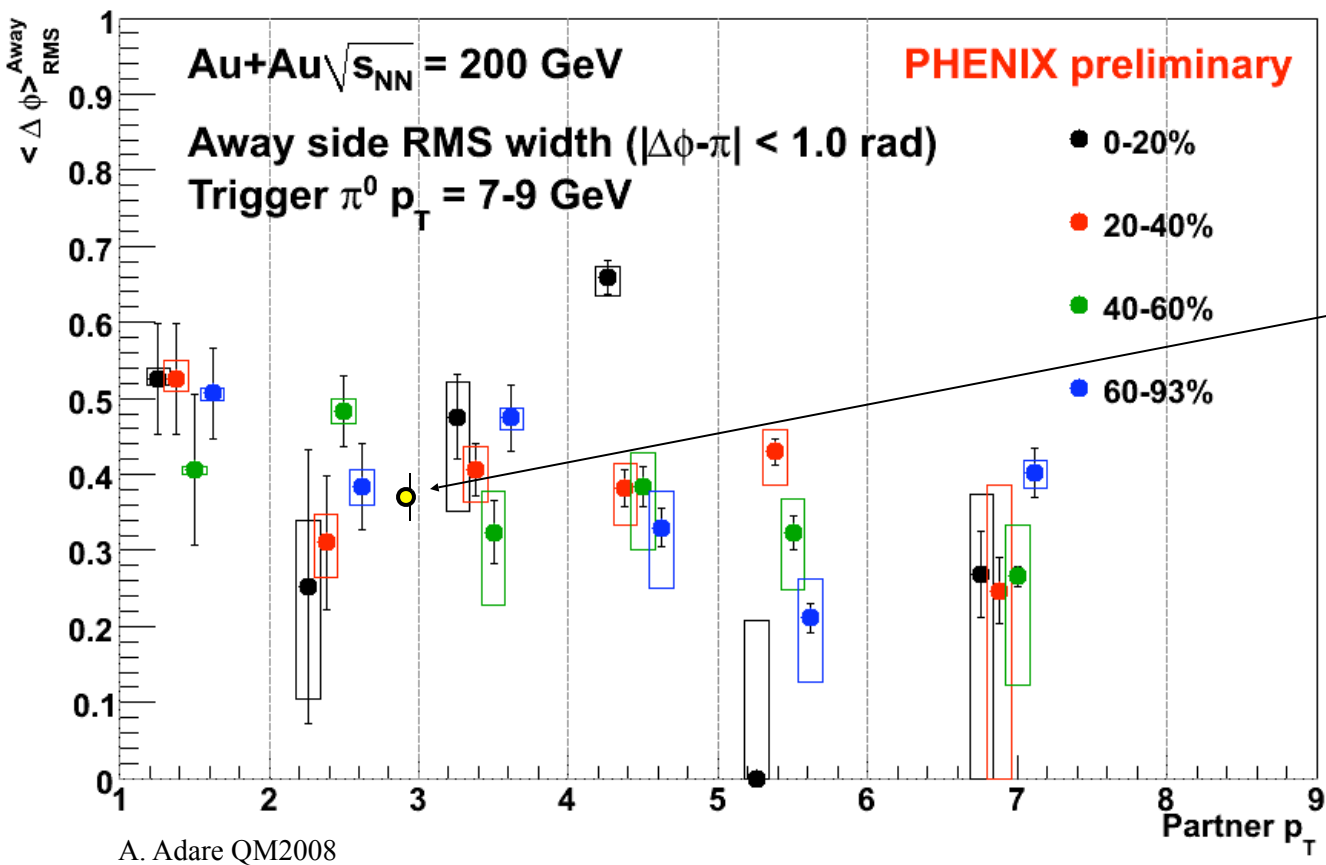


- p+p  $\pi^0$ -h:  
 $p_T^{\text{trig}} = 6.5-8$  GeV/c  
 $p_T^{\text{assoc}} = 1.4-5$  GeV/c  
**RMS =  $0.350 \pm 0.03$**   
 PHENIX:  
 Phys Rev D 74 072002

RMS Width - centrality independent

Consistent with p+p data

# High $p_T$ triggered away side RMS width



- p+p  $\pi^0$ -h:  
 $p_T^{\text{trig}} = 6.5-8$  GeV/c  
 $p_T^{\text{assoc}} = 1.4-5$  GeV/c  
**RMS =  $0.350 \pm 0.03$**   
 PHENIX:  
 Phys Rev D 74 072002

RMS Width - centrality independent

Consistent with p+p data

**Vacuum fragmentation?**

# Away-side di-hadron fragmentation functions

- Study medium-induced modification of fragmentation function due to energy loss
- Without full jet reconstruction, parton energy not measurable
  - $z$  not measured ( $z = p_{\text{hadron}}/p_{\text{parton}}$ )
  - $z_T = p_{T,\text{assoc}}/p_{T,\text{trig}}$
- Di-hadron fragmentation function - di-hadron jet-like correlated yield to single hadron yield ratio

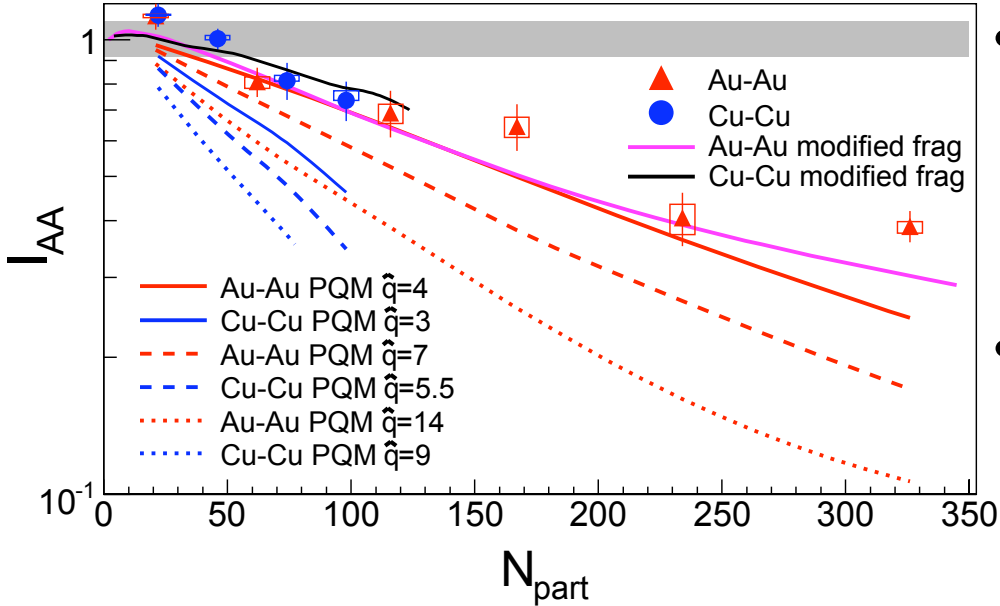
$$D^{h_1 h_2}(z_T, p_T^{\text{trig}}) = p_T^{\text{trig}} \frac{d\sigma_{AA}^{h_1 h_2} / dp_T^{\text{trig}} dp_T}{d\sigma_{AA}^{h_1} / dp_T^{\text{trig}}}$$

$$I_{AA} = \frac{D_{AA}(z_T, p_T^{\text{trig}})}{D_{pp}(z_T, p_T^{\text{trig}})}$$

# Away-side di-hadron fragmentation functions

H. Zhong et al., PRL 97 (2006) 252001  
 C. Loizides, Eur. Phys. J. C 49, 339-345 (2007)

$6 < p_{T \text{ trig}} < 10 \text{ GeV}$



- Study medium-induced modification of fragmentation function due to energy loss
- Without full jet reconstruction, parton energy not measurable
- $z$  not measured ( $z = p_{\text{hadron}}/p_{\text{parton}}$ )
- $z_T = p_{T, \text{assoc}}/p_{T, \text{trig}}$
- Di-hadron fragmentation function - di-hadron jet-like correlated yield to single hadron yield ratio

$$D^{h_1 h_2}(z_T, p_T^{\text{trig}}) = p_T^{\text{trig}} \frac{d\sigma_{AA}^{h_1 h_2} / dp_T^{\text{trig}} dp_T}{d\sigma_{AA}^{h_1} / dp_T^{\text{trig}}}$$

$$I_{AA} = \frac{D_{AA}(z_T, p_T^{\text{trig}})}{D_{pp}(z_T, p_T^{\text{trig}})}$$

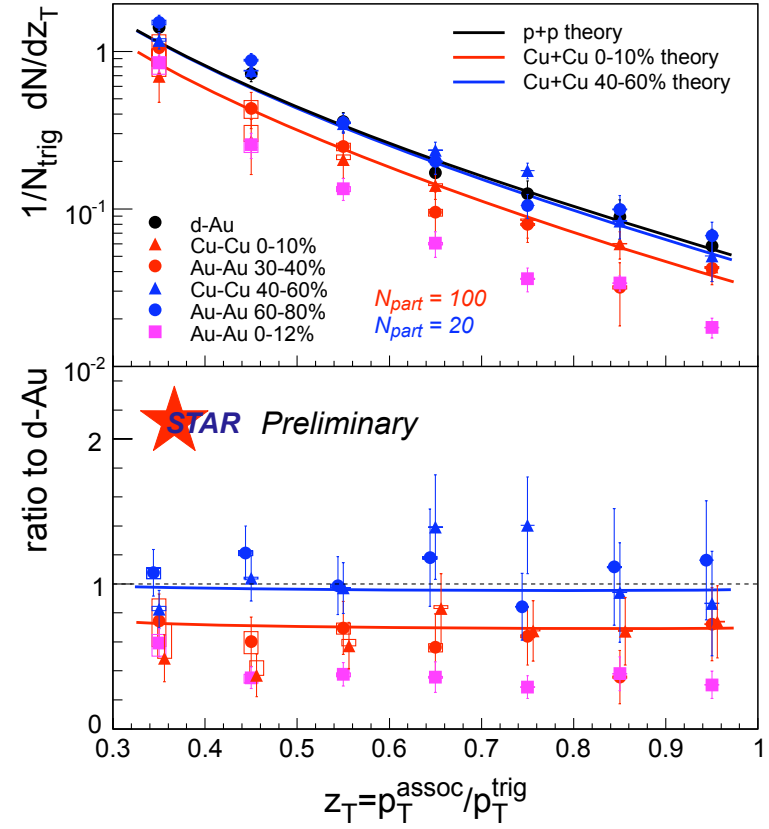
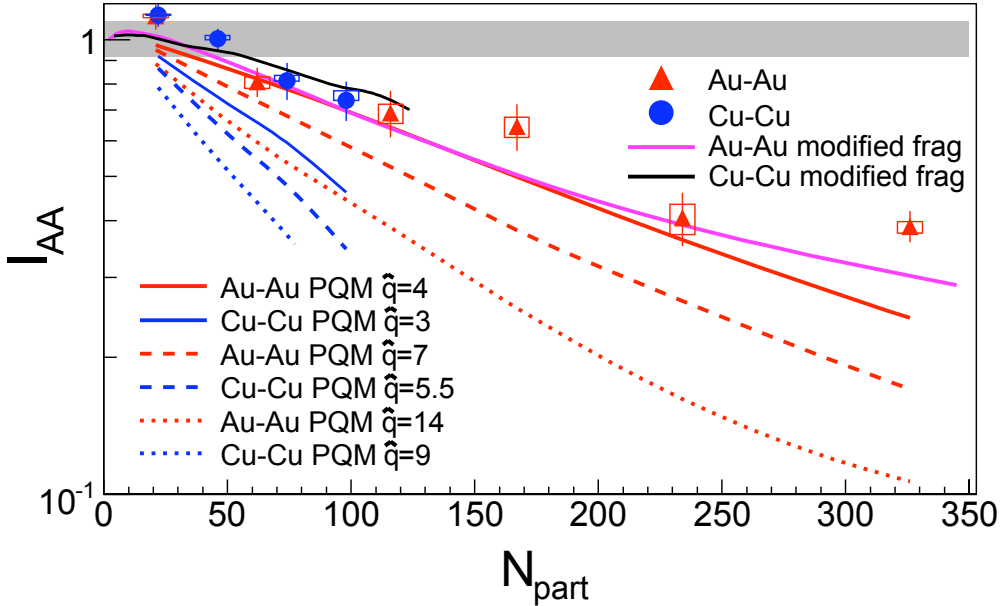
- Inconsistent with *Parton Quenching Model* calculation
- *Modified fragmentation model* better

O. Catu QM2008

# Away-side di-hadron fragmentation functions

H. Zhong et al., PRL 97 (2006) 252001  
 C. Loizides, Eur. Phys. J. C 49, 339-345 (2007)

$6 < p_{T \text{ trig}} < 10 \text{ GeV}$



- Inconsistent with *Parton Quenching Model* calculation
- *Modified fragmentation model* better

- Denser medium in central Au+Au than central Cu+Cu
- Similar medium for similar  $N_{part}$

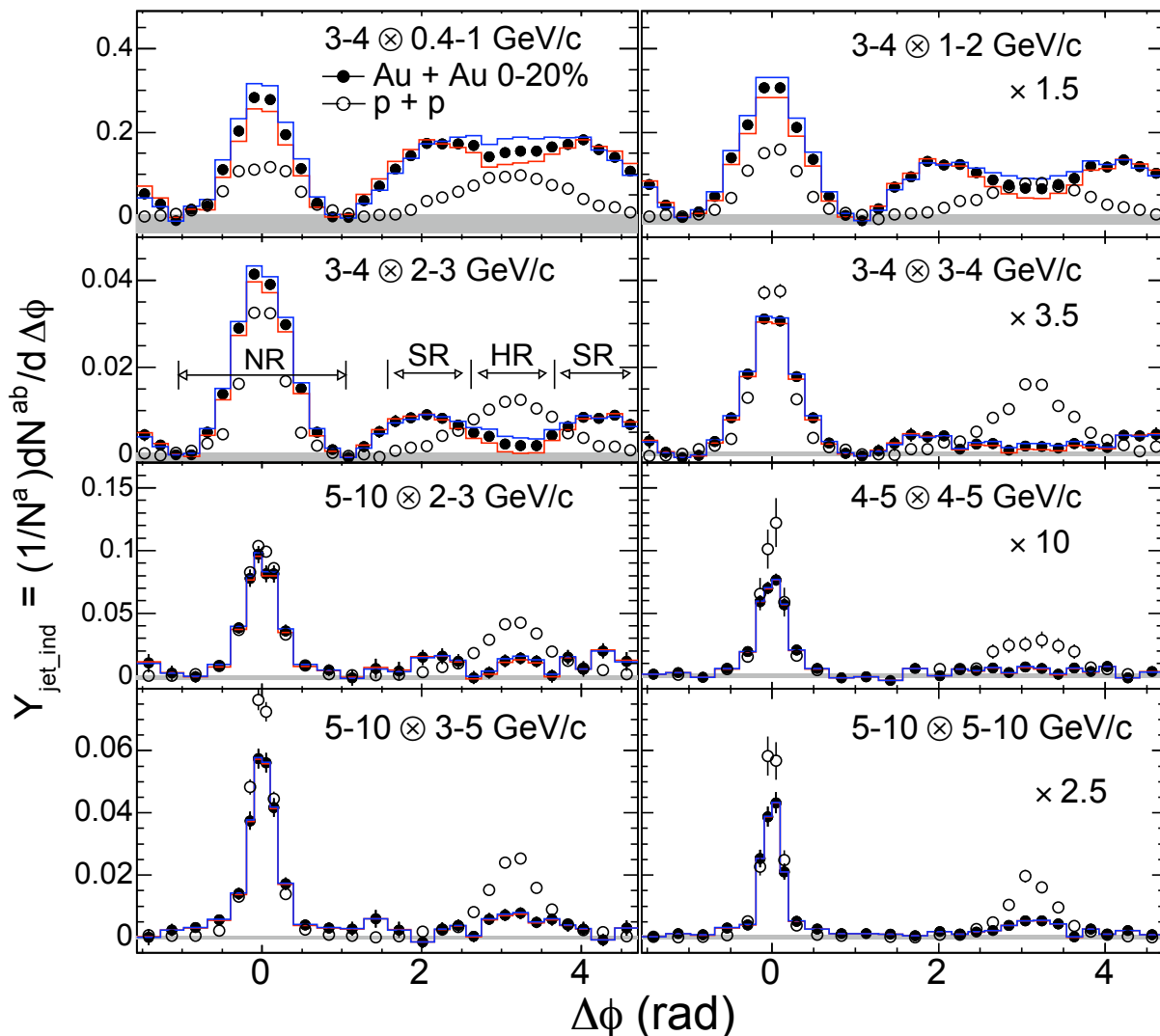
Vacuum fragmentation after parton  $E_{loss}$  in the medium

O. Catu QM2008

# $p_T$ evolution of di-hadron correlations

PHENIX Preliminary arXiv:0801.4545 [nucl-ex]

- As  $p_T$  decreases,
  - single-peak  $\rightarrow$  double peak
  - Away-side yield increases

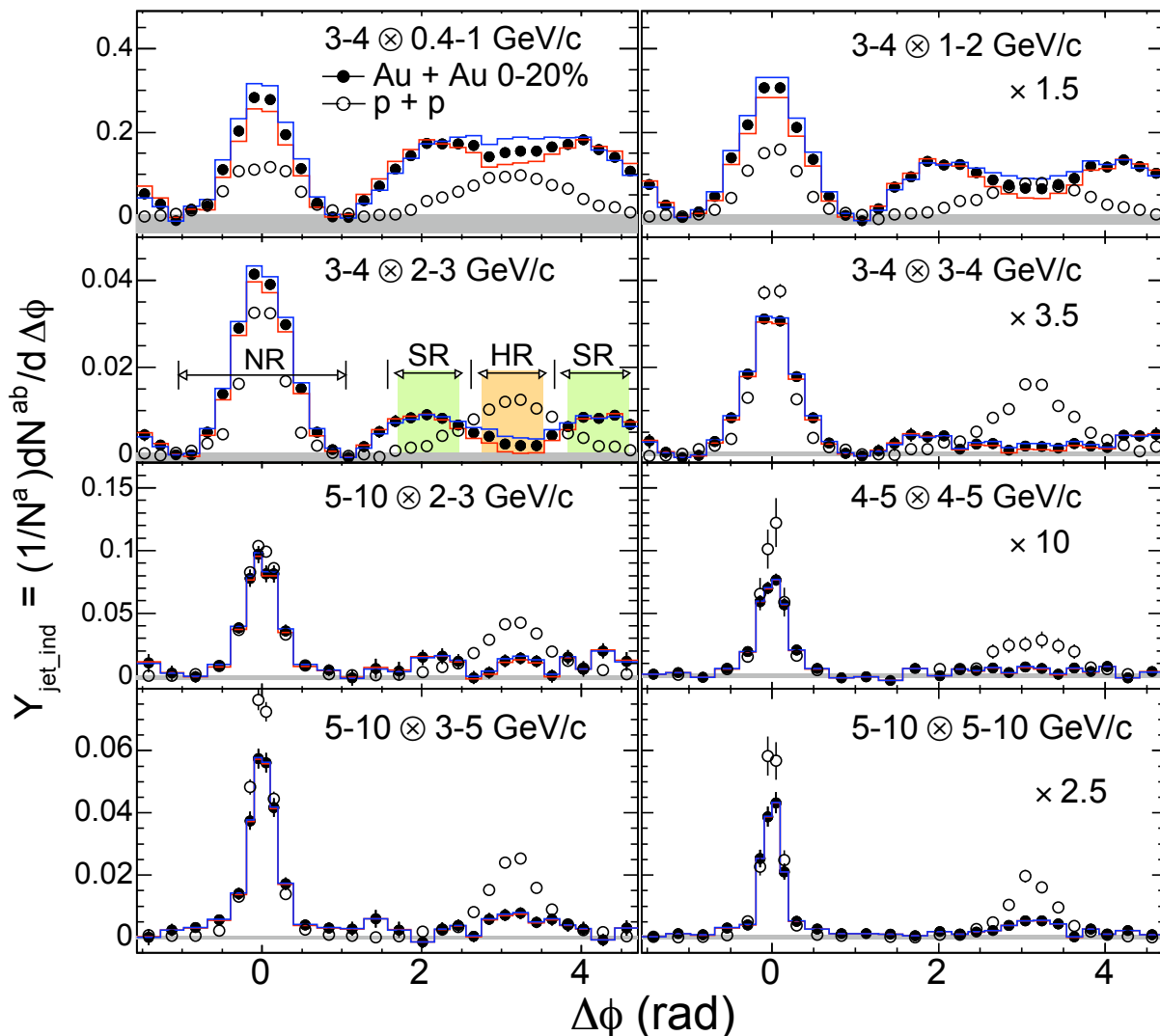




# $p_T$ evolution of di-hadron correlations

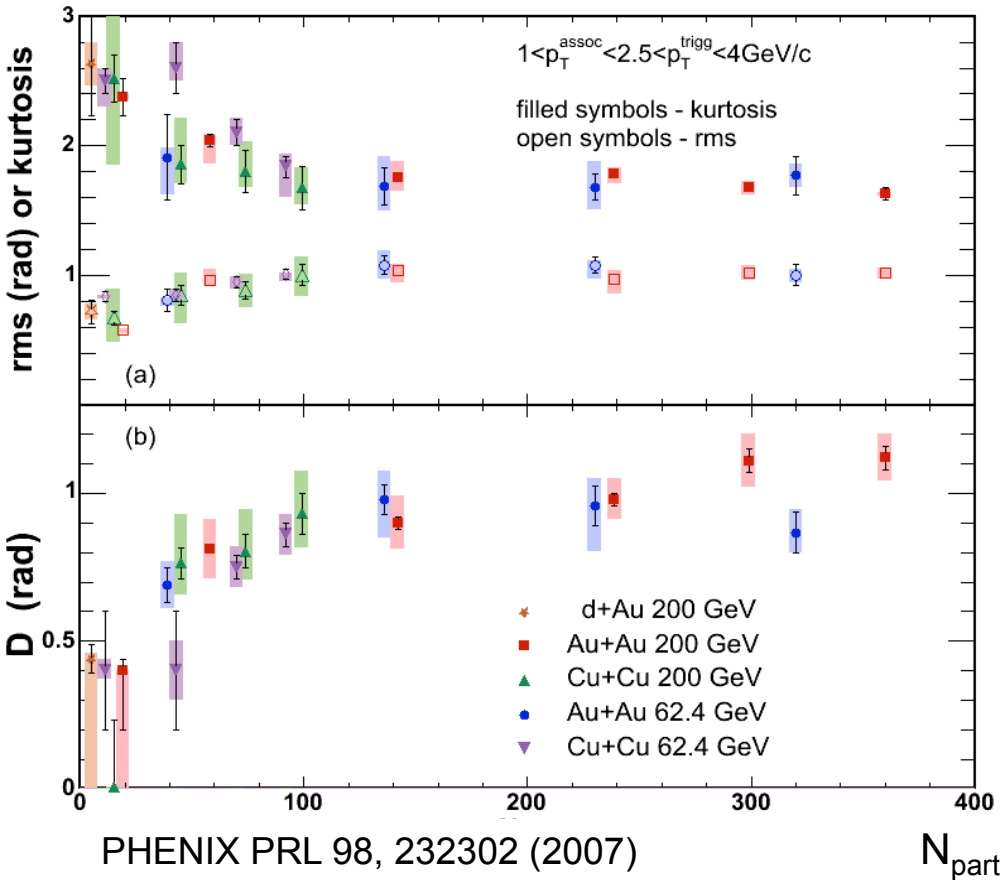
PHENIX Preliminary arXiv:0801.4545 [nucl-ex]

- As  $p_T$  decreases,
  - single-peak  $\rightarrow$  double peak
  - Away-side yield increases
- **Head region** yield begins to dominate over **shoulder region**



# Head and Shoulder evolution

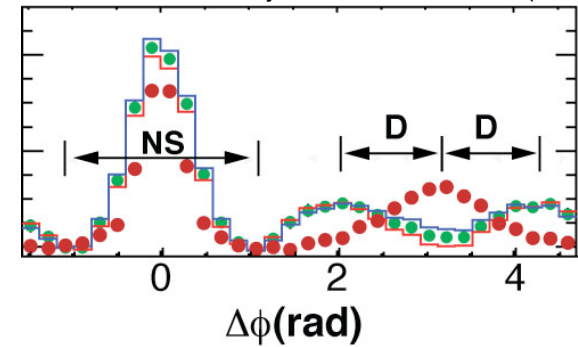
h-h correlations  $1 < p_{T_a} < 2.5 < p_{T_t} < 4 \text{ GeV}/c$



$$J(\Delta\phi) =$$

$$G(\Delta\phi) + G(\Delta\phi - \pi - D) + G(\Delta\phi - \pi + D)$$

PHENIX Phys.Rev. C 77, 011901 (2008)



$$\mu_n \equiv (\Delta\phi - \pi)^n, n = 2, 4, \dots$$

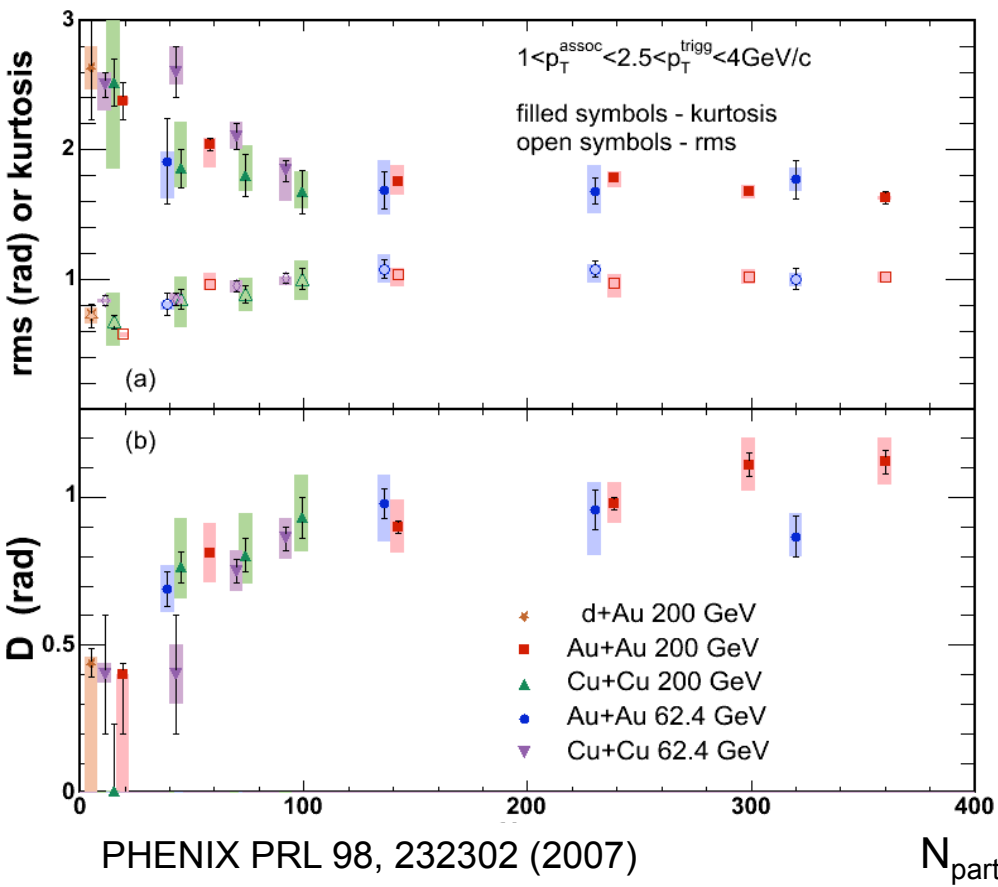
$$rms \equiv \sqrt{\mu_2}$$

$$kurtosis \equiv \frac{\mu_4}{\mu_2^2}$$

Separation and width plateau for  $N_{\text{part}} \geq 100$

# Head and Shoulder evolution

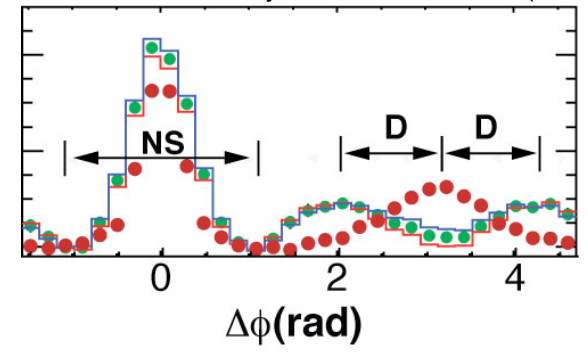
h-h correlations  $1 < p_{T_a} < 2.5 < p_{T_t} < 4 \text{ GeV}/c$



$$J(\Delta\phi) =$$

$$G(\Delta\phi) + G(\Delta\phi - \pi - D) + G(\Delta\phi - \pi + D)$$

PHENIX Phys.Rev. C 77, 011901 (2008)



$$\mu_n \equiv (\Delta\phi - \pi)^n, n = 2, 4, \dots$$

$$rms \equiv \sqrt{\mu_2}$$

$$kurtosis \equiv \frac{\mu_4}{\mu_2^2}$$

Separation and width plateau for  $N_{\text{part}} \geq 100$

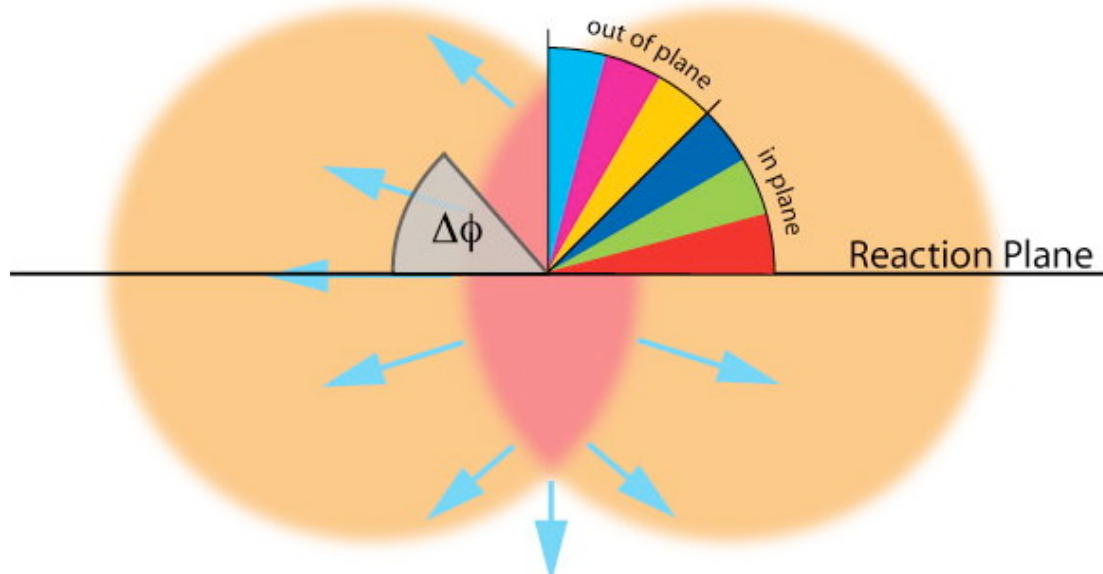
Shoulder due to away side jet interacting with medium

# Path length dependencies

Non-central events have “elliptic” overlap geometries

Measurements w.r.t reaction plane angle:

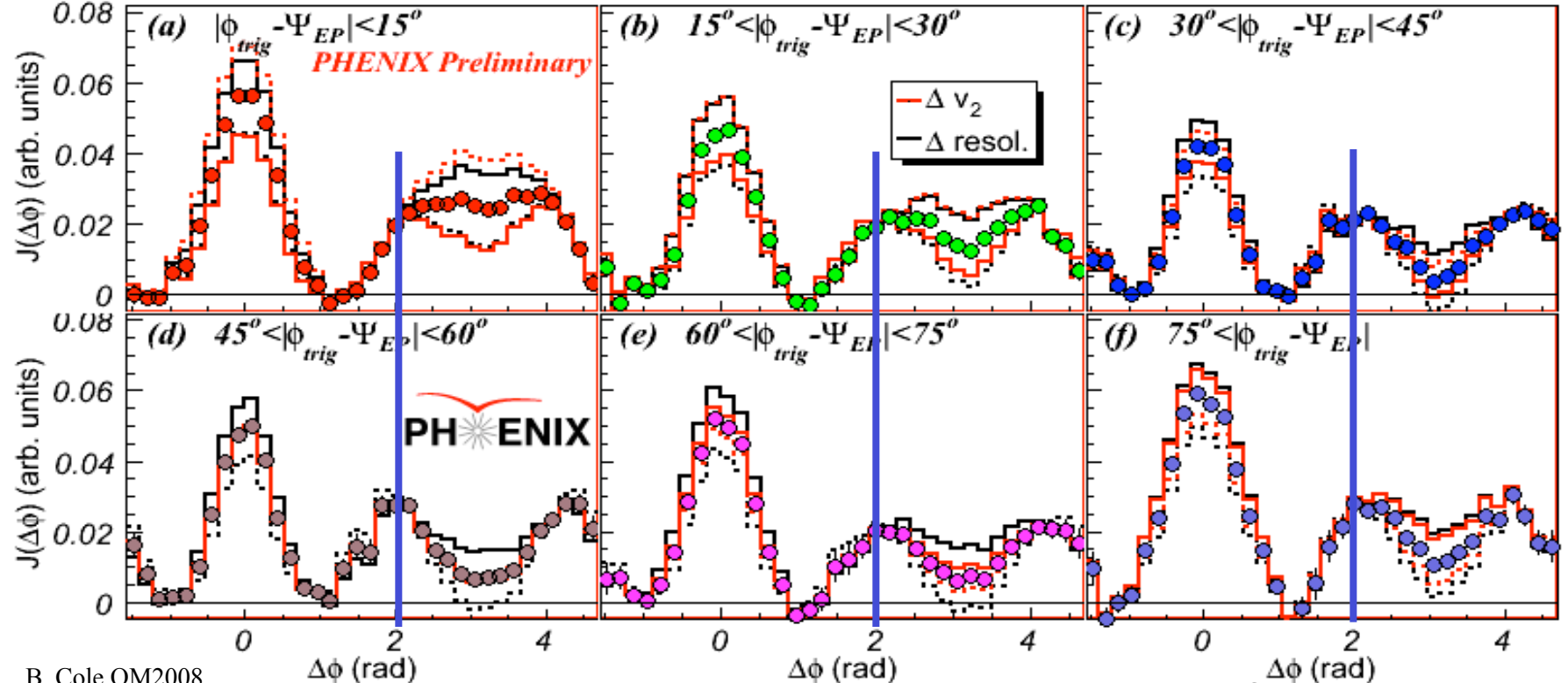
- Change path length
- Keep everything else same



Isolate effects due to path length

# Path length effect on di-hadron correlation

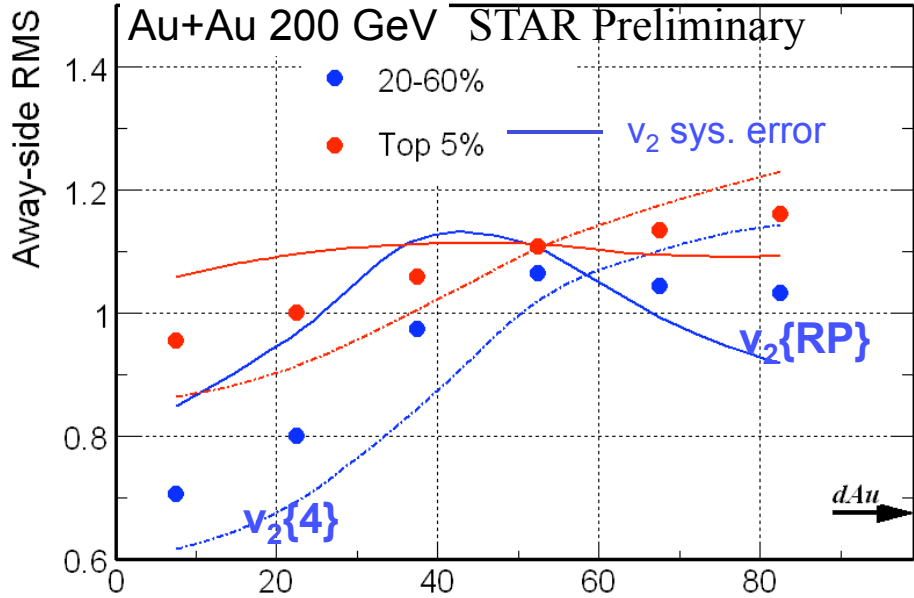
Au+Au  $\sqrt{s_{NN}}=200\text{GeV}$ , Cent=30-40%,  $1 < p_{T,assoc} < 2 \text{ GeV}/c$ ,  $2 < p_{T,trig} < 3 \text{ GeV}/c$   
 in-plane  $\phi_t - \Psi_{RP} = 0$



B. Cole QM2008

- Near side peak unchanged
- Shoulder peaks emerge as  $\phi_t - \Psi$  increases but are at fixed  $\Delta\phi$
- Head peak (di-jet remnant) decreases as  $\phi_t - \Psi_{RP}$  increases

# Centrality and path length effects



A. Feng QM2008

$$\phi_s = \phi_T - \Psi_{RP} \text{ (deg)}$$

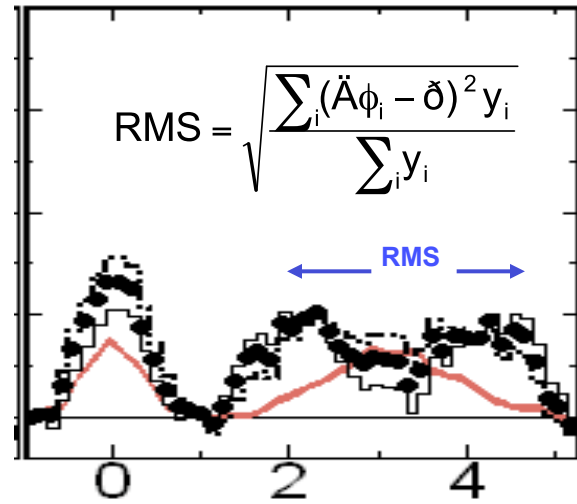
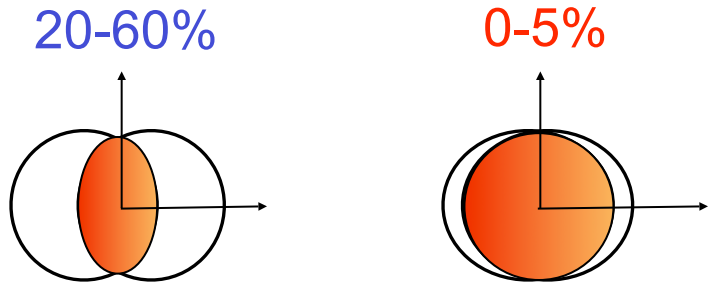
$3 < p_{T}^{trig} < 4 \text{ GeV}/c, 1.0 < p_{T}^{asso} < 1.5 \text{ GeV}/c$

In-plane: 20-60% ~ d+Au

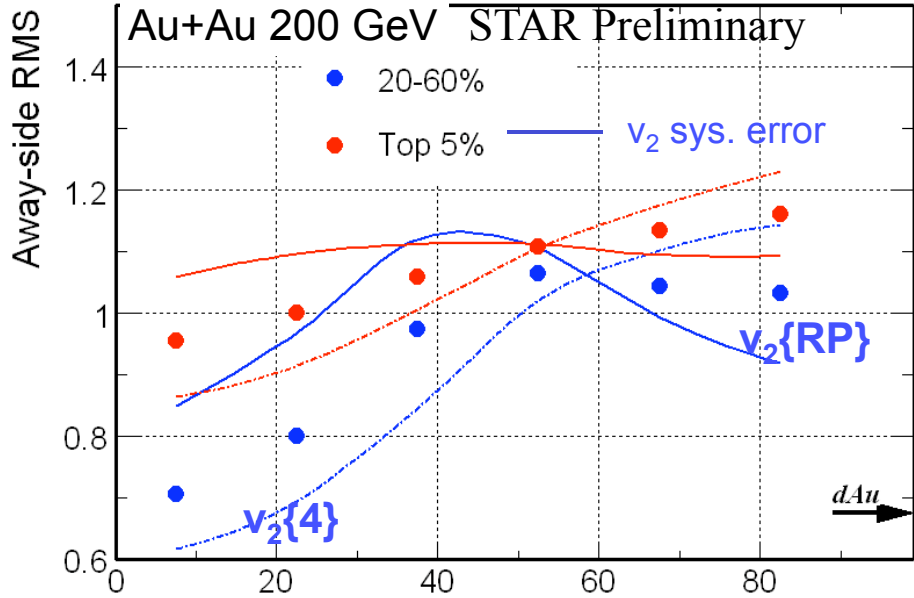
0-5% > d+Au

Out-of-plane: 20-60% ~ 0-5%

Au+Au > d+Au



# Centrality and path length effects



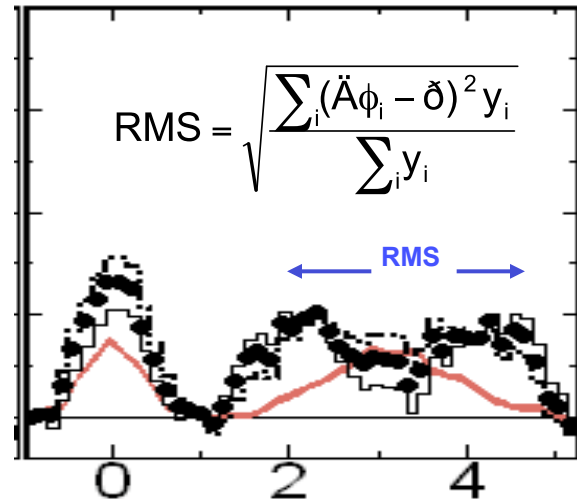
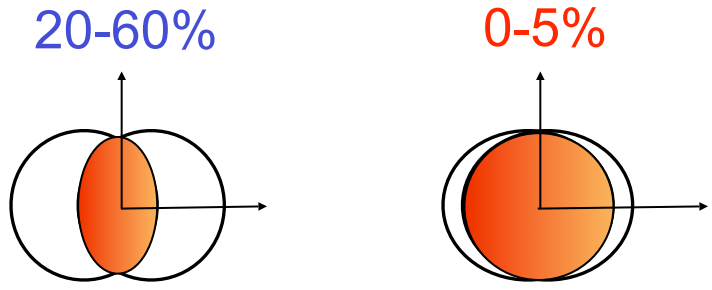
A. Feng QM2008  
 $\phi_s = \phi_T - \Psi_{RP}$  (deg)  
 $3 < p_{T}^{trig} < 4$  GeV/c,  $1.0 < p_{T}^{asso} < 1.5$  GeV/c

In-plane: 20-60% ~ d+Au

0-5% > d+Au

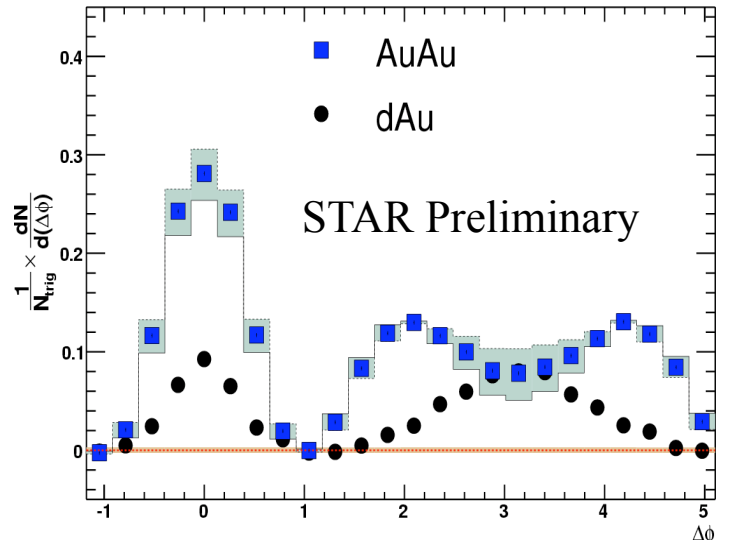
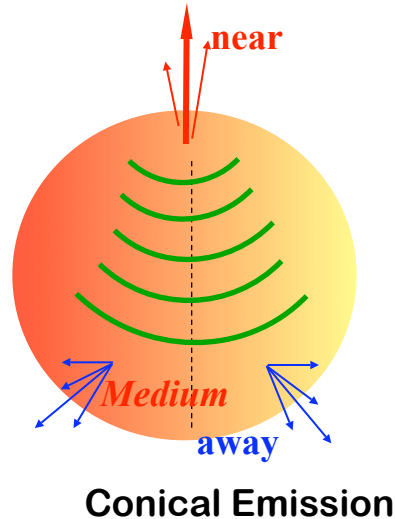
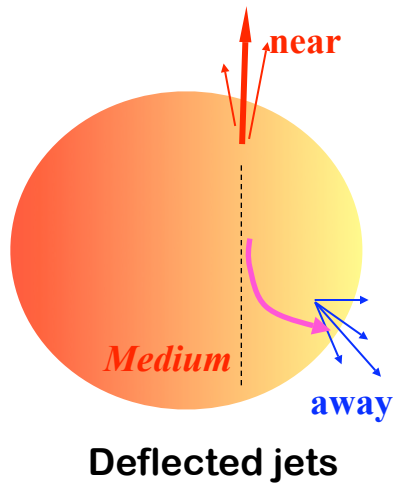
Out-of-plane: 20-60% ~ 0-5%

Au+Au > d+Au



Away-side features reveal path length effects

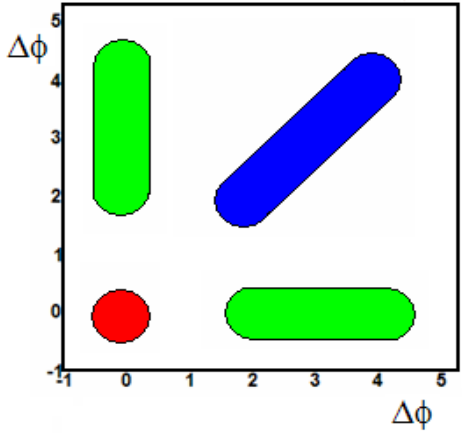
# What causes shoulder peaks?



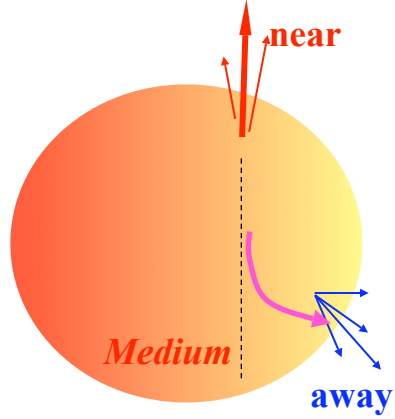
Conical emission or deflected jets?



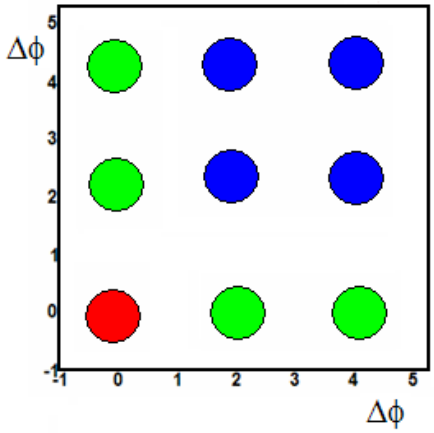
# What causes shoulder peaks?



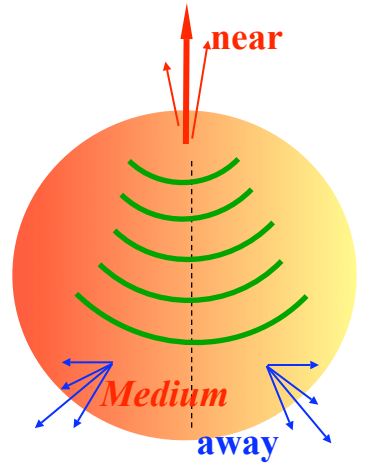
Deflected jets



Deflected jets



Conical Emission

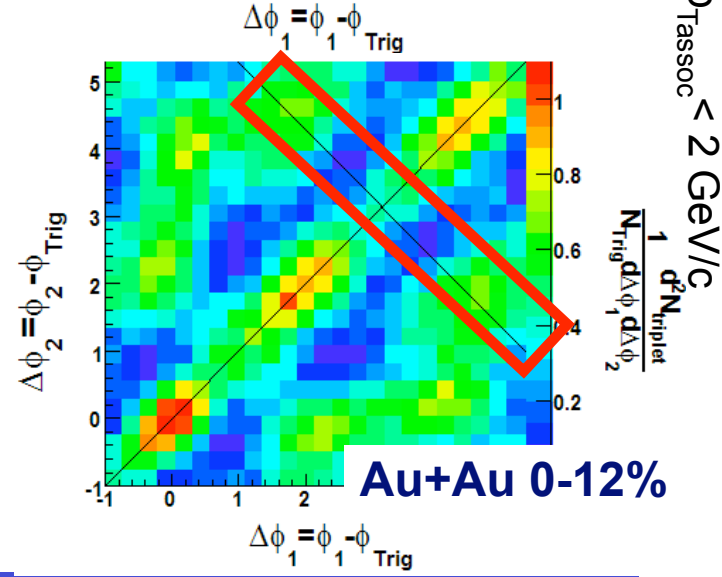
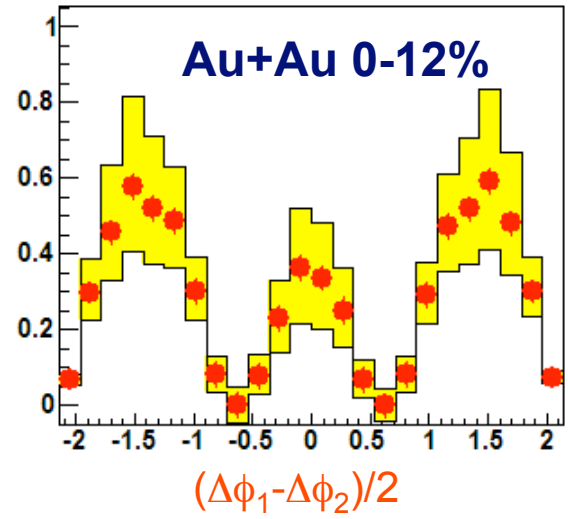
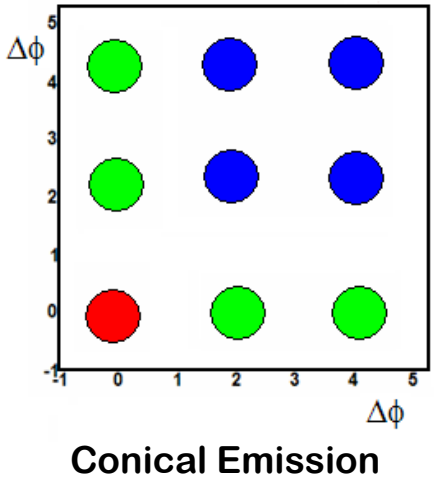
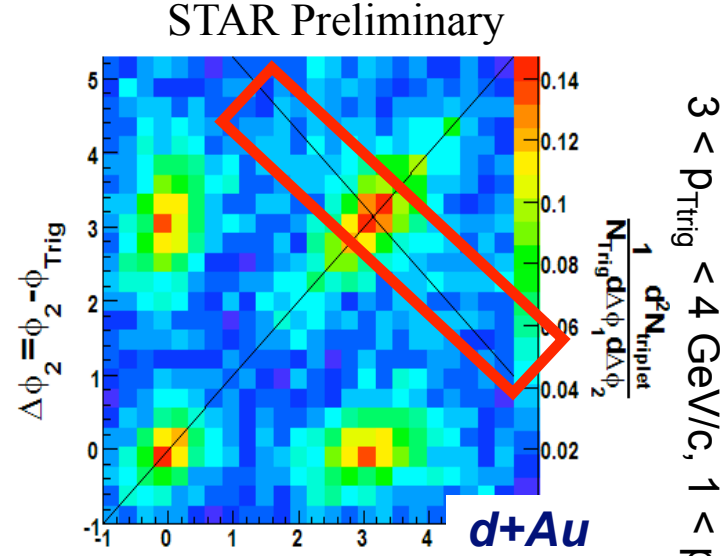
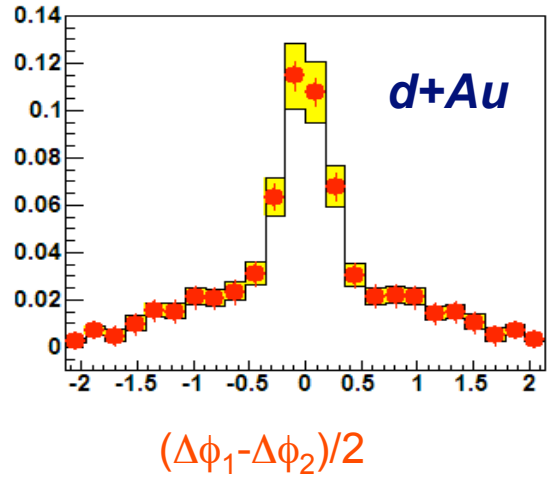
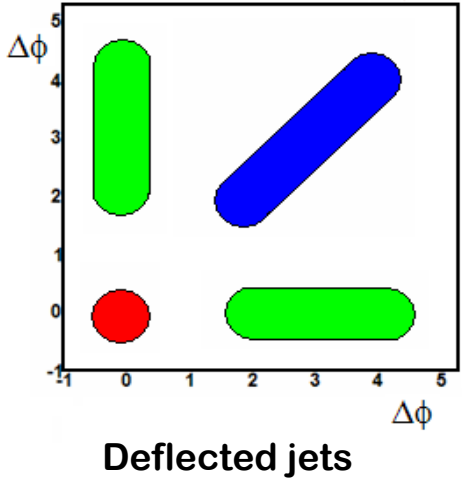


Conical Emission

- Two component approach:
- Correlated to trigger (jets..)
  - Uncorrelated to trigger (except via anisotropic flow)
- Bkg normalization 3-particle ZYAM

B. Mohnaty QM2008

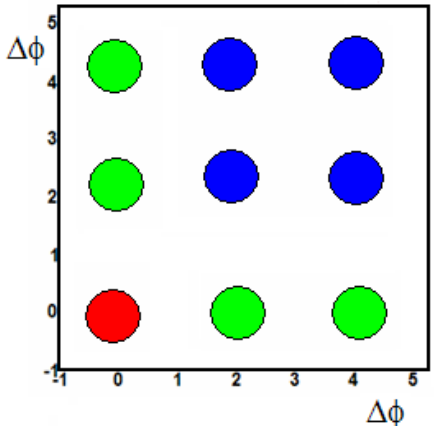
# What causes shoulder peaks?



B. Mohnaty QM2008

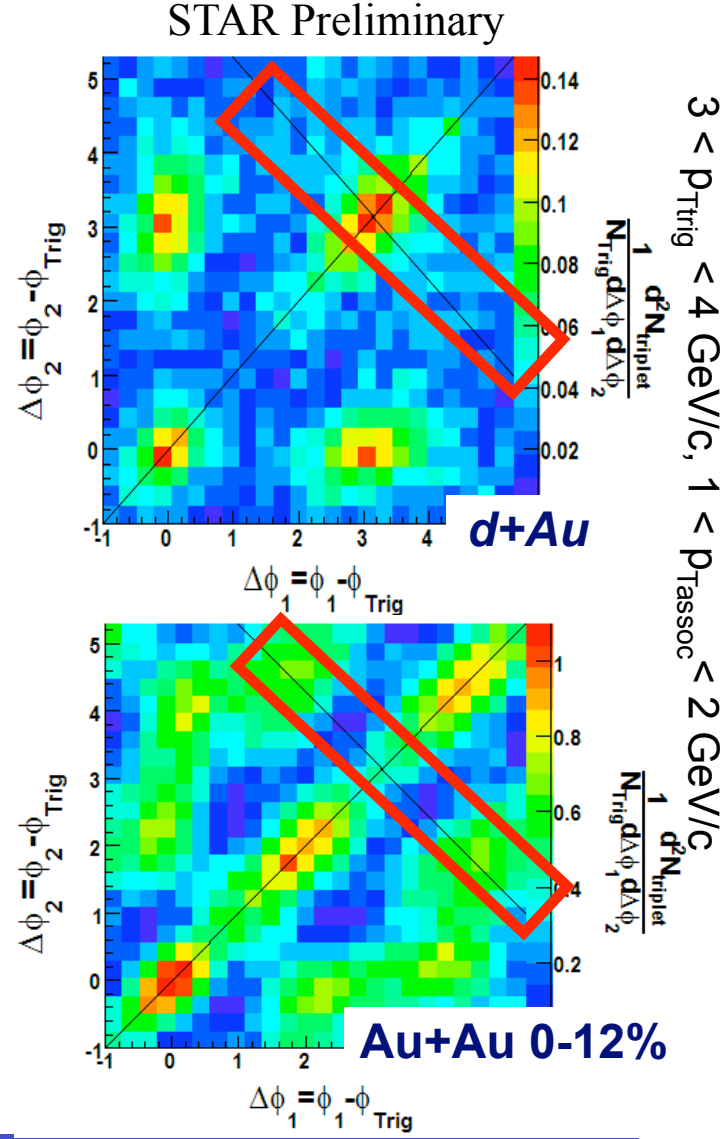
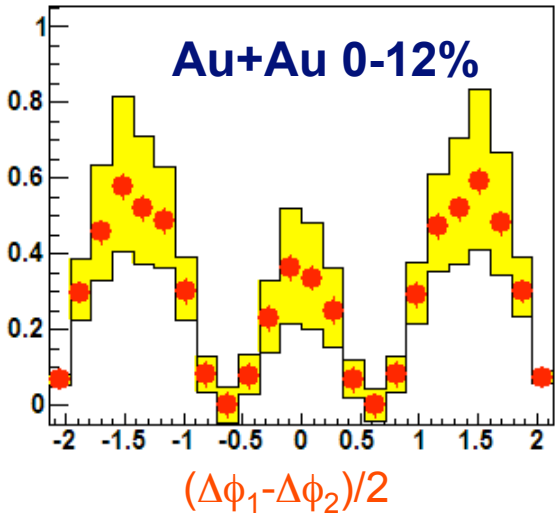
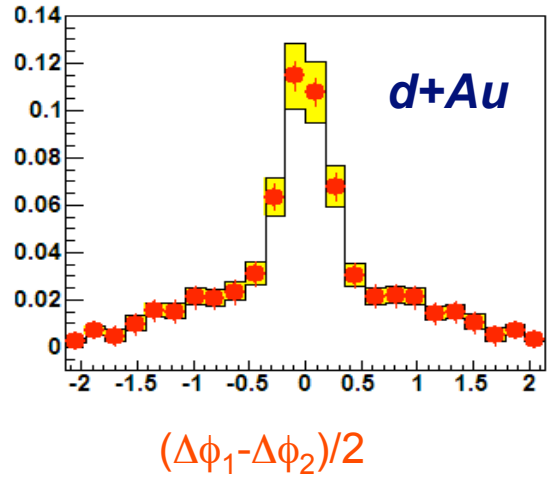
# What causes shoulder peaks?

Au+Au data consistent with Conical emission



Conical Emission

B. Mohnaty QM2008



$3 < p_{Trig} < 4 \text{ GeV}/c, 1 < p_{Tassoc} < 2 \text{ GeV}/c$

# Possible causes of conical emission

## Mach Cone

Similar to jet creating sonic boom in air.



- Energy radiated from parton deposited in collective hydrodynamic modes.

- Mach angle depends on  $C_s$ 
  - T dependent

$$\frac{c_s}{v_{parton}} = \cos(\theta_M)$$

- Angle independent of  $p_T^{assoc}$

# Possible causes of conical emission

## Mach Cone

Similar to jet creating sonic boom in air.



- Energy radiated from parton deposited in collective hydrodynamic modes.

- Mach angle depends on  $C_s$

- T dependent

$$\frac{C_s}{v_{parton}} = \cos(\theta_M)$$

- Angle independent of  $p_T^{assoc}$

## Čerenkov Gluon Radiation

- Gluons radiated by superluminal parton.

$$\begin{aligned} \frac{C_n}{v_{parton}} &= \cos(\theta_c) \\ &= \frac{c}{n(p)v_{parton}} \\ &\approx \frac{1}{n(p)} \end{aligned}$$

- Angle dependent on  $p_T^{assoc}$

# Conical Emission Theories

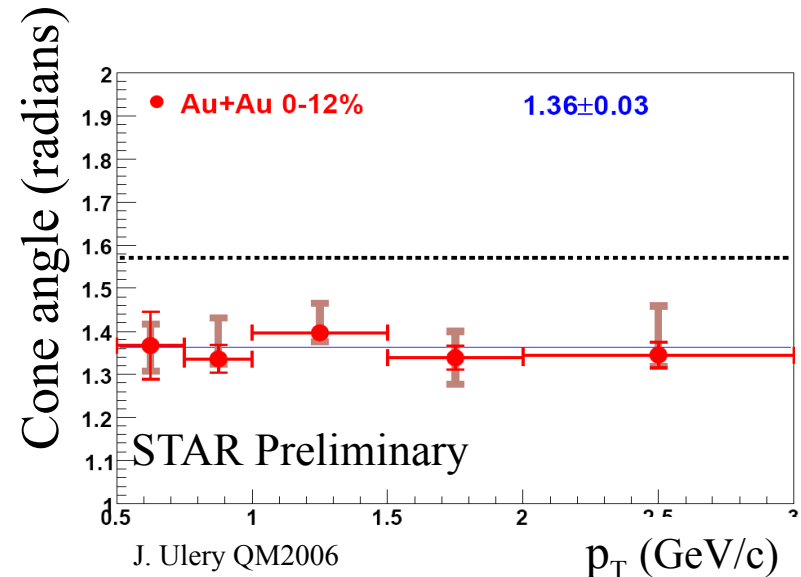
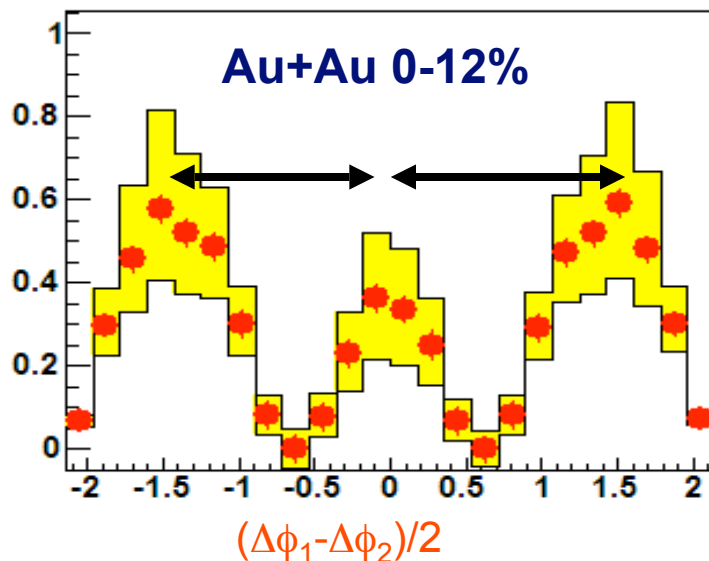
- **Mach-cone:**
  - Can be produced in different theories:
    - **Hydrodynamics**
      - H. Stöcker et al. (Nucl.Phys.A750:121,2005)
      - J. Casalderra-Solana et. al. (Nucl.Phys.A774:577,2006)
      - T. Renk & J. Ruppert (Phys.Rev.C73:011901,(2006))
    - **Colored plasma**
      - J. Ruppert & B. Müller (Phys.Lett.B618:123,2005)
    - **AdS/CFT**
      - S. Gubser, S. Pufu, A. Yarom. (arXiv:0706.4307v1, 2007)
- **Čerenkov Gluon Radiation:**
  - I.M. Dremin (Nucl. Phys. A750: 233, 2006)
  - V. Koch et. al. (Phys. ReV. Lett. 96, 172302, 2006)
- **Parton Cascade:**
  - G. L. Ma et. al. (Phys. Lett. B647, 122, 2007)

References are only a small subset of those existing. Apologies to those not included.

# Mach cone or Čerenkov gluons?

Angle predictions:

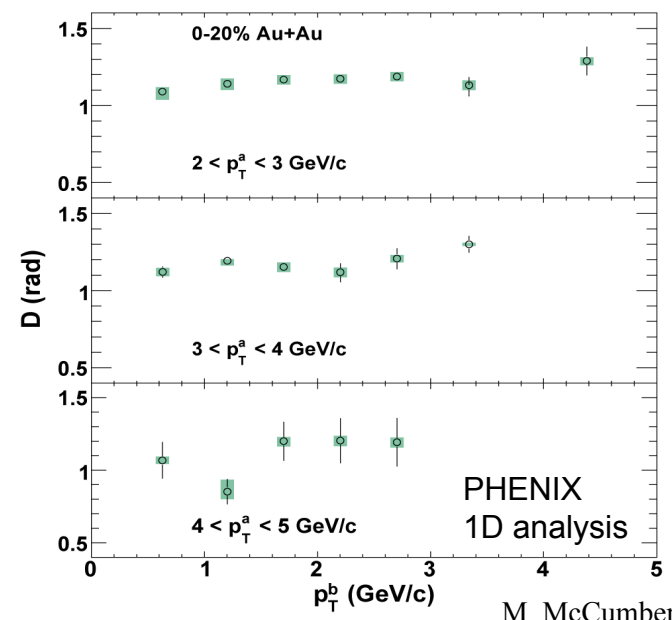
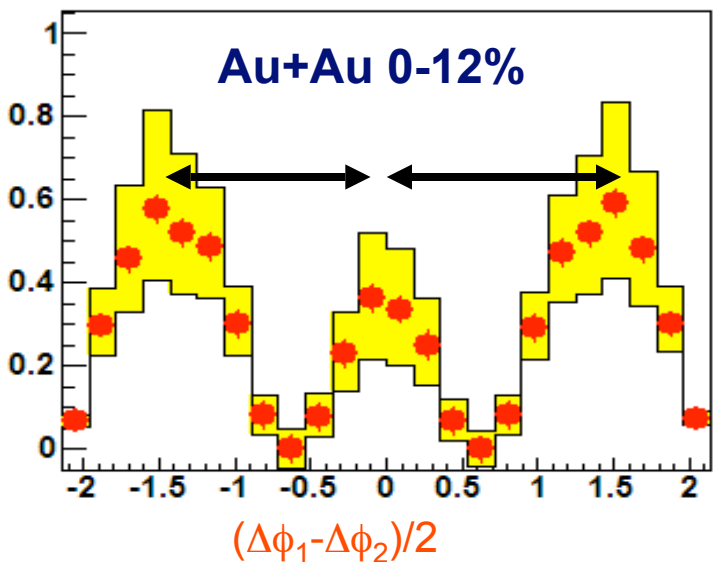
- Mach-cone:  
Angle **independent** of associated  $p_T$
- Čerenkov gluon radiation:  
Angle **decreases** with associated  $p_T$



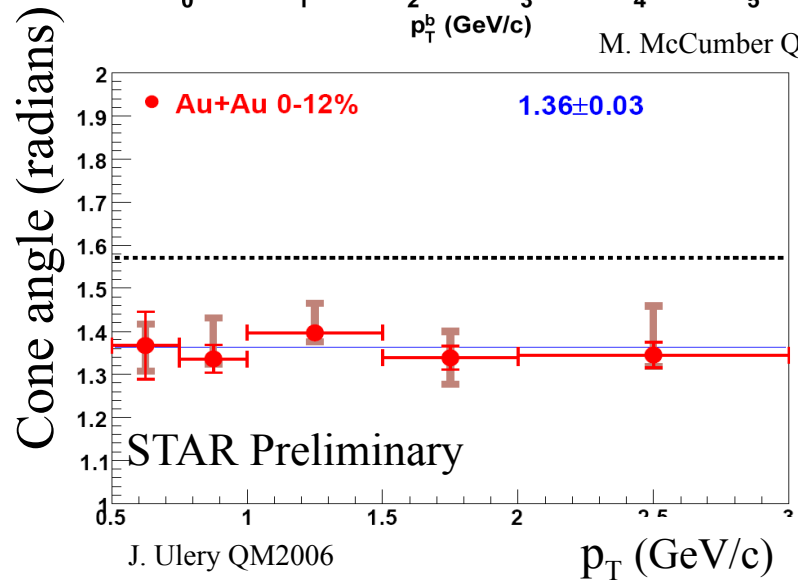
# Mach cone or Čerenkov gluons?

Angle predictions:

- **Mach-cone:**  
Angle **independent** of **associated  $p_T$**
- **Čerenkov gluon radiation:**  
Angle **decreases** with **associated  $p_T$**



M. McCumber QM2008



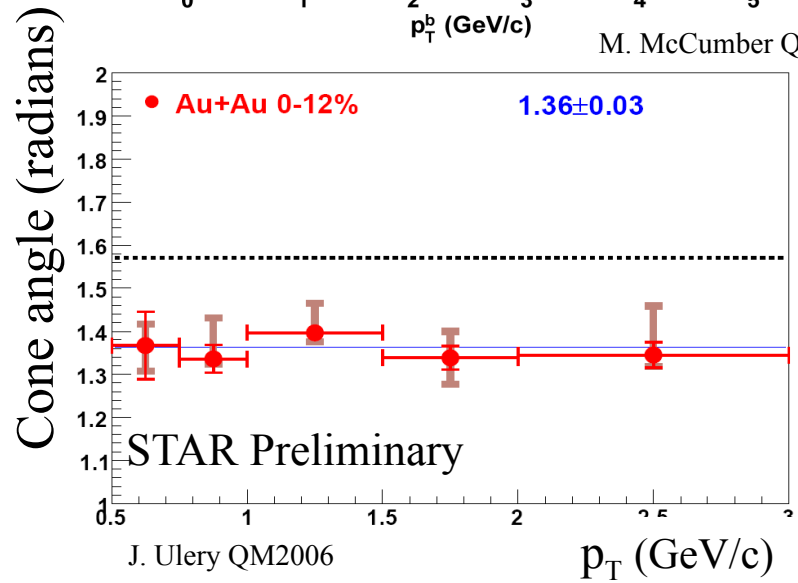
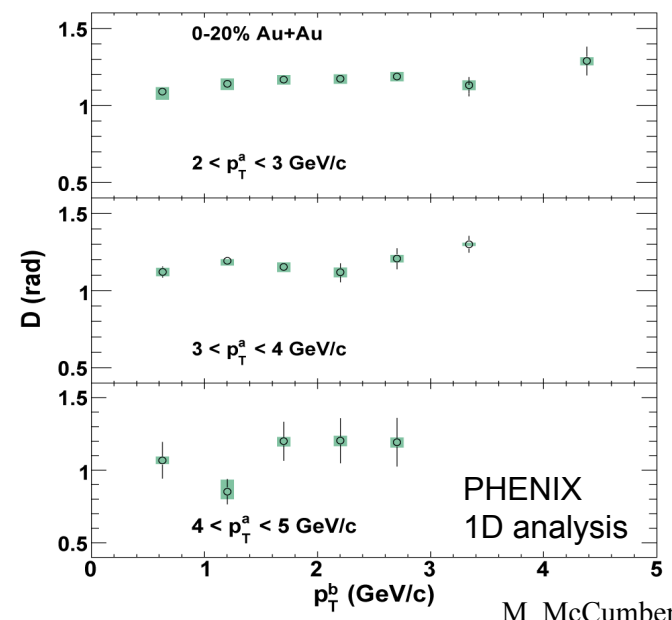
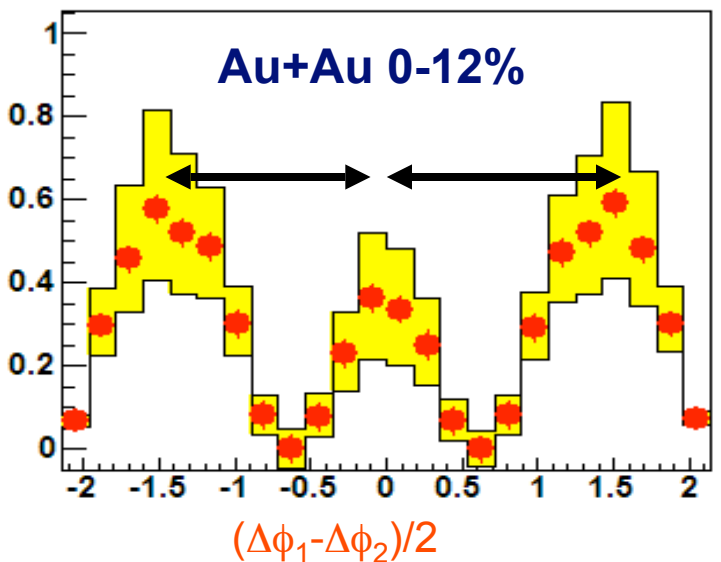
J. Ulery QM2006



# Mach cone or Čerenkov gluons?

Angle predictions:

- Mach-cone: ✓  
Angle independent of associated  $p_T$
- Čerenkov gluon radiation: ✗  
Angle decreases with associated  $p_T$

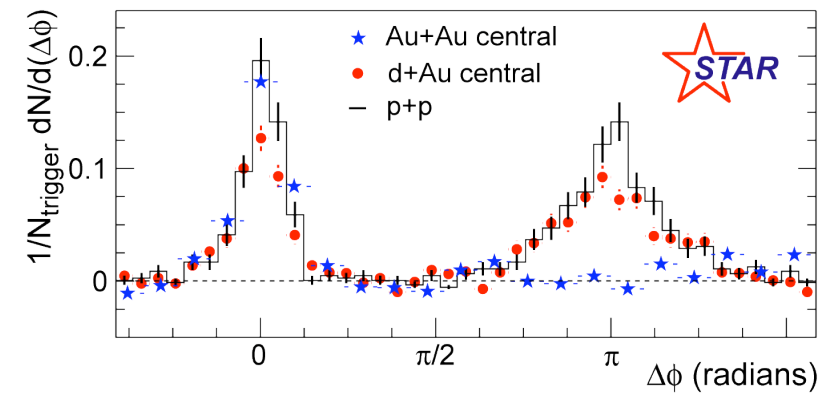


M. McCumber QM2008

J. Ulery QM2006

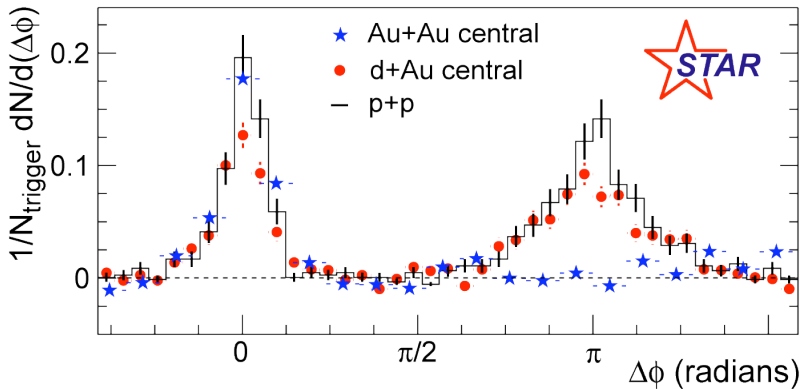
# Parton interactions on near side

## $\Delta(\phi)$ correlations

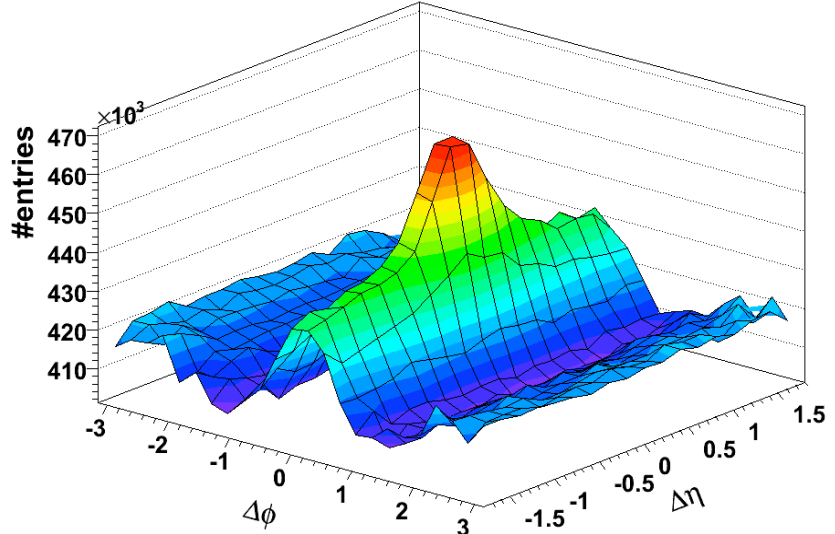


# Parton interactions on near side

## $\Delta(\phi)$ correlations



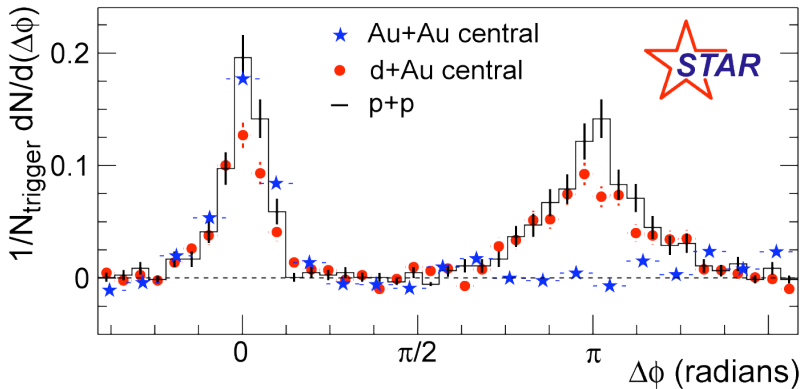
## $\Delta(\eta) - \Delta(\phi)$ correlations



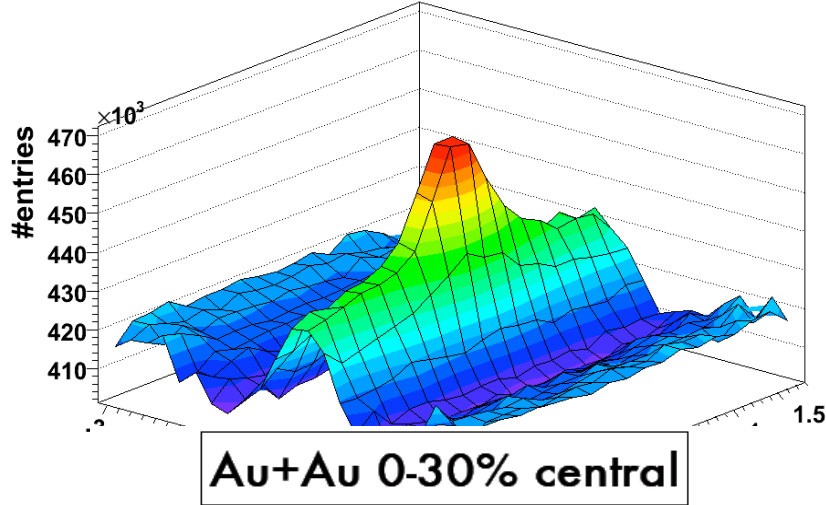
Long range  $\Delta(\eta)$  correlation  
– the “Ridge”

# Parton interactions on near side

$\Delta(\phi)$  correlations

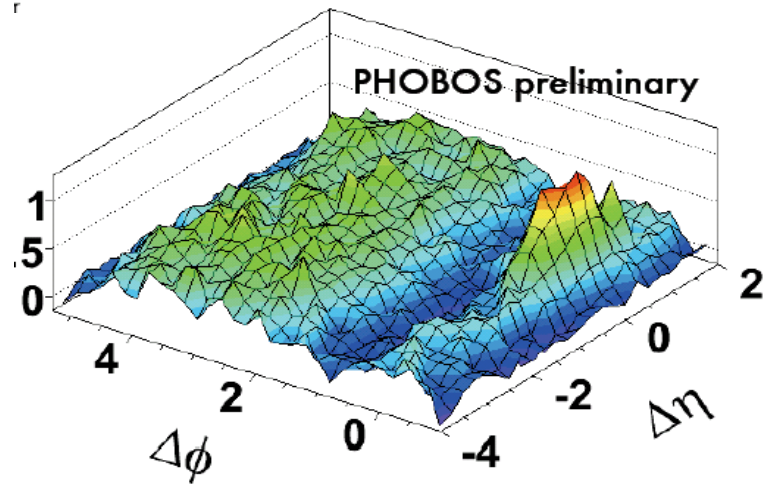


$\Delta(\eta) - \Delta(\phi)$  correlations



Long range  $\Delta(\eta)$  correlation  
 – the “Ridge”

Persists out to very large  $\Delta(\eta) > 2$



# Some possible explanations of the ridge

Recombination between thermal and shower partons at intermediate  $p_T$

R.C. Hwa & C.B. Chiu Phys. Rev. C 72 (2005) 034903

QCD bremsstrahlung radiation boosted by transverse flow

S.A. Voloshin, Phys. Lett. B 632 (2007) 490

E. Shuryak, hep-ph:0706.3531

In medium radiation and longitudinal flow push

N. Armesto et.al Phys. Rev. Lett. 93 (2007) 242301

Broadening of quenched jets in turbulent color fields

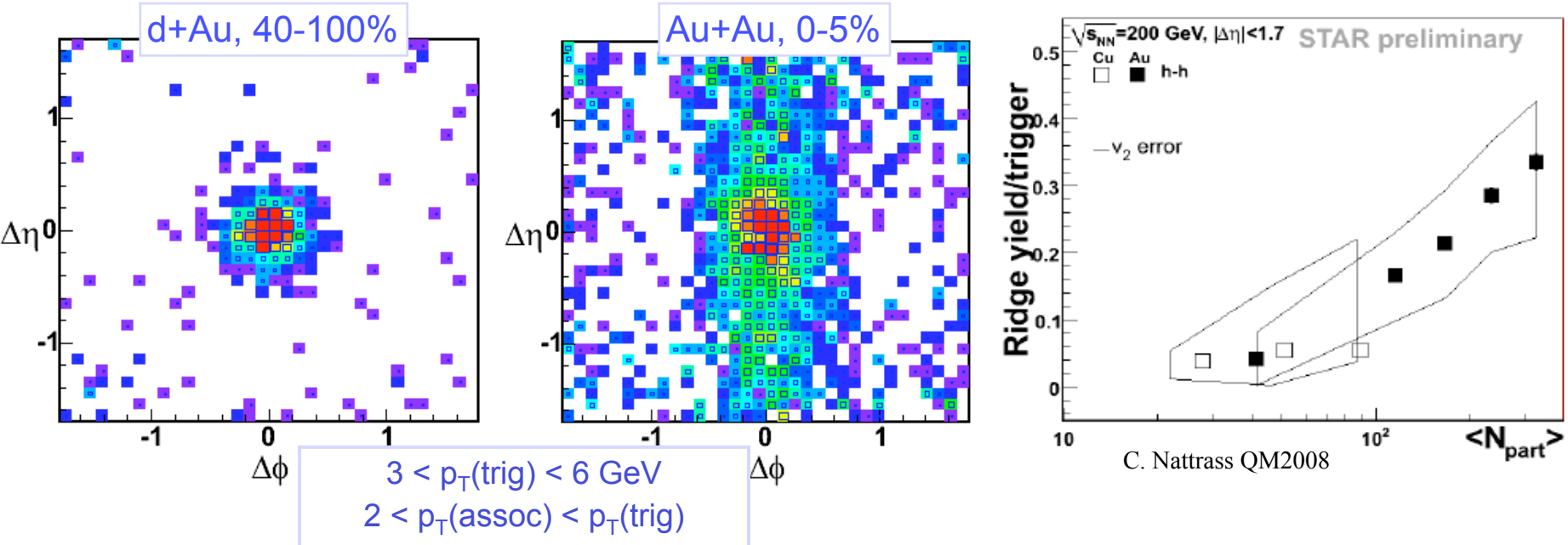
A. Majumder et.al Phys. Rev. Lett. 99 (2004) 042301

Momentum Kick Model

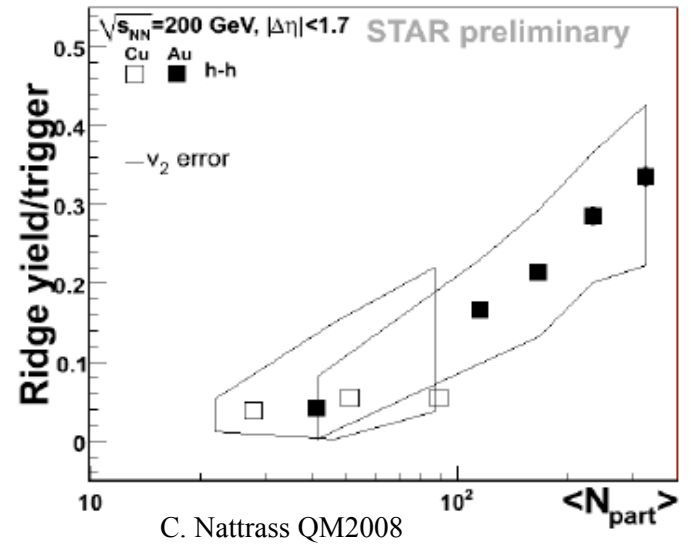
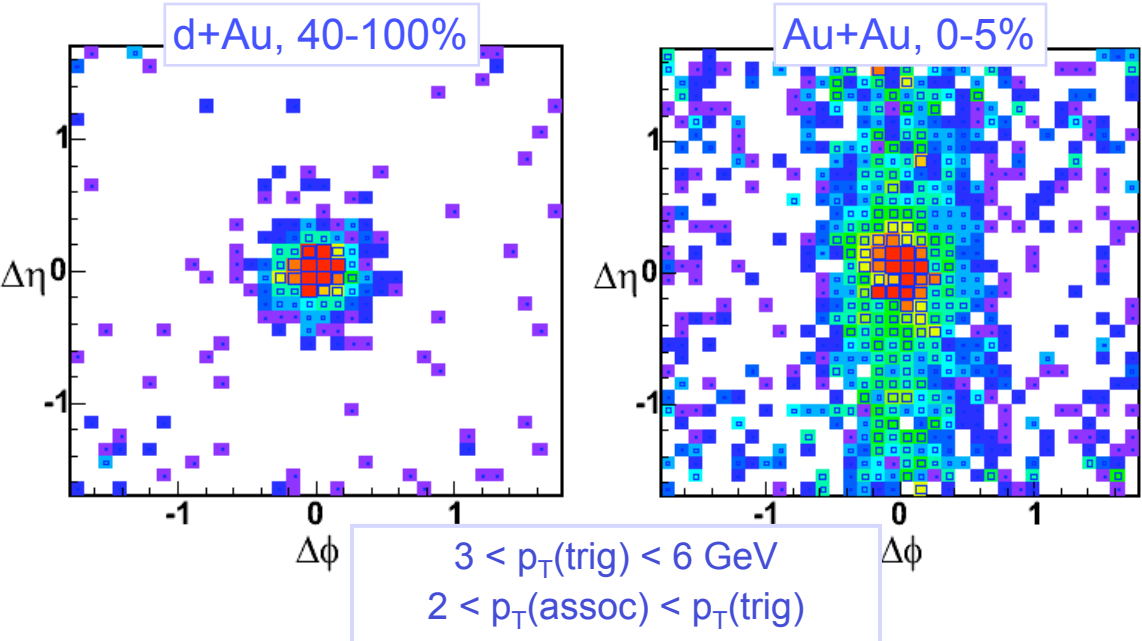
C.Y. Wong hep-ph:0712.3282

All qualitatively consistent with the features of the ridge

# Path length dependence of the ridge

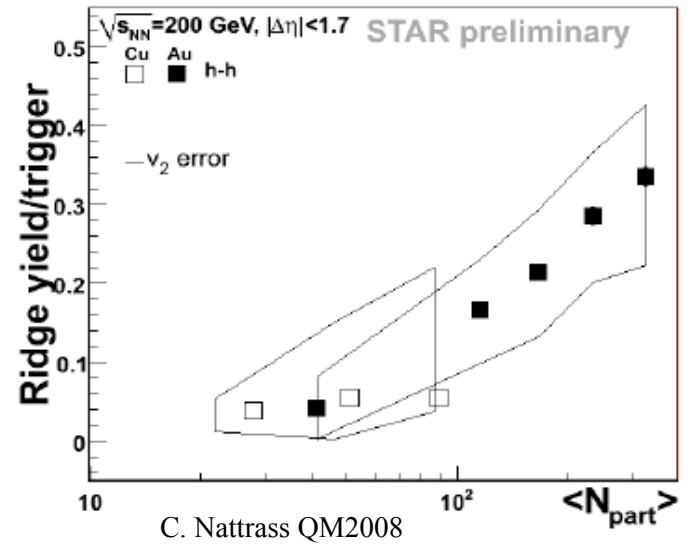
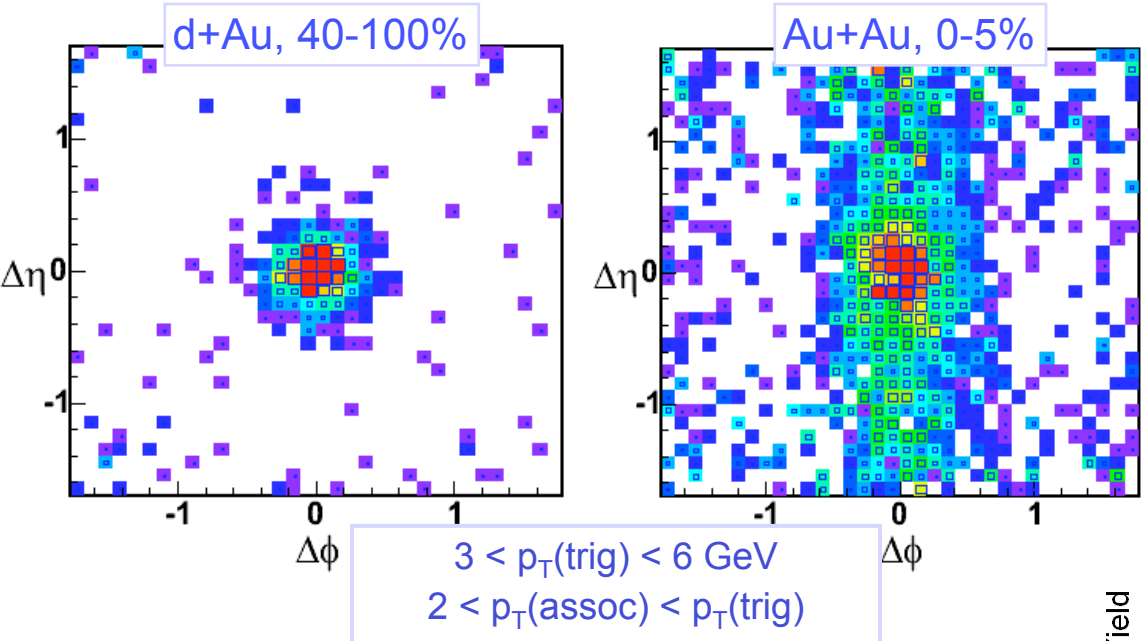


# Path length dependence of the ridge



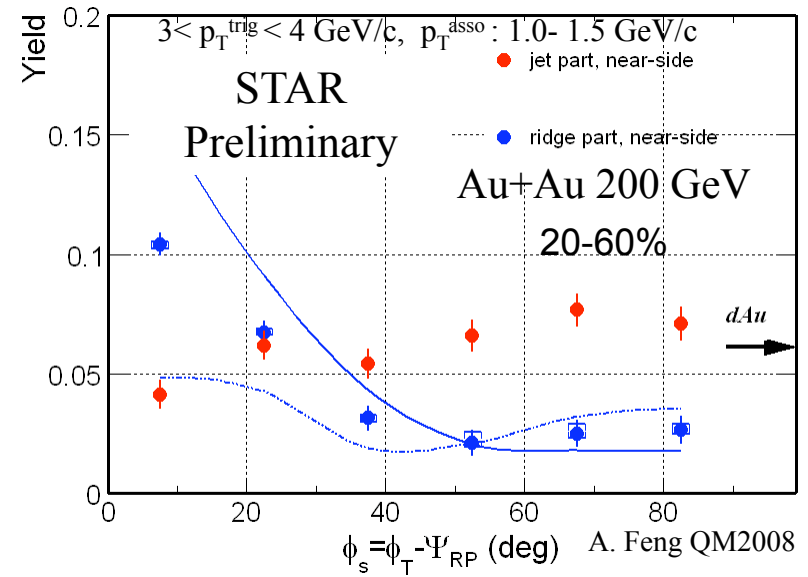
Ridge: Increases with  $N_{\text{part}}$   
 Independent of colliding system  
 Decreases with  $\phi_t - \Psi$

# Path length dependence of the ridge



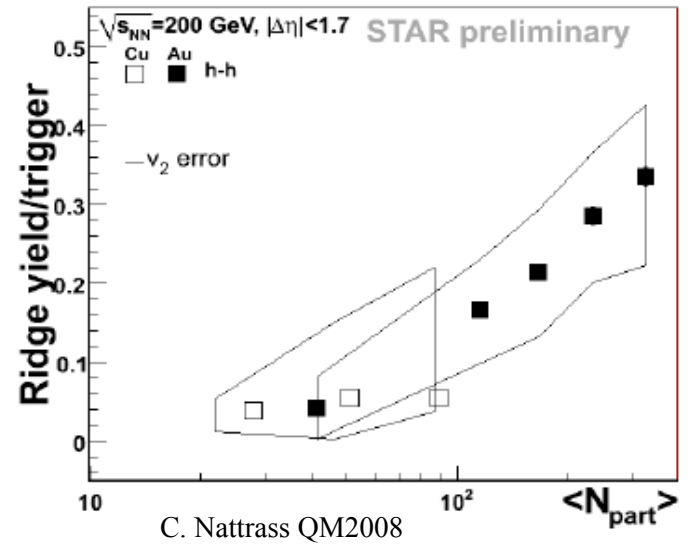
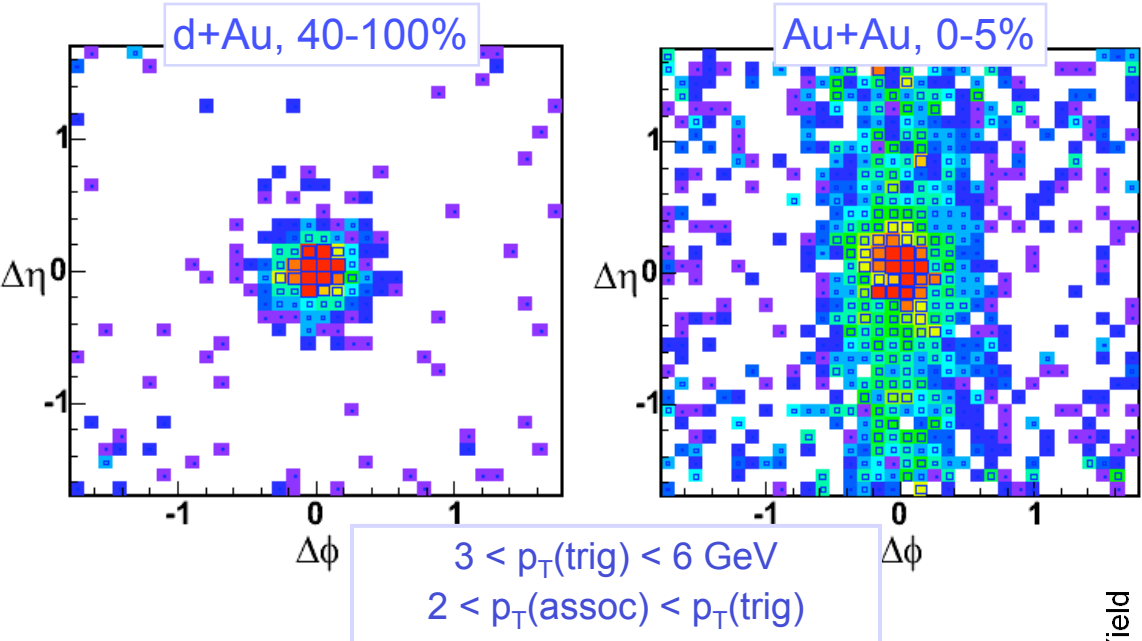
Ridge: Increases with  $N_{\text{part}}$   
 Independent of colliding system  
 Decreases with  $\phi_t - \Psi$

Jet: Slight to no increases with  $\phi_t - \Psi$   
 Au+Au  $\sim$  d+Au





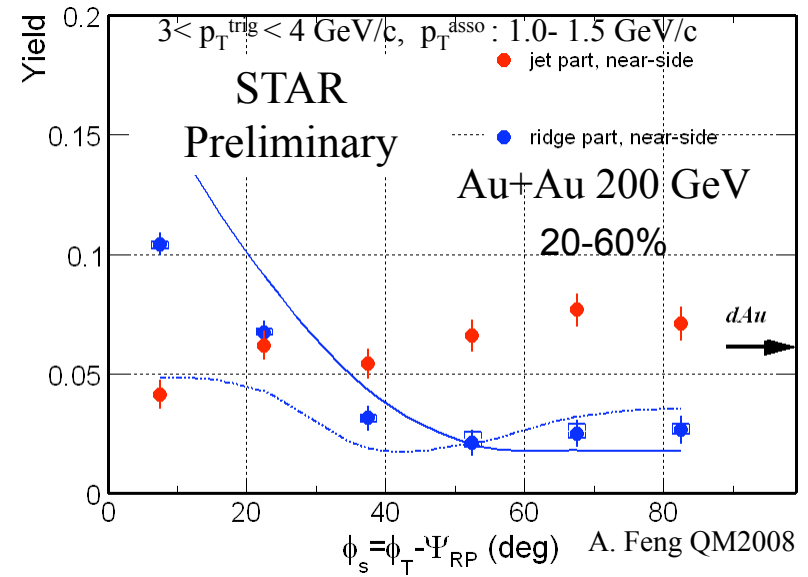
# Path length dependence of the ridge



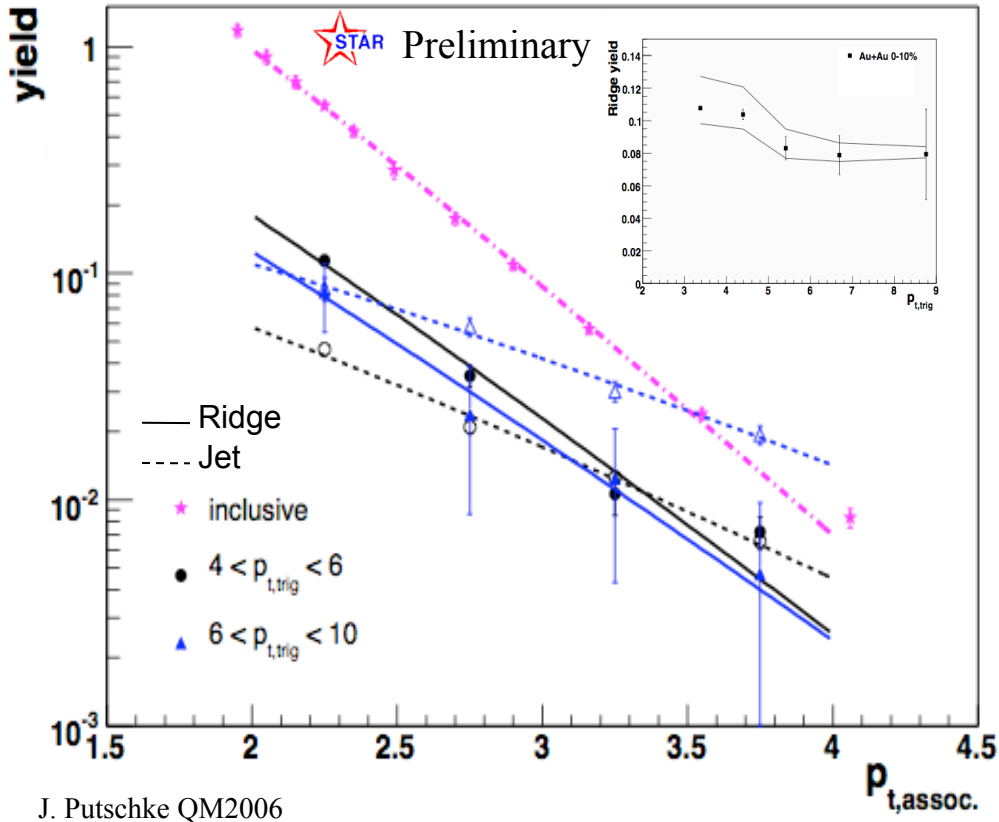
Ridge: Increases with  $N_{\text{part}}$   
 Independent of colliding system  
 Decreases with  $\phi_t - \Psi$

Jet: Slight to no increases with  $\phi_t - \Psi$   
 Au+Au  $\sim$  d+Au

Parton interacts with medium (ridge)  
 and then vacuum fragments (jet)?



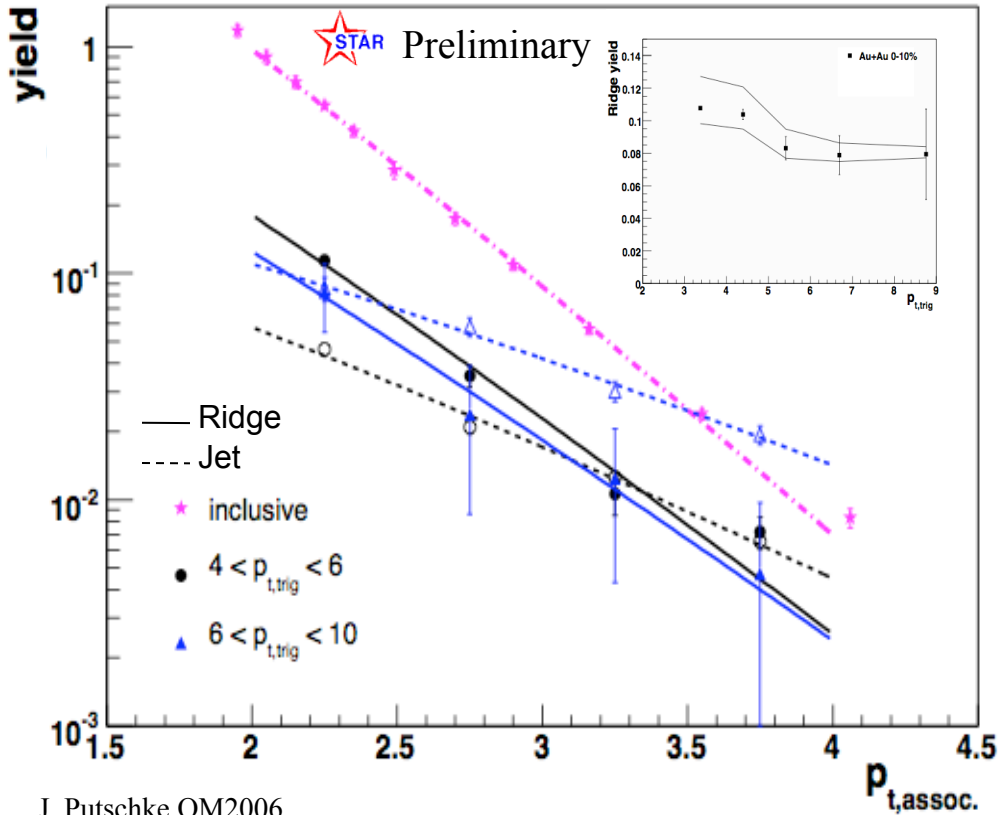
# Spectra of ridge and shoulder particles



J. Putschke QM2006

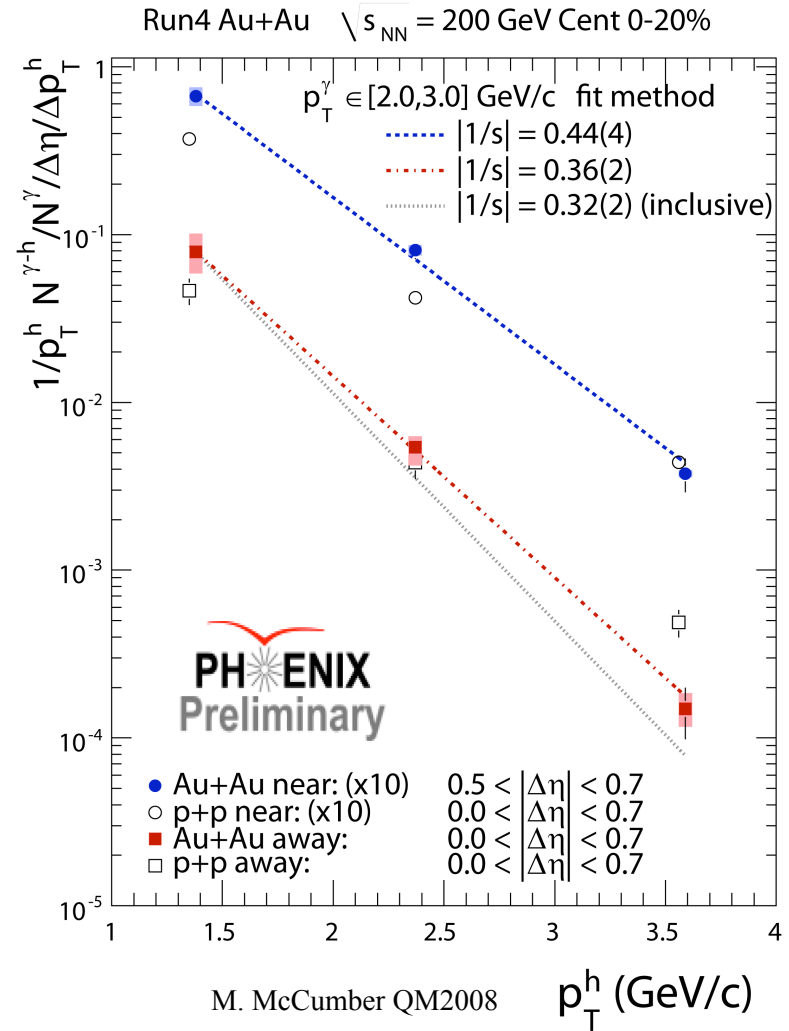
$\text{slope}_{\text{ridge}} > \text{slope}_{\text{jet}}$   
 $\sim \text{slope}_{\text{inclusive}}$

# Spectra of ridge and shoulder particles

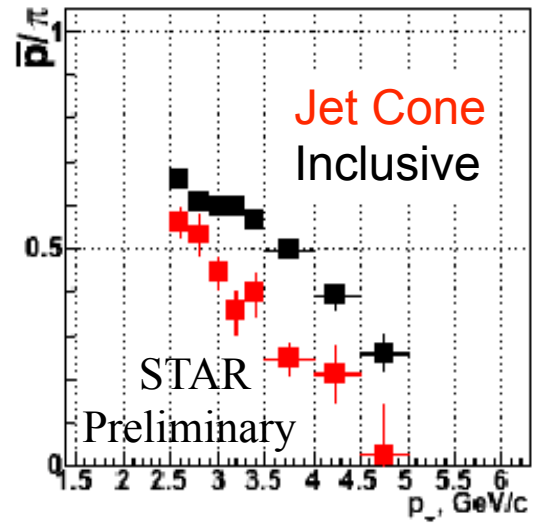
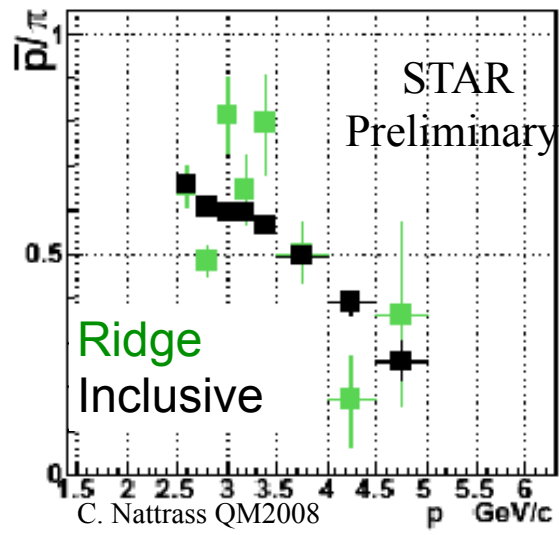


J. Putschke QM2006

$\text{slope}_{\text{ridge}} > \text{slope}_{\text{jet}}$   
 $\sim \text{slope}_{\text{inclusive}}$   
 $\geq \text{slope}_{\text{shoulder}}$

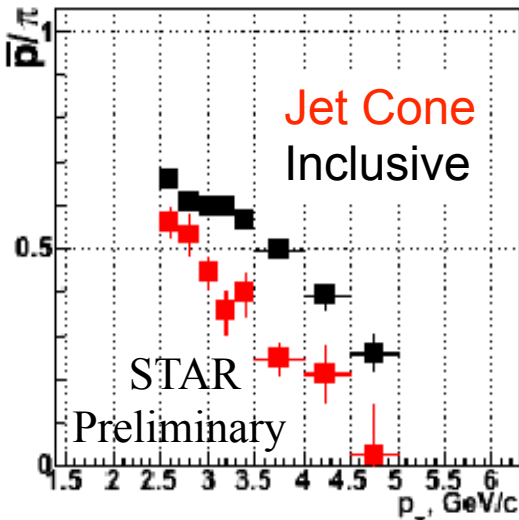
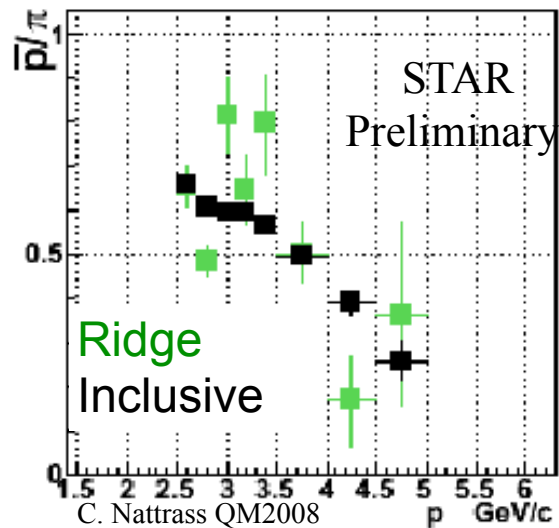


# Composition of ridge and shoulders



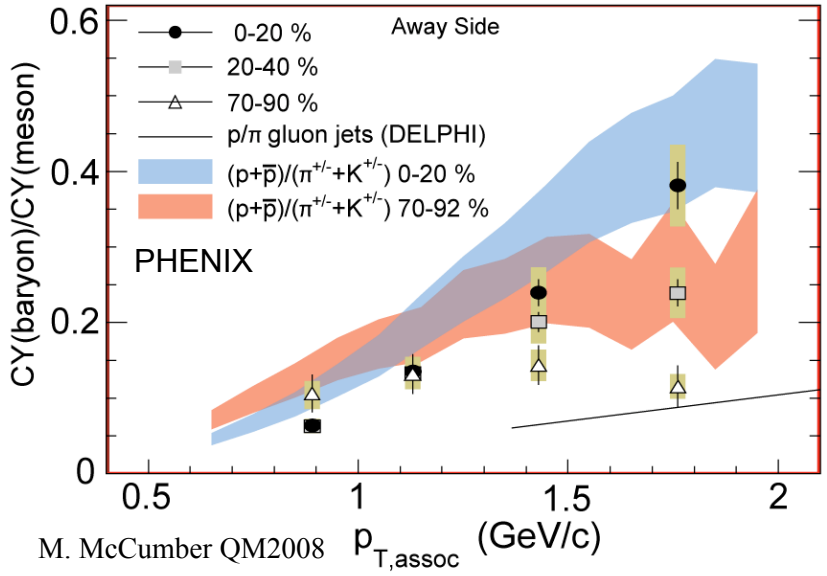
ridge ratio  $\sim$  inclusive ratio  
> jet ratio

# Composition of ridge and shoulders

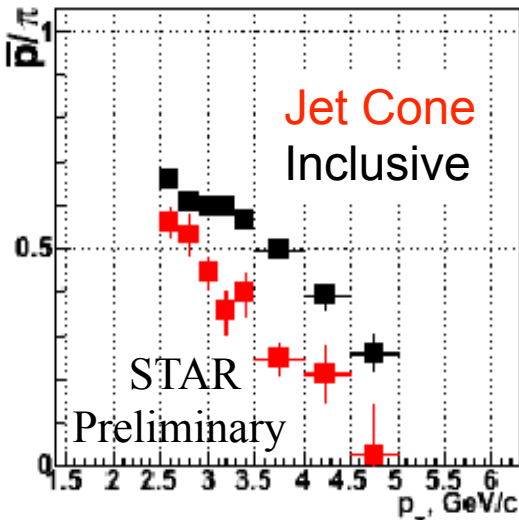
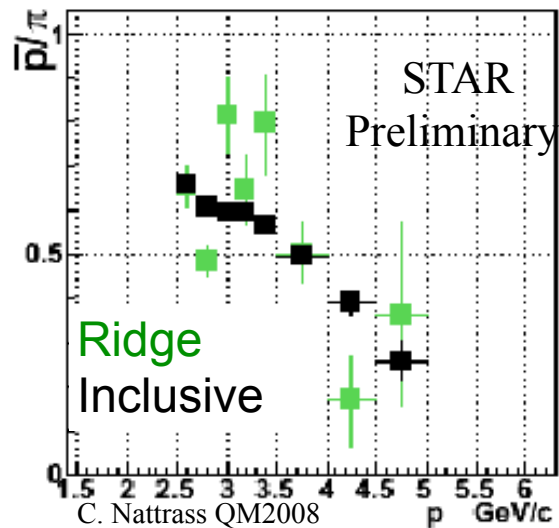


ridge ratio ~ inclusive ratio  
> jet ratio

shoulder ratio ~ inclusive ratio  
> jet ratio

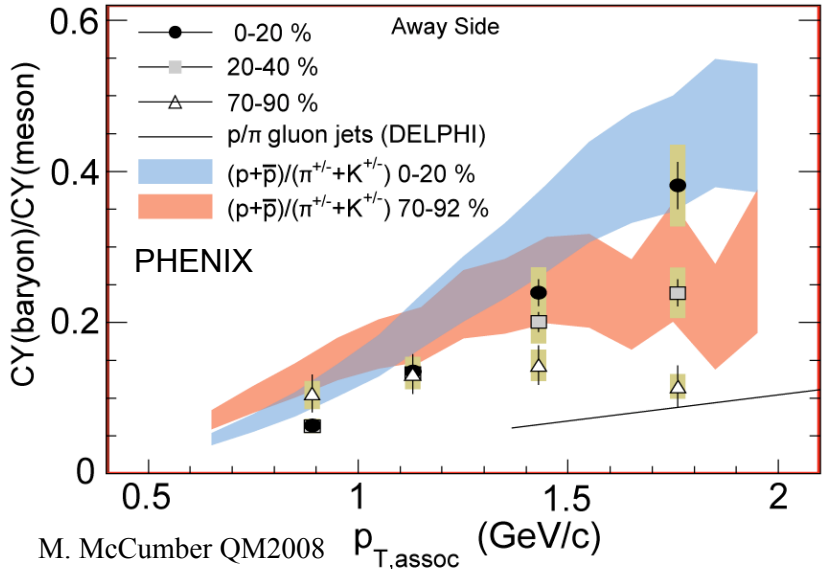


# Composition of ridge and shoulders



ridge ratio ~ inclusive ratio  
> jet ratio

shoulder ratio ~ inclusive ratio  
> jet ratio



Ridge and Shoulder similar properties  
**NOT** vacuum fragmentation

Energy lost by jet partons seems to be re-distributed into the medium and freezes out in similar fashion

# Un-triggered pair correlations

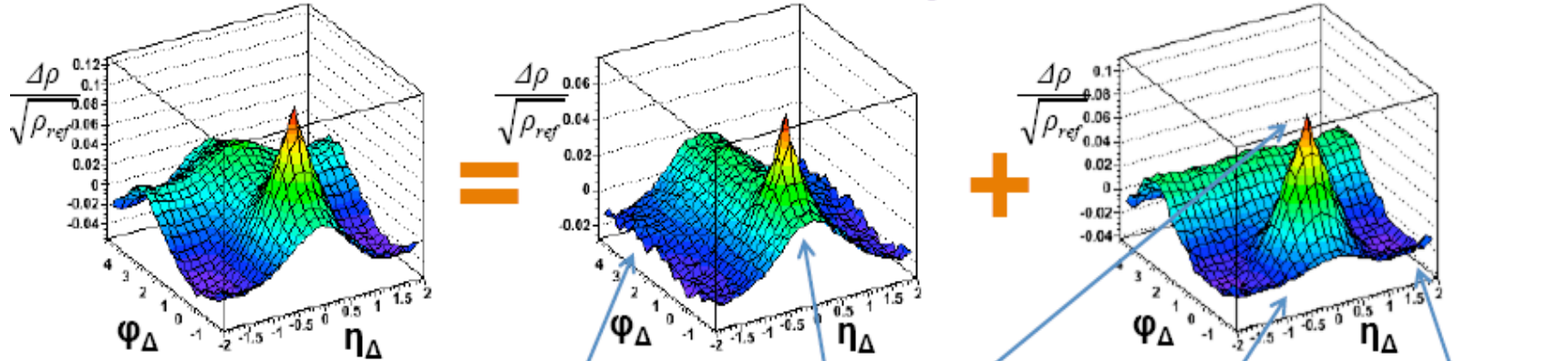
Method: measure pair densities  $\rho(\eta_1-\eta_2, \varphi_1-\varphi_2)$  for all possible pairs in same and mixed events.

Define correlation measure as:

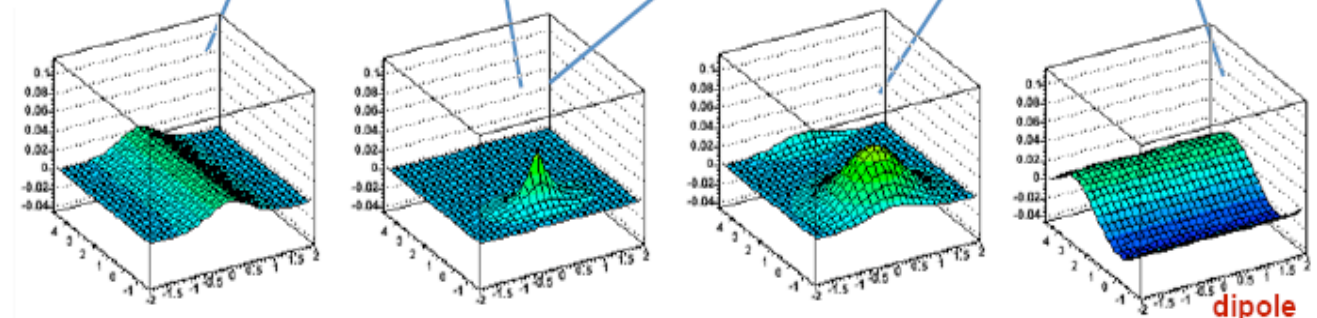
$$\frac{\rho_{same} - \rho_{mixed}}{\sqrt{\rho_{mixed}}} \equiv \frac{\Delta\rho}{\sqrt{\rho_{ref}}} \propto \frac{\# \text{ correlated pairs}}{\text{particle}}$$

Proton-Proton fit function

STAR Preliminary



Minijet:  
Same-side jet-like  
correlations with no  
trigger particle

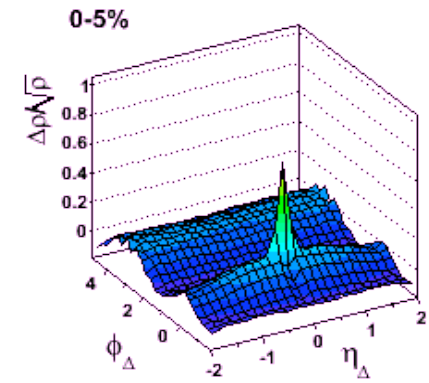
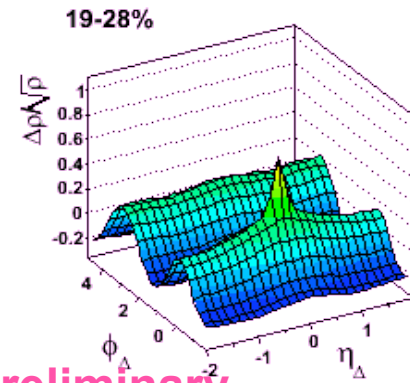
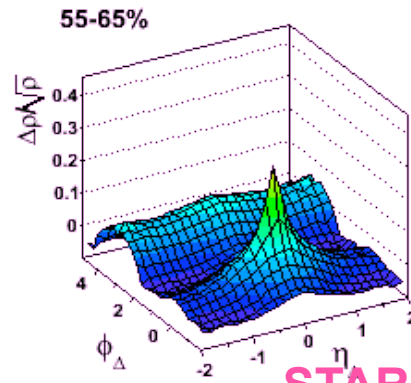
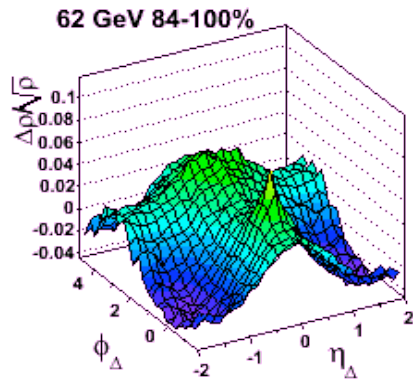
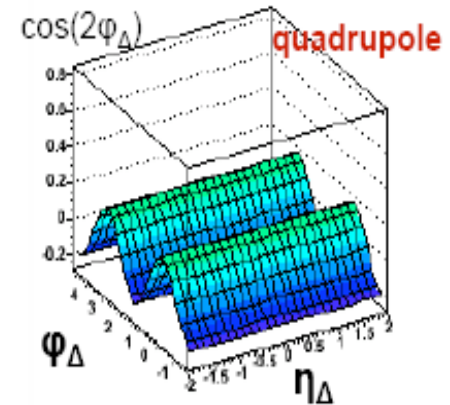


longitudinal fragmentation 1D gaussian    HBT and e+e- 2D exponential    Minijet Peak 2D gaussian    Away-side -cos(φ) dipole

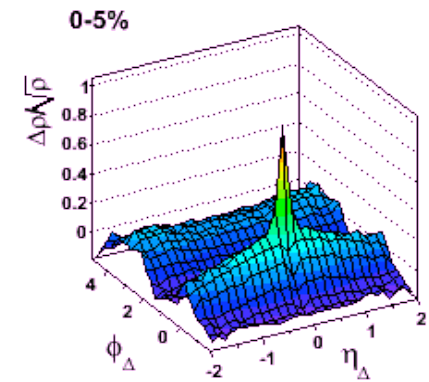
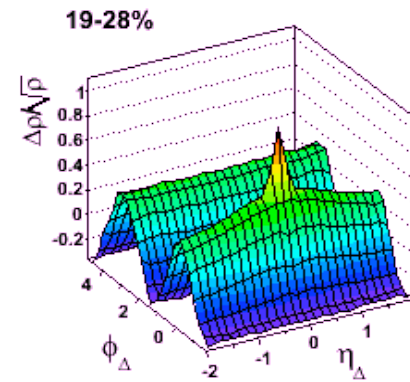
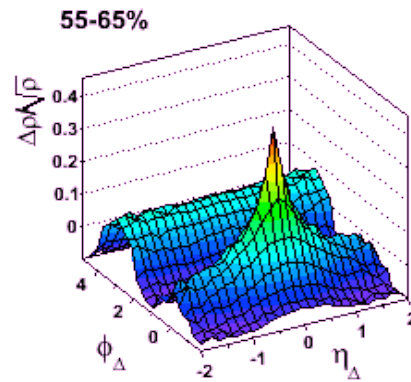
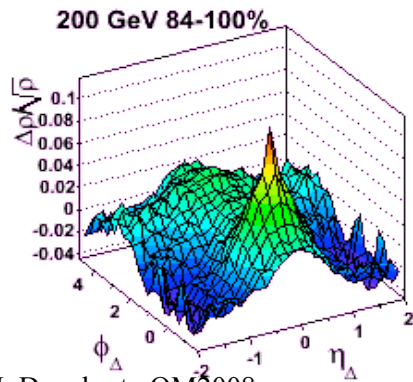
# Un-triggered pair correlations

## Au-Au fit function

Use proton-proton fit function +  $\cos(2\phi_\Delta)$  quadrupole term ("flow").  
 This gives the *simplest possible* way to describe Au+Au data.



STAR Preliminary



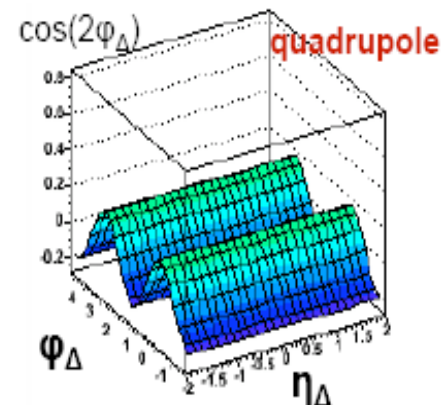


# Un-triggered pair correlations

## Au-Au fit function

Use proton-proton fit function +  $\cos(2\varphi_{\Delta})$  quadrupole term (“flow”).  
This gives the *simplest possible* way to describe Au+Au data.

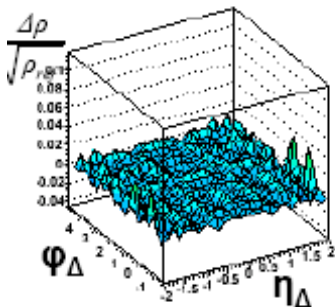
Small residual indicates goodness of fit



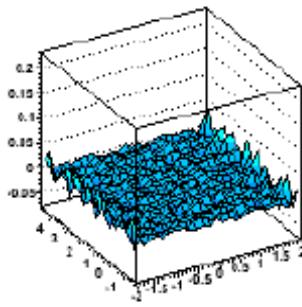
Fit residual = data - model

STAR Preliminary

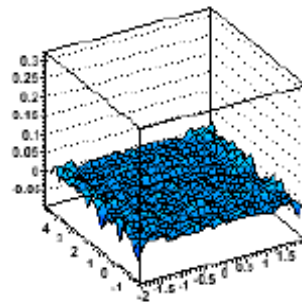
84-93%



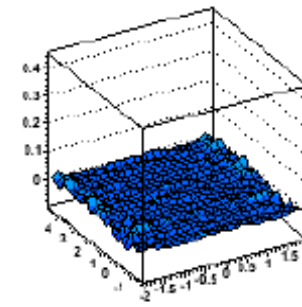
75-84%



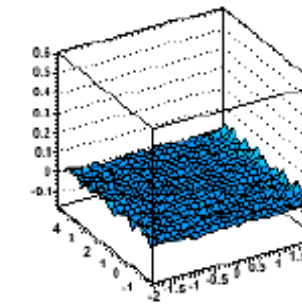
65-75%



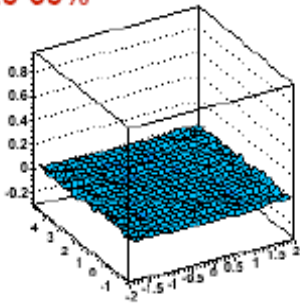
55-65%



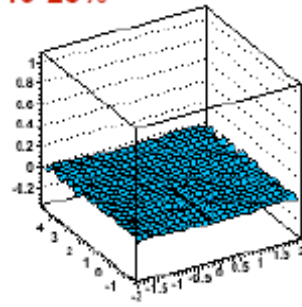
46-55%



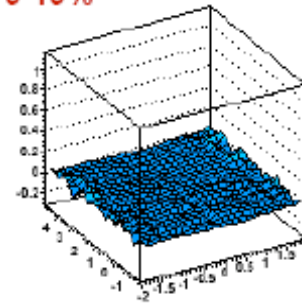
28-38%



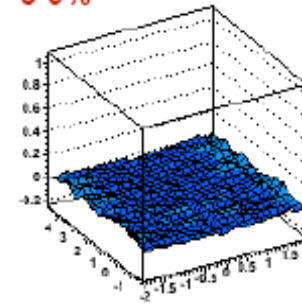
19-28%



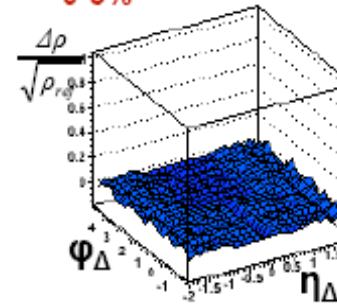
9-19%



5-9%

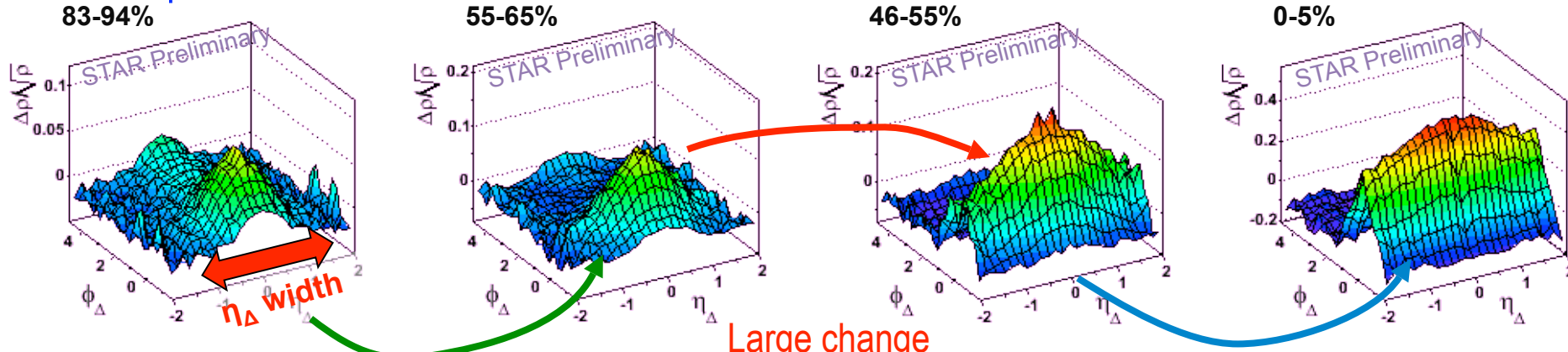


0-5%



# Evolution of mini-jet with centrality

Same-side peak



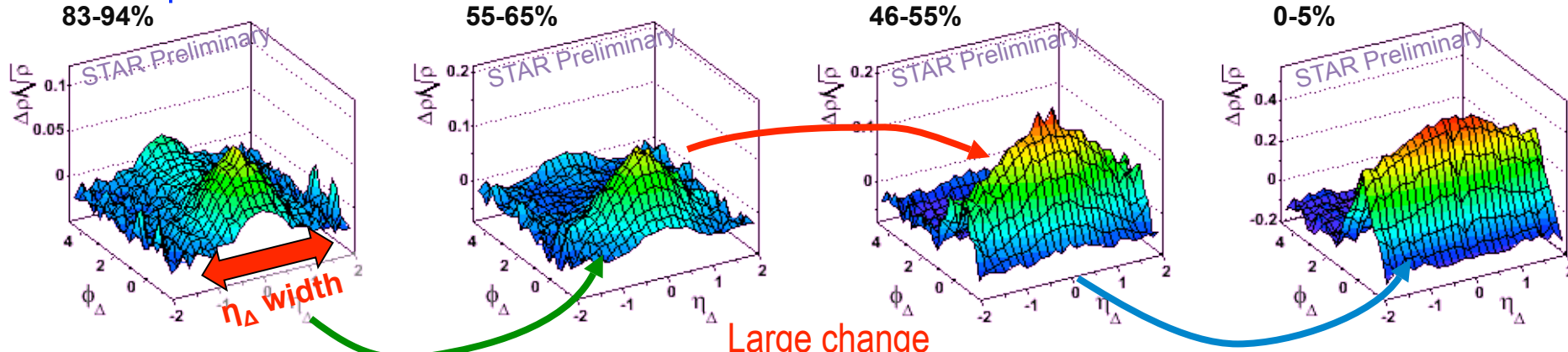
Little shape change from peripheral to 55% centrality

Large change within ~10% centrality

Smaller change from transition to most central

# Evolution of mini-jet with centrality

Same-side peak

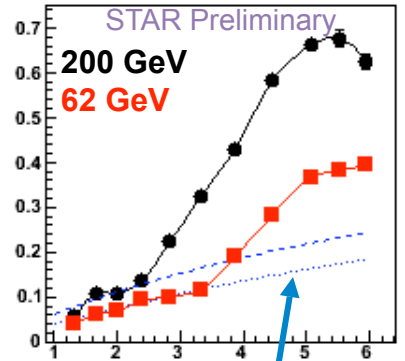


Little shape change from peripheral to 55% centrality

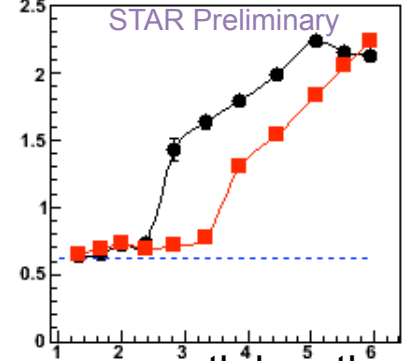
Large change within ~10% centrality

Smaller change from transition to most central

peak amplitude



peak  $\eta$  width



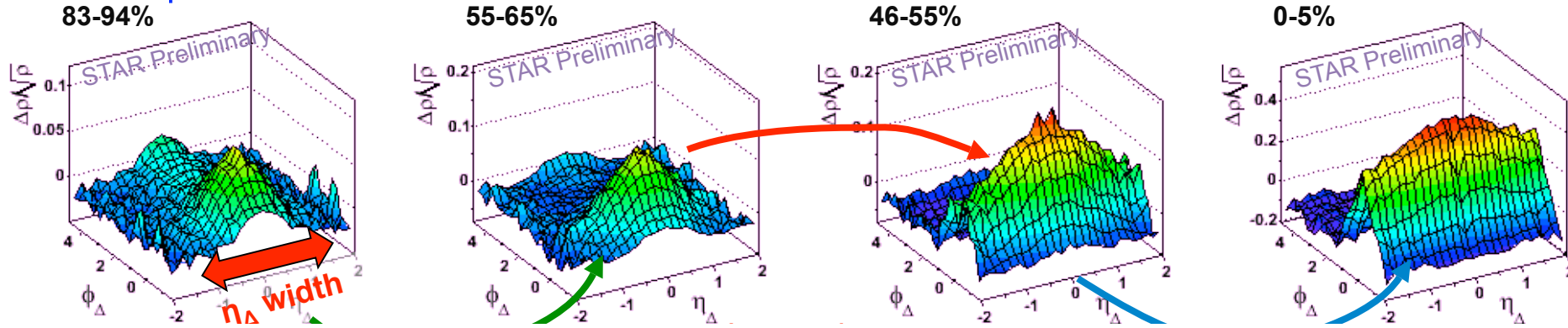
binary scaling references

$$v \equiv \frac{\langle N_{bin} \rangle}{\langle N_{part} / 2 \rangle}$$

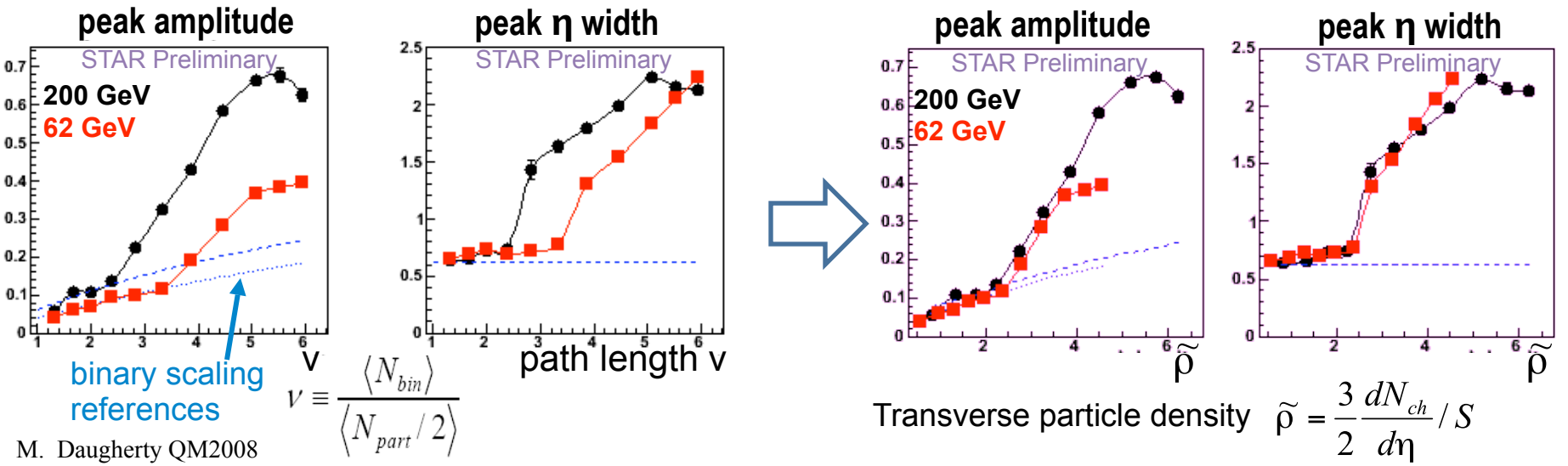
M. Daugherty QM2008

# Evolution of mini-jet with centrality

Same-side peak



Binary scaling reference followed until sharp transition at  $\rho \sim 2.5$   
 $\sim 30\%$  of the hadrons in central Au+Au participate in the same-side correlation

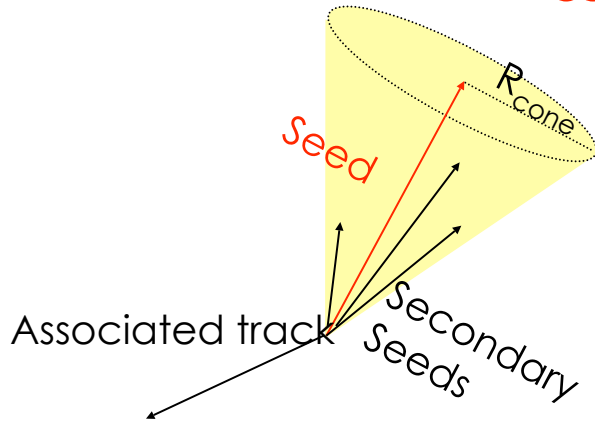


M. Daugherty QM2008

# Towards true jet reconstruction

- Reduce leading trigger particle biases from di-hadron correlations
- First step to jet reconstruction in A+A

## Multi-hadron trigger



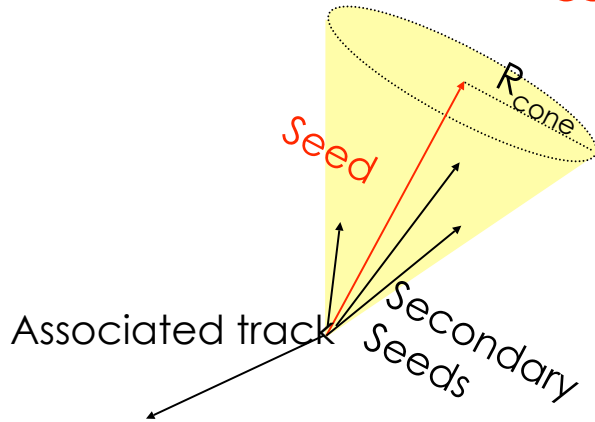
Use “cluster energy” as trigger:

- $R_{\text{cone}} = 0.3$
- $p_{T,\text{seed}} > 5 \text{ GeV}$
- $p_{T,\text{sec seed}} > 3 \text{ GeV}$

# Towards true jet reconstruction

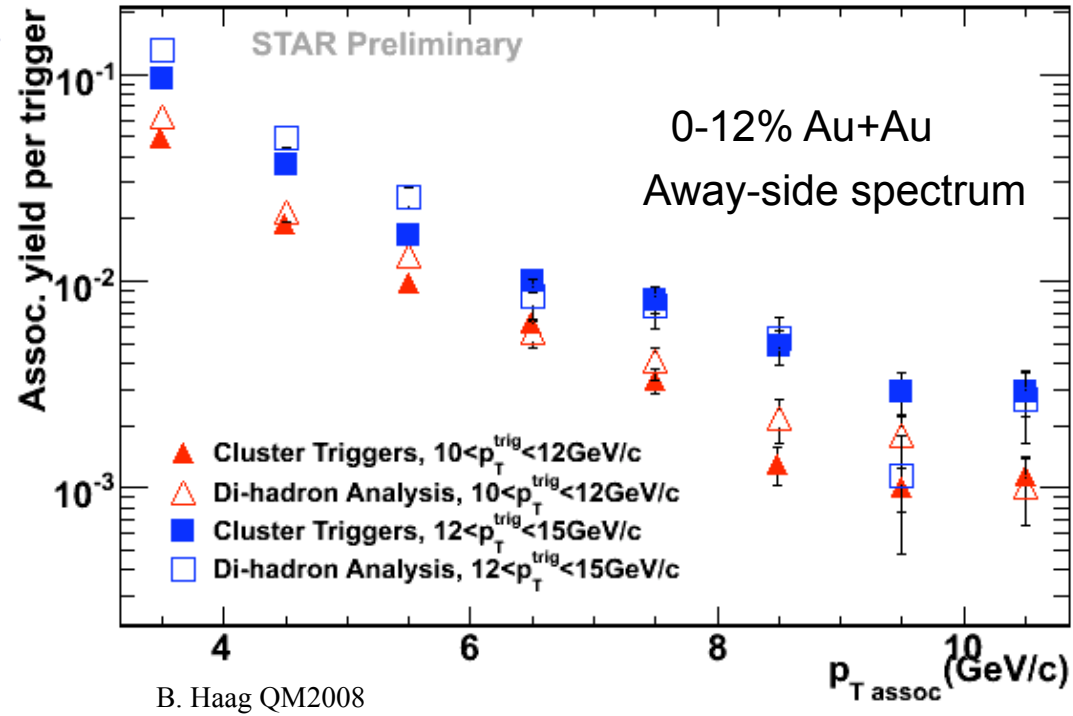
- Reduce leading trigger particle biases from di-hadron correlations
- First step to jet reconstruction in A+A

## Multi-hadron trigger



Use “cluster energy” as trigger:

- $R_{\text{cone}} = 0.3$
- $p_{T,\text{seed}} > 5 \text{ GeV}$
- $p_{T,\text{sec seed}} > 3 \text{ GeV}$



- **Single-hadron trig.**  $\approx$  **multi-hadron trig.**
- **Single high  $p_T$  triggered correlations** probe **jet-like correlations**

# Conclusions

Jets have been observed at RHIC

Making important steps towards constraining models that try to explain:

- How partons interact with and lose energy in the medium
- Where that energy goes
- How the medium changes with  $\sqrt{s}$ , centrality, and ion collided

# Conclusions

Jets have been observed at RHIC

Making important steps towards constraining models that try to explain:

- How partons interact with and lose energy in the medium
- Where that energy goes
- How the medium changes with  $\sqrt{s}$ , centrality, and ion collided

RHIC on the threshold of new era of quantitative comparison between theory and experiment that will characterize the properties of the remarkable new matter discovered at here



# BACKUPS

# At RHIC we've created a new state of matter

- The QGP is the:

*hottest* ( $T=200-400 \text{ MeV} \sim 2.5 \cdot 10^{12} \text{ K}$ )

*densest* ( $\varepsilon = 30-60 \varepsilon_{\text{nuclear matter}}$ )

matter ever studied in the lab.

It flows as a

*(nearly) perfect fluid*

with systematic patterns, consistent with

*quark degree of freedom*

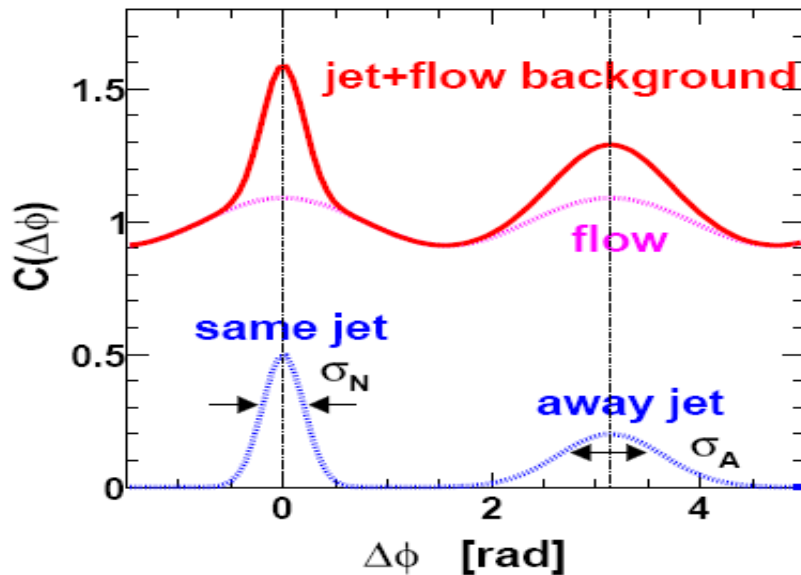
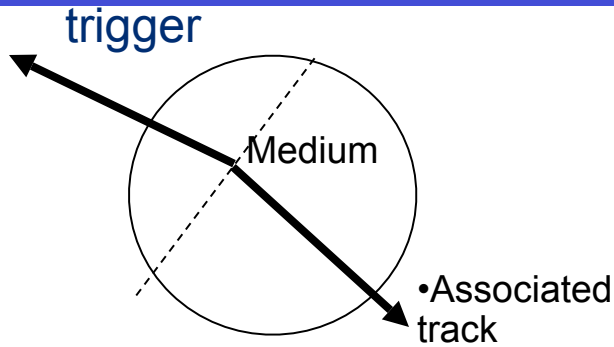
and a viscosity to entropy density ratio

*lower*

than any other known fluid.

Now want to learn more about properties

# 2 particle angular correlations

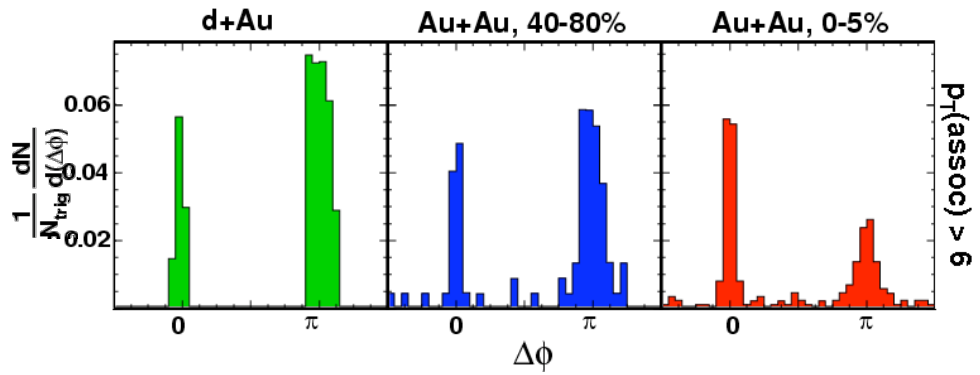


$$C(\Delta\phi) \equiv \frac{Y_{same}(\Delta\phi)}{Y_{mixed}(\Delta\phi)} \times \frac{\int Y_{mixed}(\Delta\phi) d\phi}{\int Y_{same}(\Delta\phi) d\phi}$$

$$C(\Delta\phi) \equiv b_0 \left[ 1 + 2 \langle v_2^{assoc} \rangle \langle v_2^{trig} \rangle \cos(2\Delta\phi) \right] + J(\Delta\phi)$$

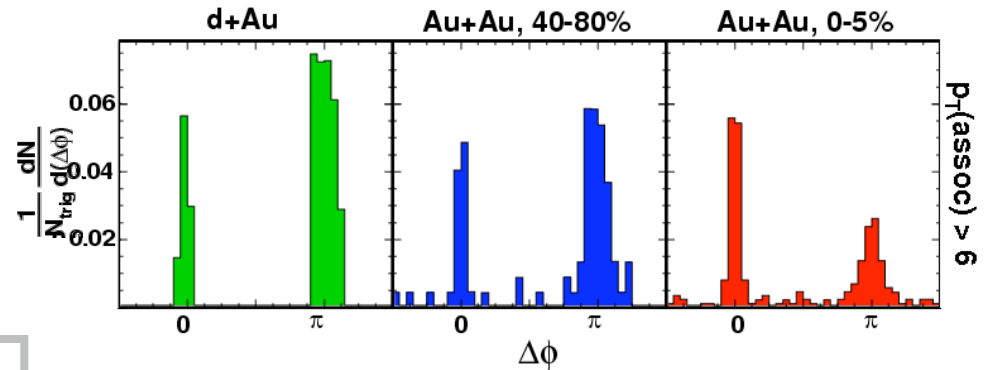
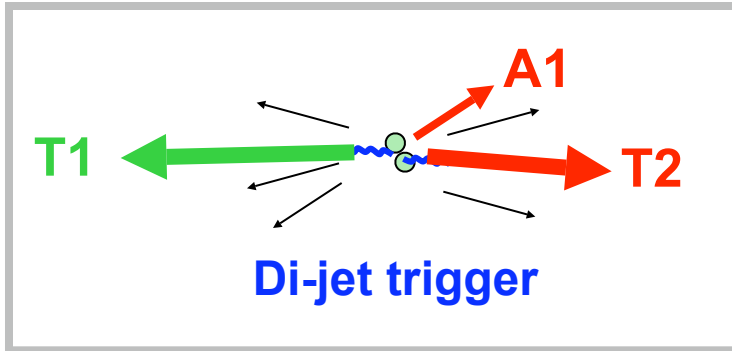
# Di-jet triggered correlations

Observation of di-jets:  
punch through



# Di-jet triggered correlations

Observation of di-jets:  
punch through

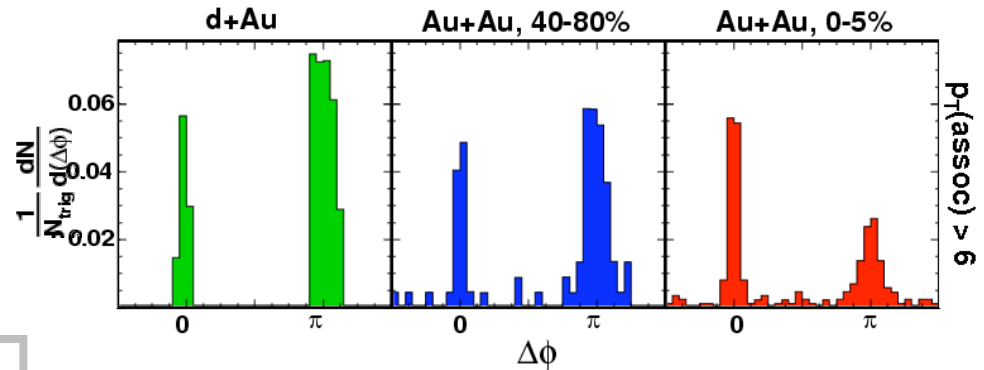
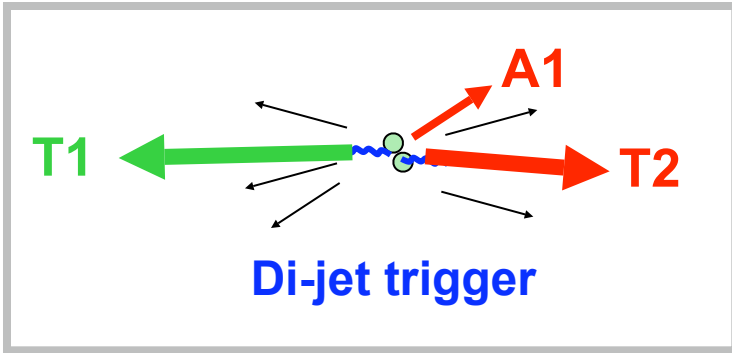


Select di-jets events:  
Require **T1** and **T2** b-to-b

**T1**:  $p_T > 5 \text{ GeV}/c$     **T2**:  $p_T > 4 \text{ GeV}/c$   
**A1**:  $p_T > 1.5 \text{ GeV}/c$

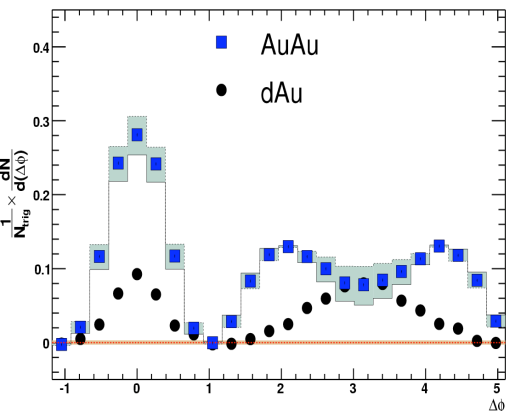
# Di-jet triggered correlations

Observation of di-jets:  
punch through

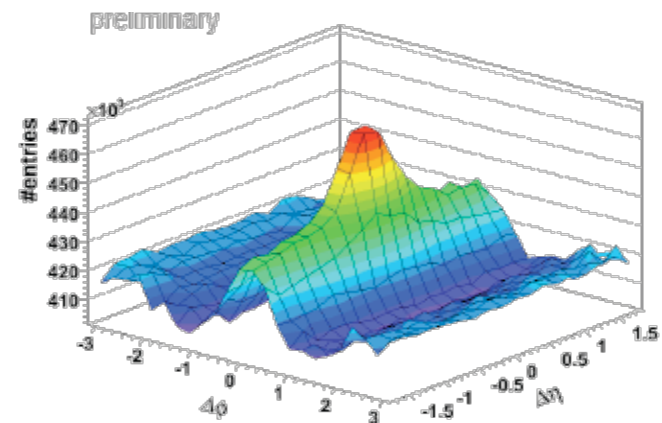


Select di-jets events:  
Require **T1** and **T2** b-to-b

**T1**:  $p_T > 5 \text{ GeV}/c$     **T2**:  $p_T > 4 \text{ GeV}/c$   
**A1**:  $p_T > 1.5 \text{ GeV}/c$

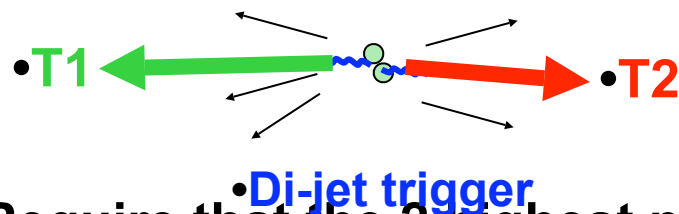


What happens to  
away-side hump and  
near-side ridge if we  
trigger on di-jets ?



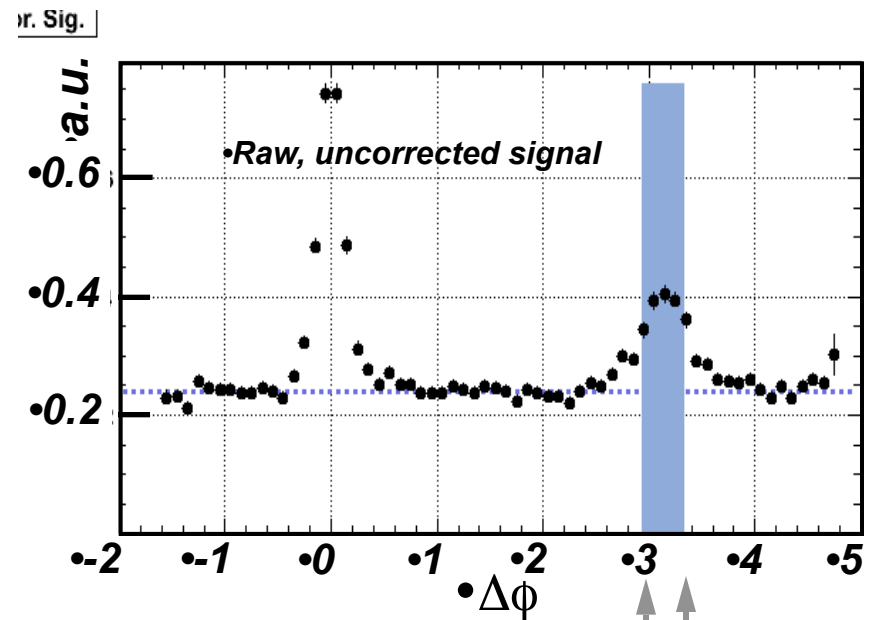
# Data analysis: di-jet selection

- Correlation between primary trigger (**T1**) and “away-jet-axis trigger” (**T2**).



- Require that the 2 highest  $p_T$  particles are back-to-back in  $\phi$ .
- Assume this defines the jet-axis, look in 2D-space about the second trigger.

## T1T2 correlation

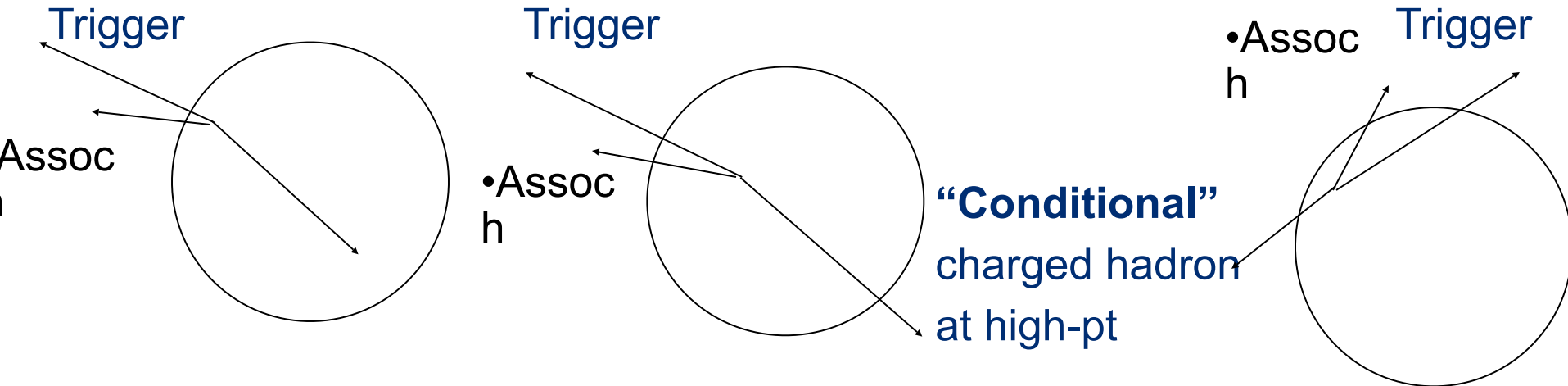


- T1:  $p_T > 5$  GeV/c
- T2:  $p_T > 4$  GeV/c

$\pi \pm 0.2$

# Change the surface-bias of near-side?

- Hope to shift distribution of hard scattering towards center of medium. Near-side parton travels through more medium

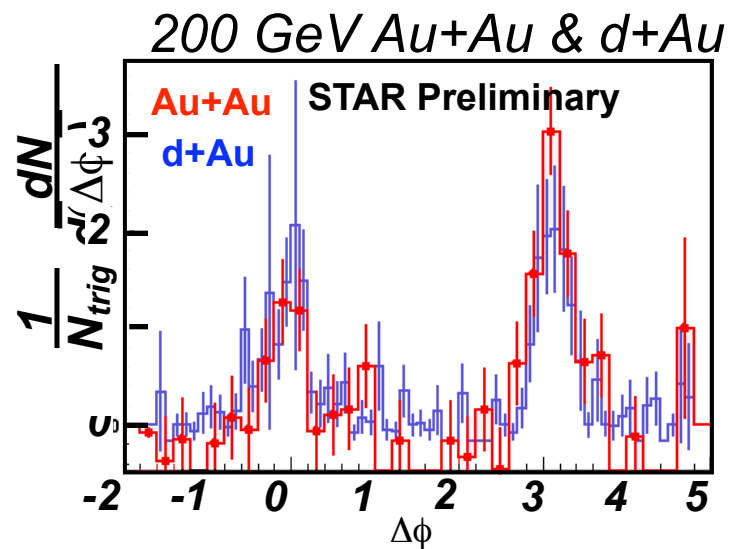


Create path lengths **comparable** in dense medium.

However not always from center could be tangential



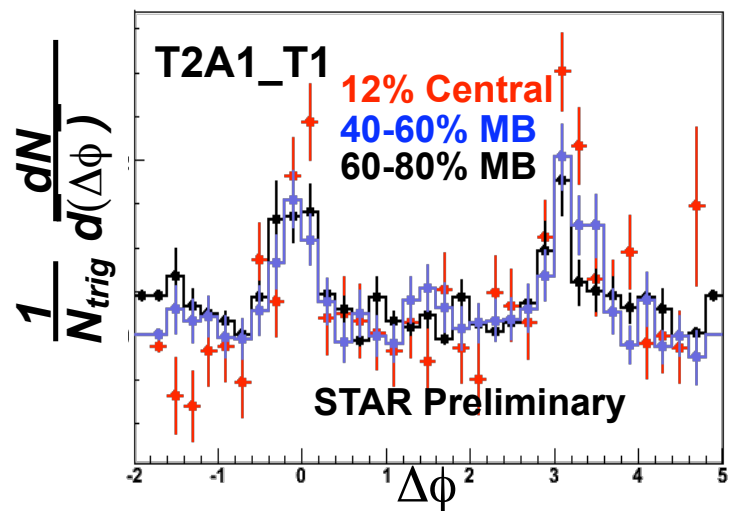
# Di-jet triggered correlations



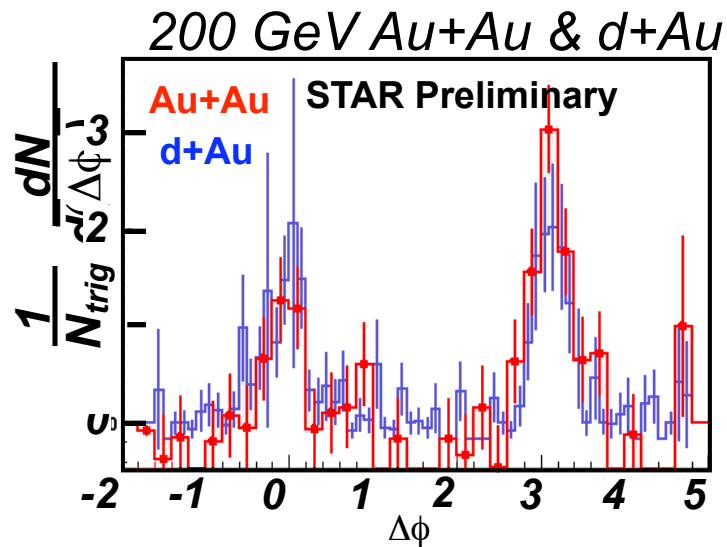
Di-jets are suppressed

Once selected:

- No Away-side suppression  
Au+Au  $\sim$  d+Au
- No Away-side shape modification



# Di-jet triggered correlations



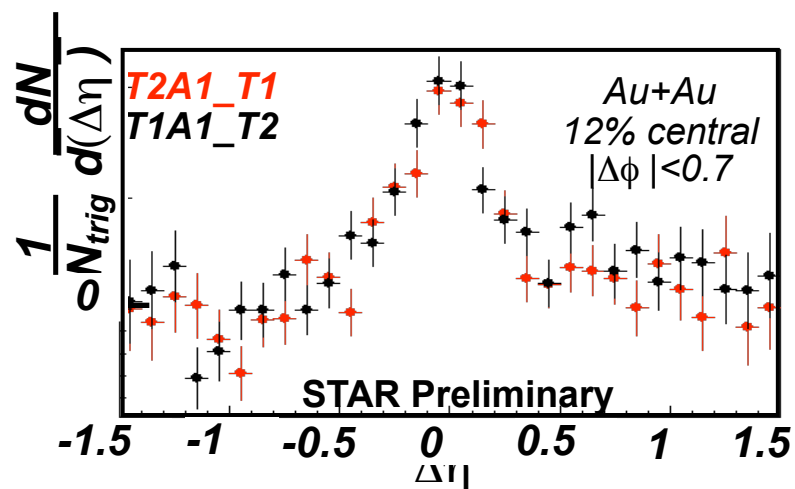
Di-jets are suppressed

Once selected:

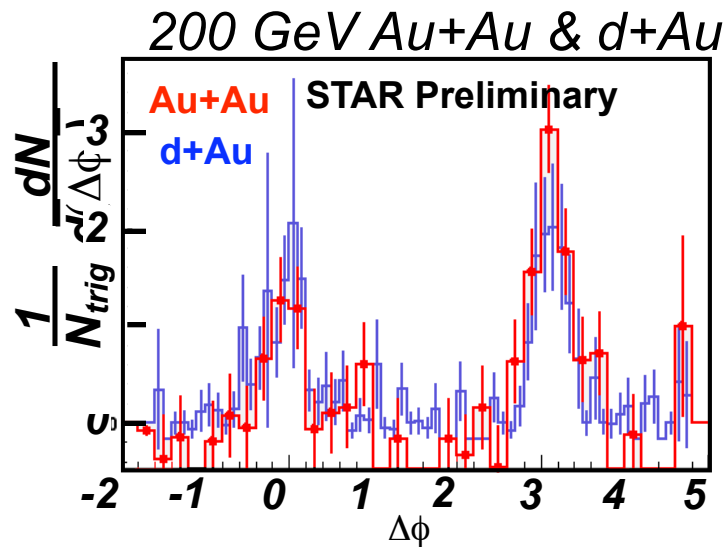
- No Away-side suppression  
Au+Au  $\sim$  d+Au

- No Away-side shape modification

- No Ridge



# Di-jet triggered correlations



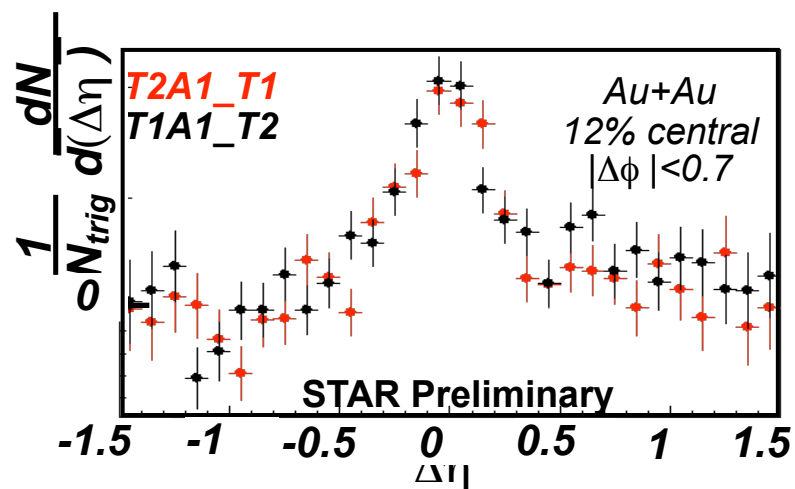
Di-jets are suppressed

Once selected:

- No Away-side suppression  
Au+Au  $\sim$  d+Au

- No Away-side shape modification

- No Ridge



Di-Jets don't interact with medium. Tangential jets or punch through without interaction ?

# Mach-Cone

•Trigger

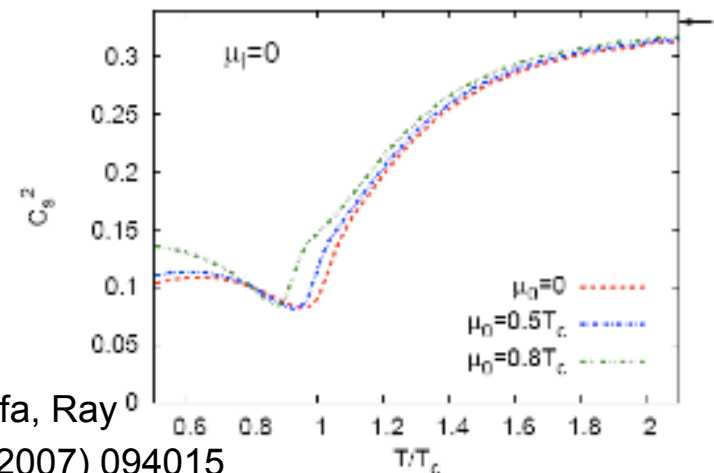
$$\frac{c_s}{v_{parton}} = \cos(\theta_M)$$

$$c_s^2 = \frac{\partial p}{\partial \epsilon}; \quad v_{parton} \approx c$$

- Mach angle depends on speed of sound in medium
  - T dependent
- Angle independent of associated  $p_T$ .

•Away-side

•PNJL Model



•Mikherjee, Mustafa, Ray

•Phys. Rev. D75 (2007) 094015

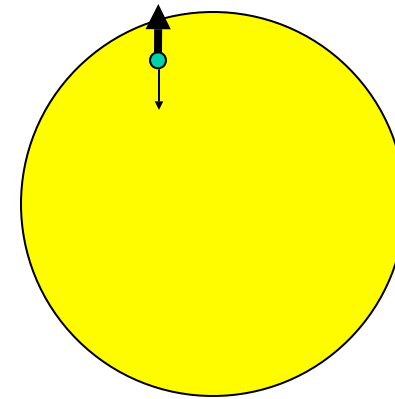
# Mach-Cone

$$\frac{c_s}{v_{parton}} = \cos(\theta_M)$$

$$c_s^2 = \frac{\partial p}{\partial \epsilon}; \quad v_{parton} \approx c$$

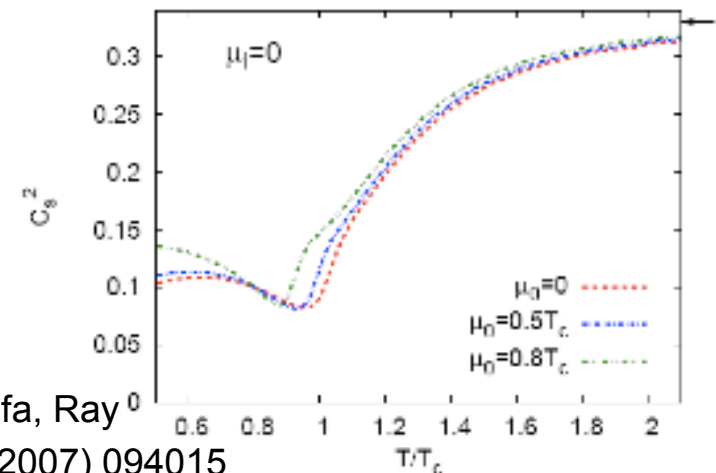
- Mach angle depends on speed of sound in medium
  - T dependent
- Angle independent of associated  $p_T$ .

•Trigger



•Away-side

•PNJL Model



•Mikherjee, Mustafa, Ray

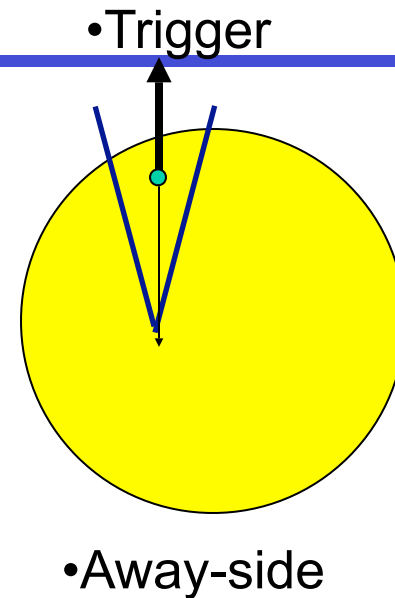
•Phys. Rev. D75 (2007) 094015

# Mach-Cone

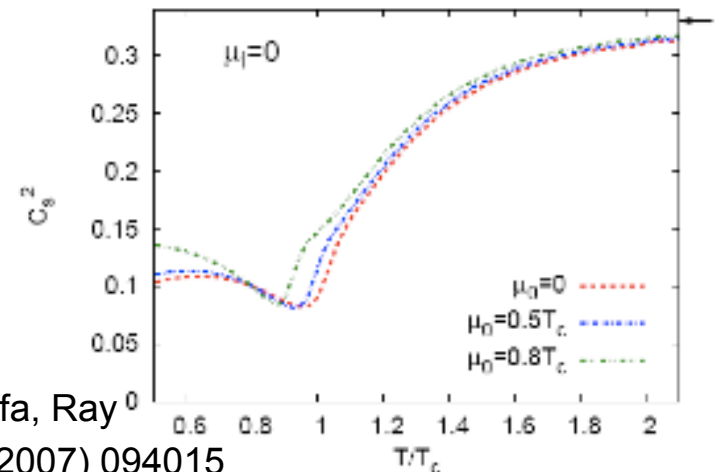
$$\frac{c_s}{v_{parton}} = \cos(\theta_M)$$

$$c_s^2 = \frac{\partial p}{\partial \epsilon}; \quad v_{parton} \approx c$$

- Mach angle depends on speed of sound in medium
  - T dependent
- Angle independent of associated  $p_T$ .



•PNJL Model



•Mikherjee, Mustafa, Ray

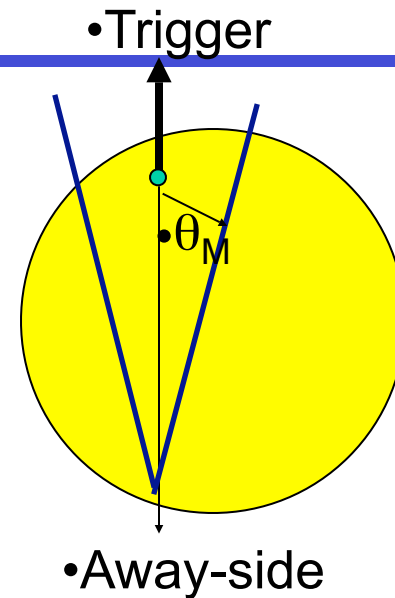
•Phys. Rev. D75 (2007) 094015

# Mach-Cone

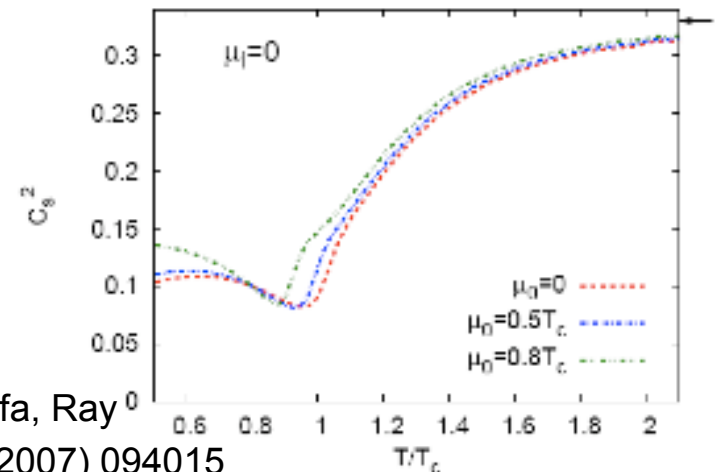
$$\frac{c_s}{v_{parton}} = \cos(\theta_M)$$

$$c_s^2 = \frac{\partial p}{\partial \epsilon}; \quad v_{parton} \approx c$$

- Mach angle depends on speed of sound in medium
  - T dependent
- Angle independent of associated  $p_T$ .



## •PNJL Model



•Mikherjee, Mustafa, Ray

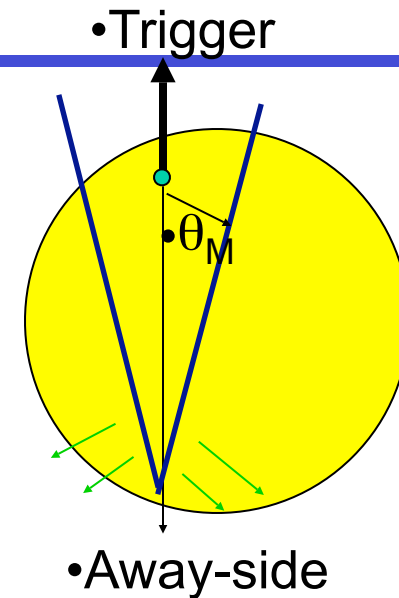
•Phys. Rev. D75 (2007) 094015

# Mach-Cone

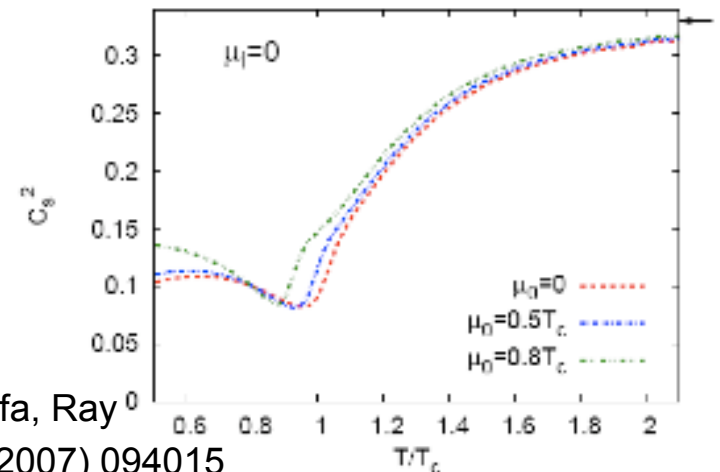
$$\frac{c_s}{v_{parton}} = \cos(\theta_M)$$

$$c_s^2 = \frac{\partial p}{\partial \epsilon}; \quad v_{parton} \approx c$$

- Mach angle depends on speed of sound in medium
  - T dependent
- Angle independent of associated  $p_T$ .



## •PNJL Model



•Mikherjee, Mustafa, Ray

•Phys. Rev. D75 (2007) 094015

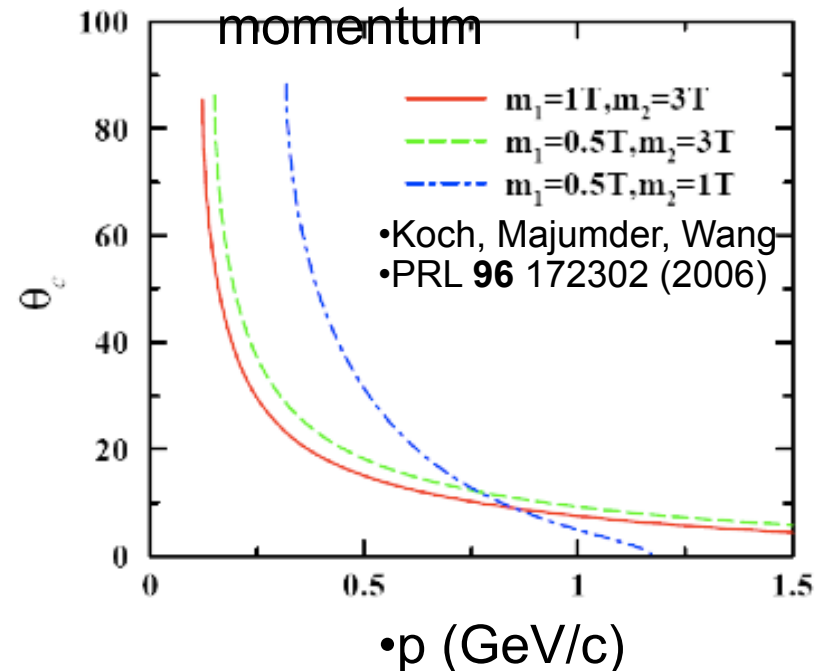


# Čerenkov Gluon Radiation

- Gluons radiated by **superluminal** partons.
- Angle is **dependent** on emitted **momentum**.

$$\frac{c_n}{v_{parton}} = \cos(\theta_c) = \frac{c}{n(p)v_{parton}} \approx \frac{1}{n(p)}$$

- Čerenkov angle vs emitted particle momentum



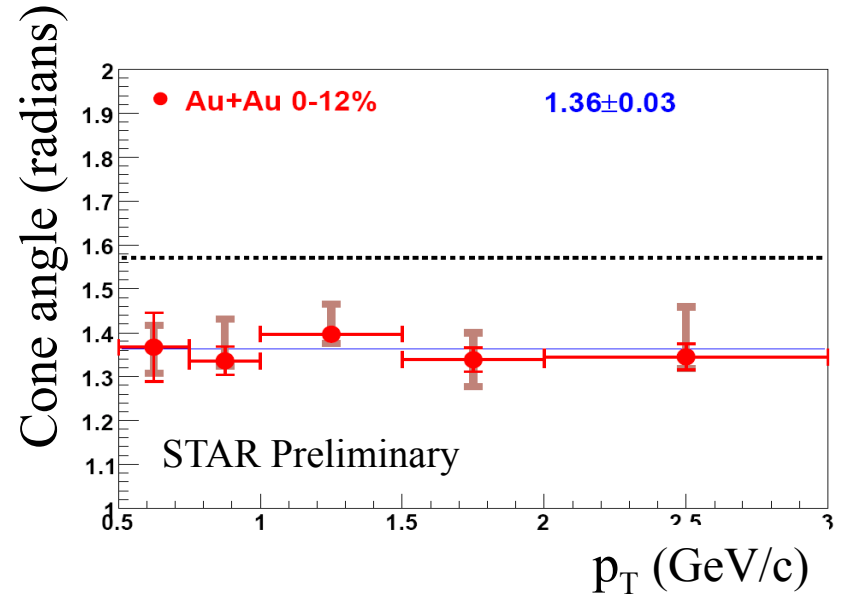
# Mach cone?

Naive calc. of time averaged velocity of sound in medium:

$$\frac{c_s}{v_{\text{parton}}} = \cos(\theta_M), v_{\text{parton}} = c$$

Cone angle  $\sim 1.36$  radians

$$c_s = 0.2c!$$



# Mach cone?

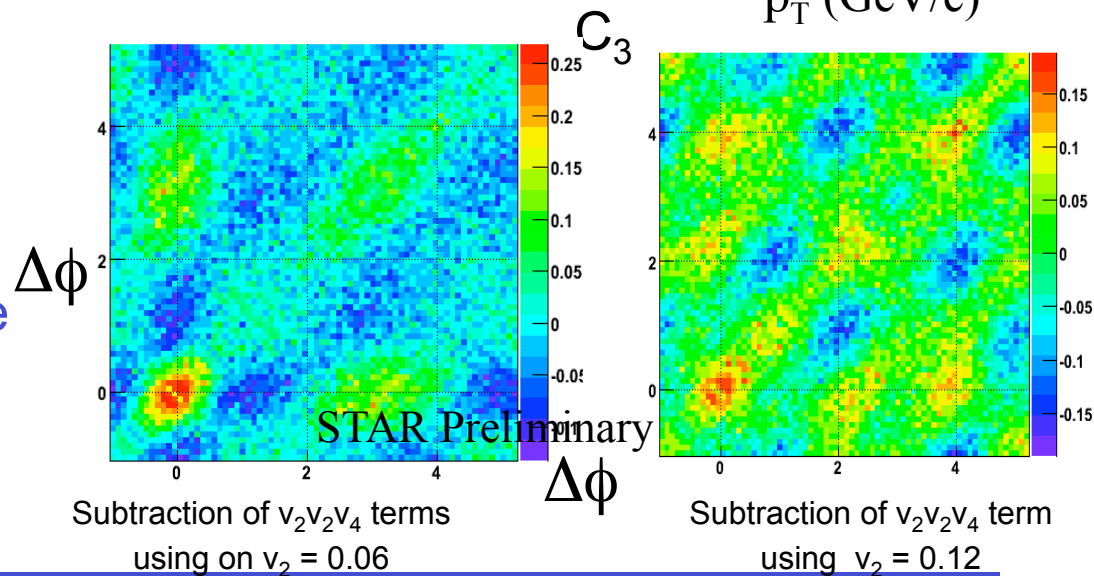
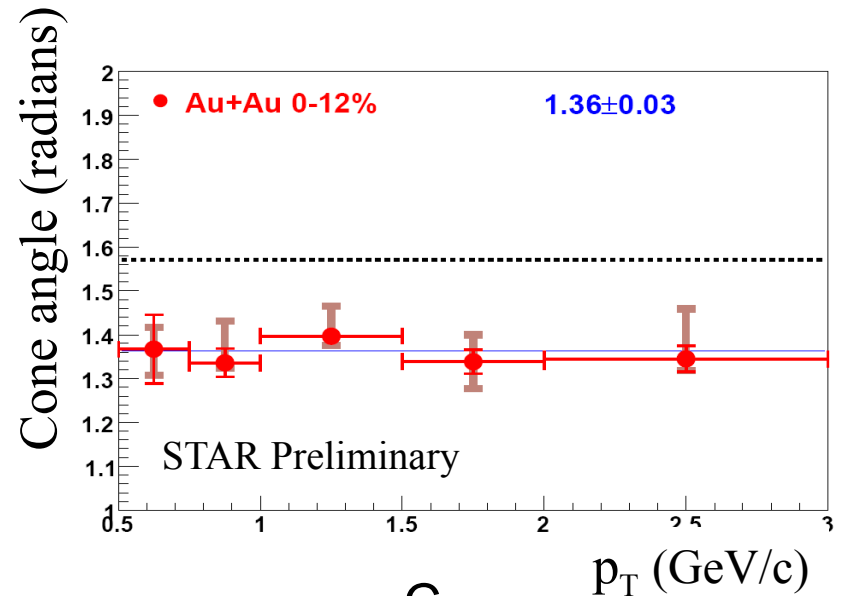
Naive calc. of time averaged velocity of sound in medium:

$$\frac{c_s}{v_{\text{parton}}} = \cos(\dot{\epsilon}_M), v_{\text{parton}} = c$$

Cone angle  $\sim 1.36$  radians

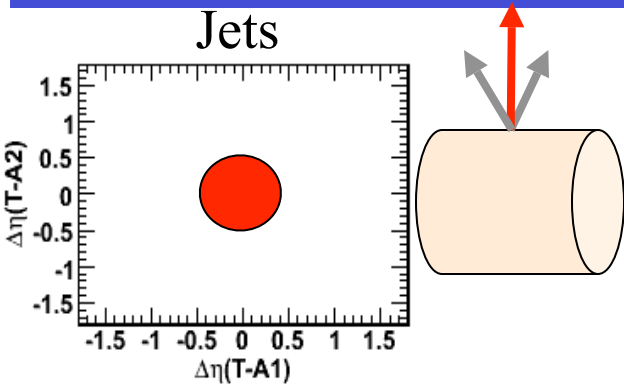
$$c_s = 0.2c!$$

- In cumulant approach: **no conclusive evidence for conical emission** so far
- Strength and shape of away side structures observed depends on assumed magnitude of flow coefficients**

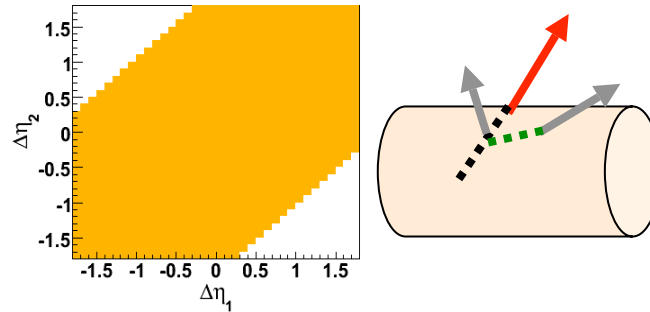


# Ridge : 3-particle Correlation

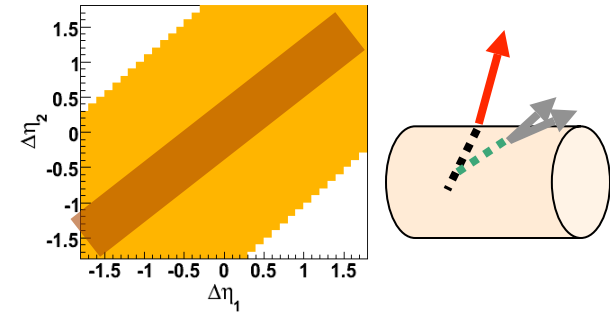
Jets



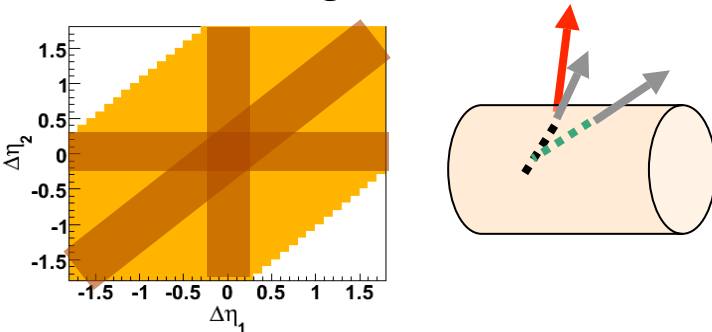
In-medium radiated  
gluons diffused in  $\eta$



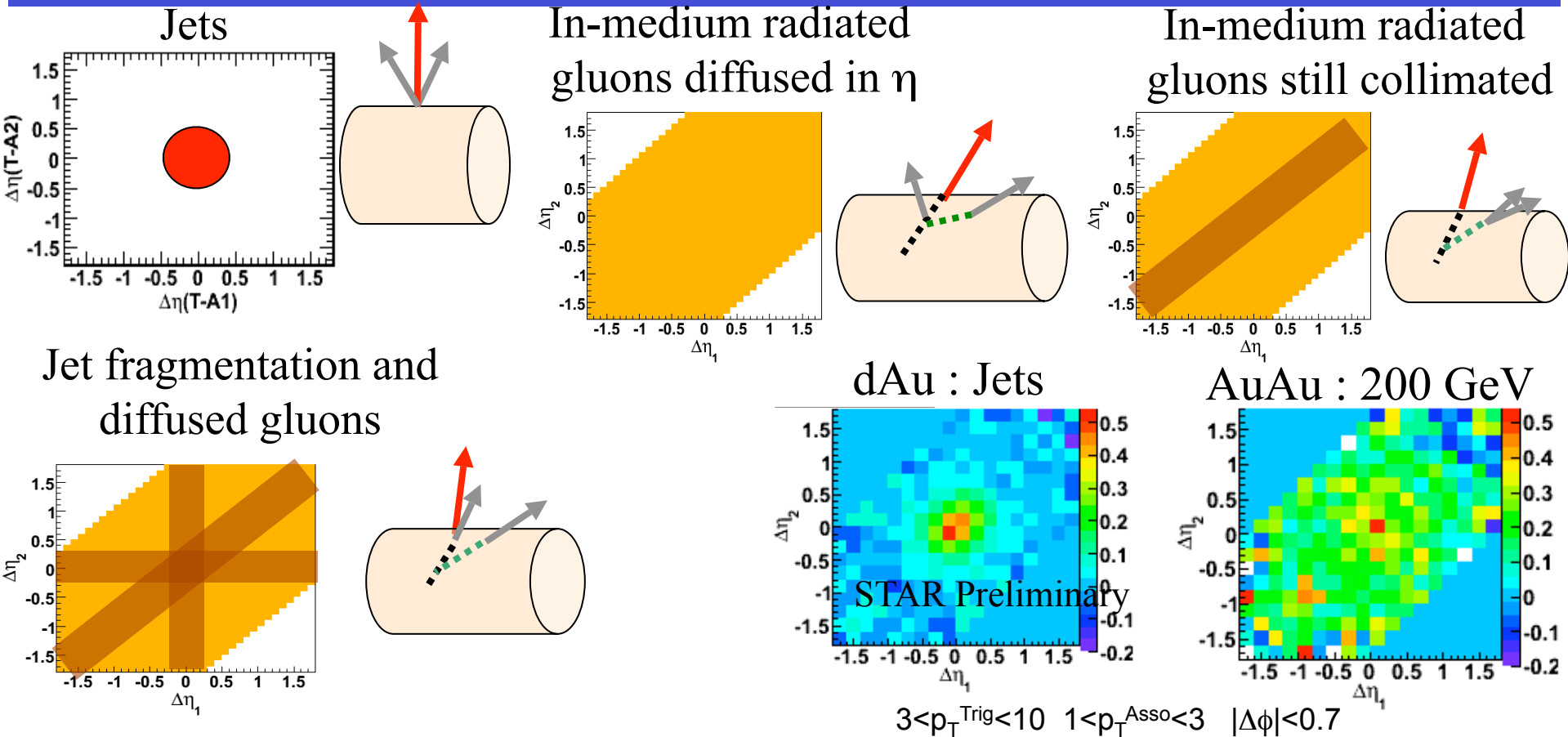
In-medium radiated  
gluons still collimated



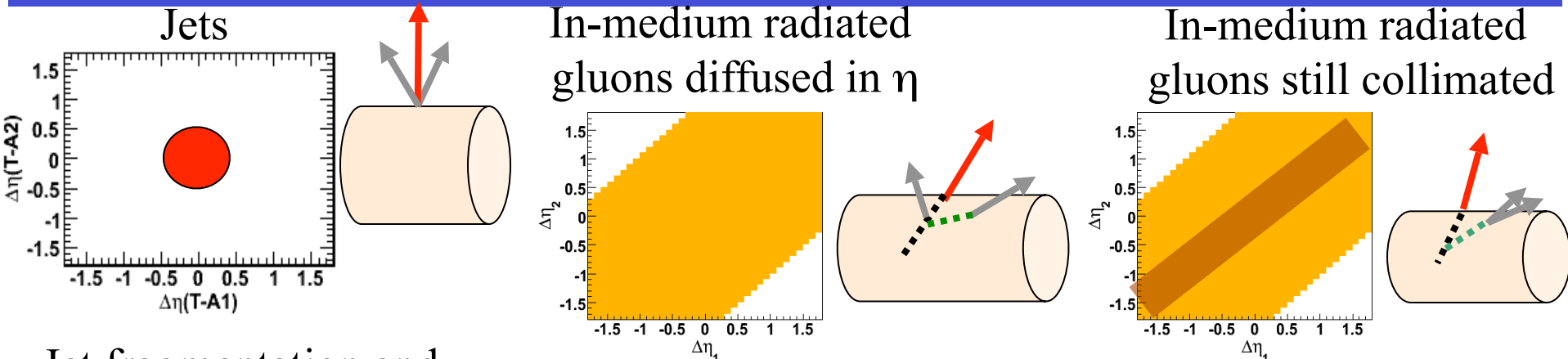
Jet fragmentation and  
diffused gluons



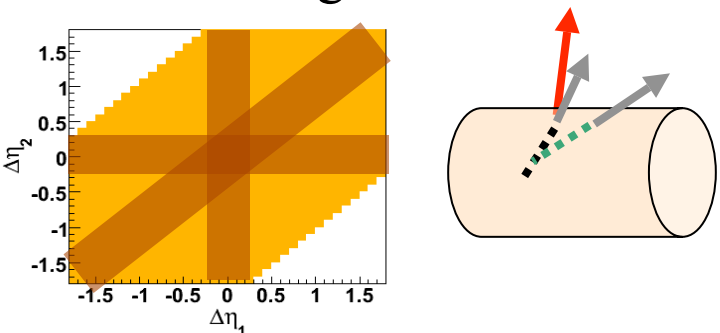
# Ridge : 3-particle Correlation



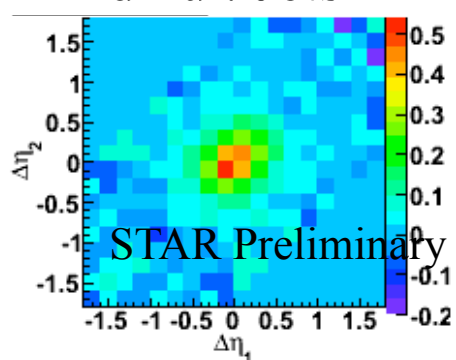
# Ridge : 3-particle Correlation



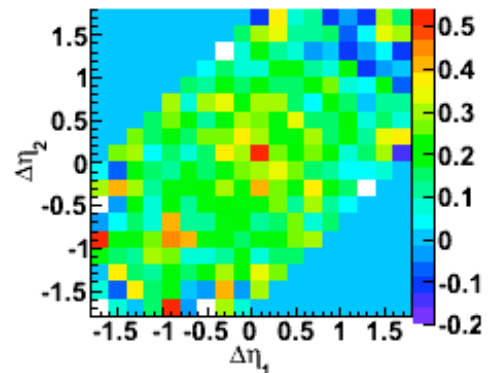
Jet fragmentation and diffused gluons



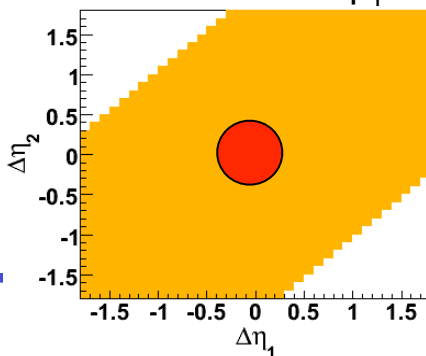
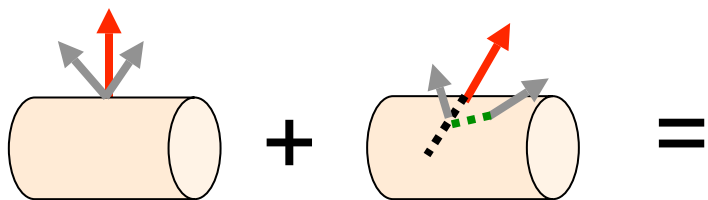
dAu : Jets



AuAu : 200 GeV

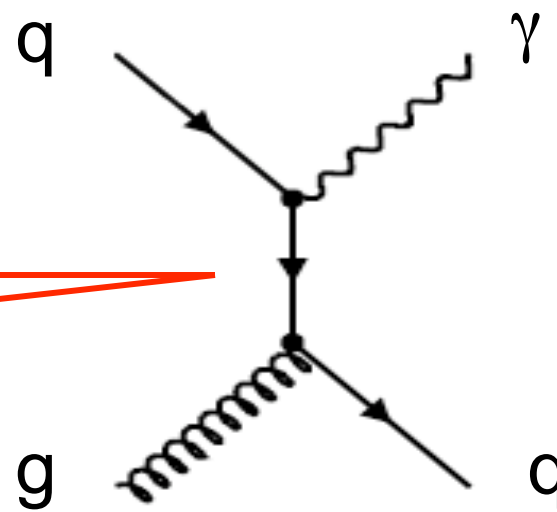
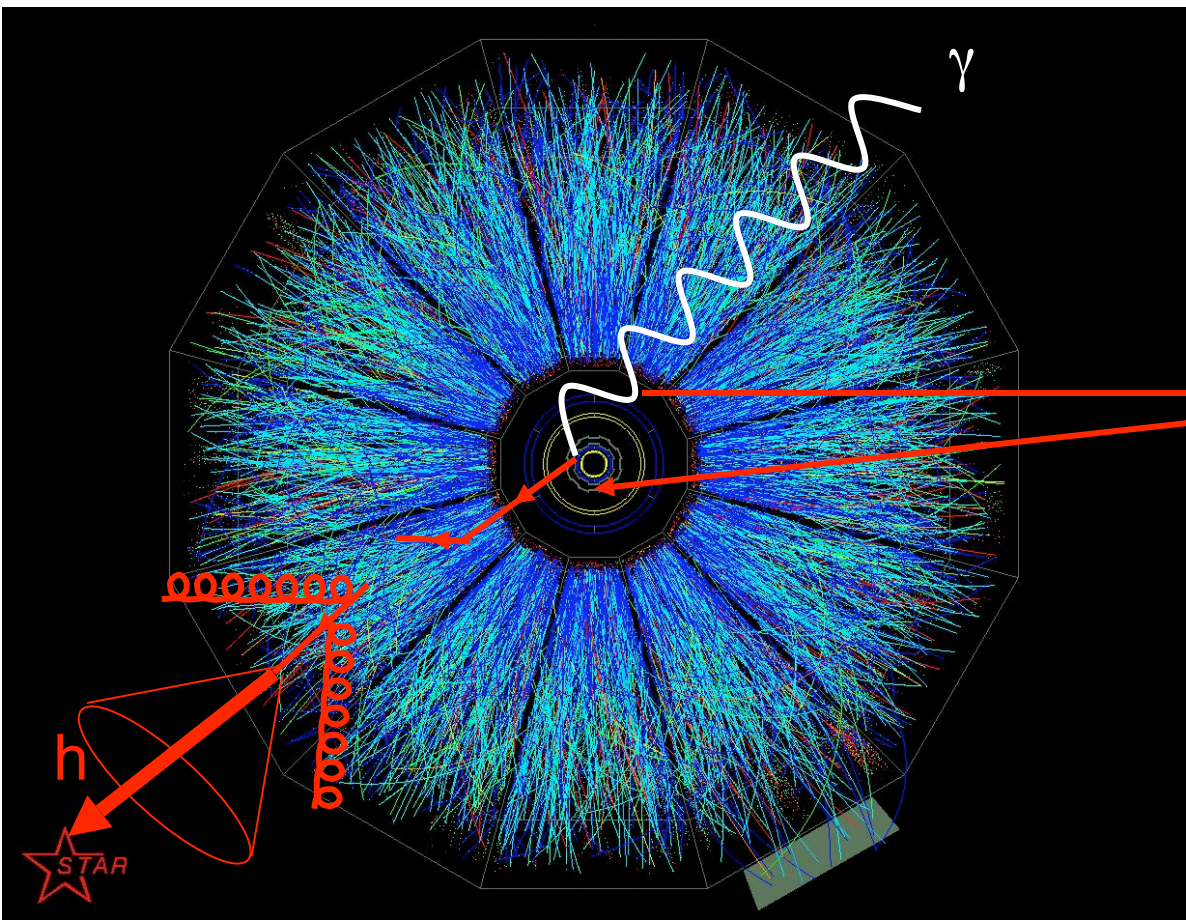


$$3 < p_T^{\text{Trig}} < 10 \quad 1 < p_T^{\text{Asso}} < 3 \quad |\Delta\phi| < 0.7$$



Uniform overall excess of associated particles not due to correlated emission

# Golden Probe of QCD Energy Loss - $\gamma$ -Jet

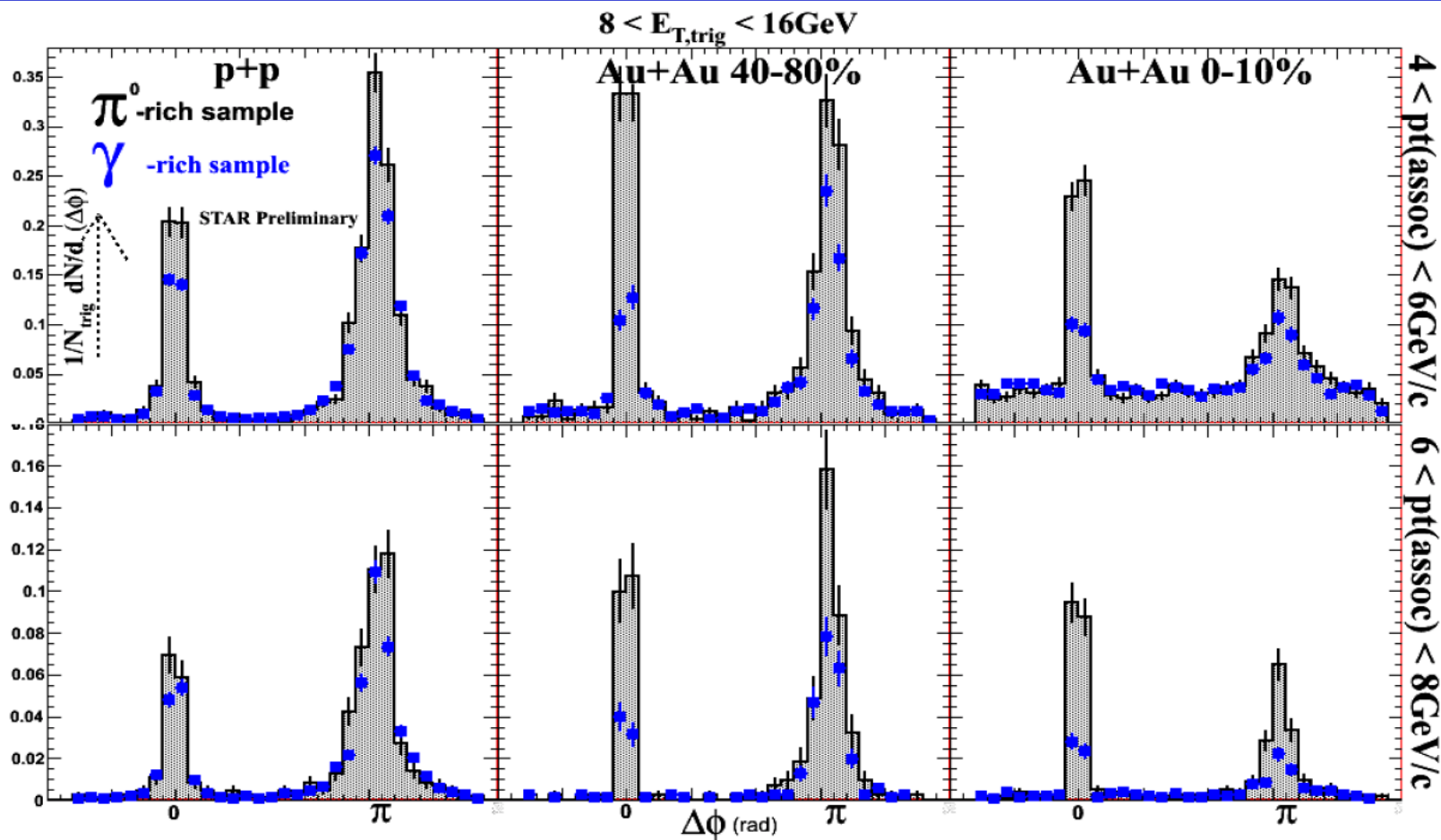


QCD analog of Compton Scattering

$\gamma$  emerges “unscathed” from medium

- Full reconstruction of kinematics: real fragmentation function ( $D(z)$ )

# $\gamma$ -hadron and $\pi^0$ -hadron correlations



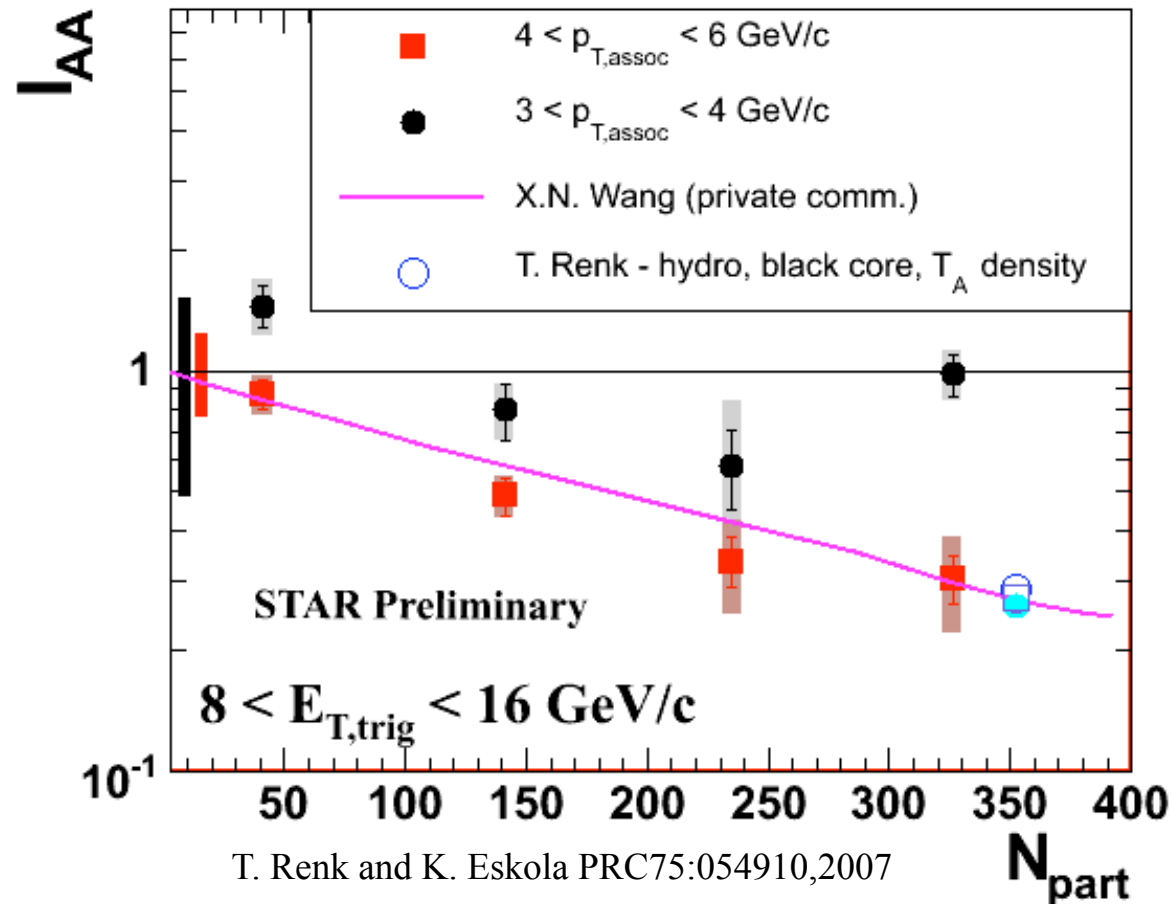
Shower shape in Shower Maximum Detector gives  $\gamma$ -,  $\pi^0$ -enriched samples

The  $\gamma$ -rich sample has lower near-side yield than  $\pi^0$ .



# First measure of away-side $I_{AA}$ for $\gamma$ -h

## (Direct) $\gamma$ triggers



$$E_{jet} = E_{\gamma} = E_{trig}$$

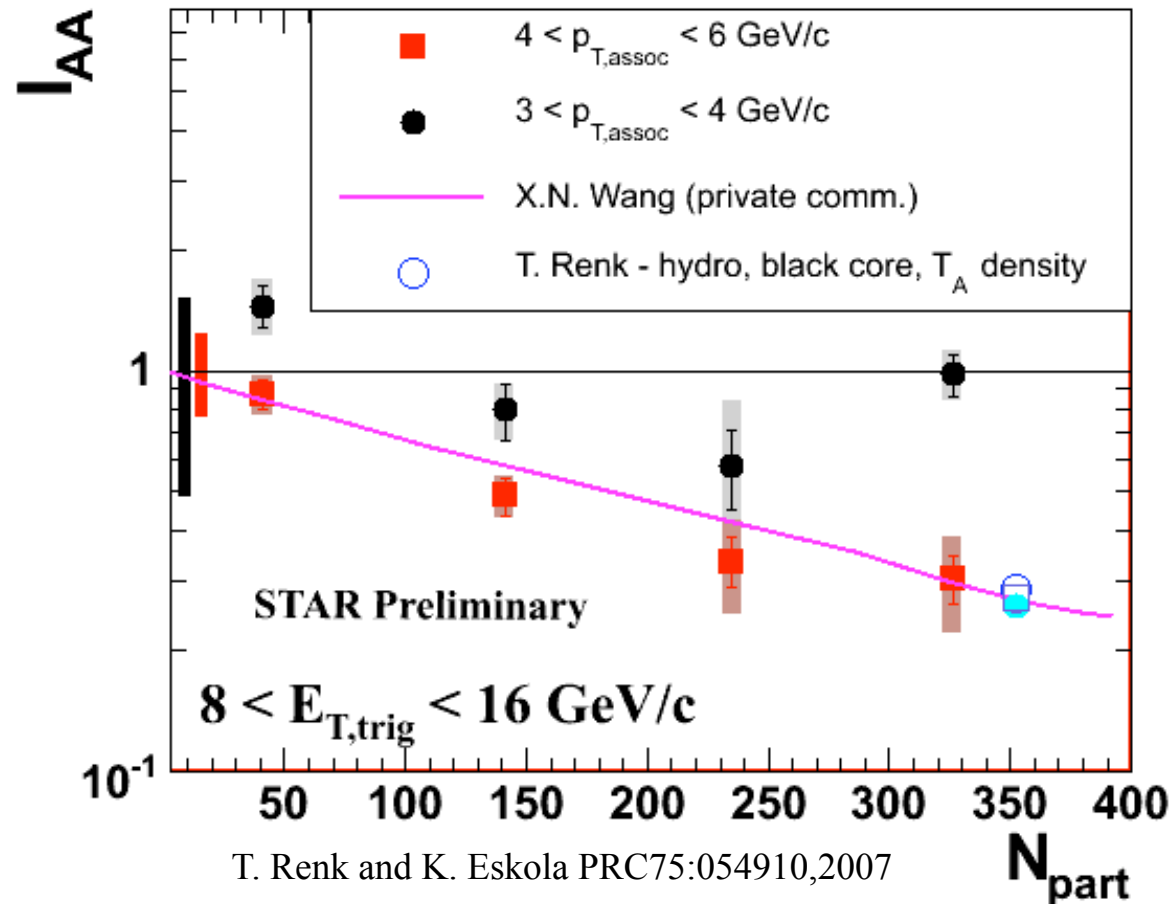
$$I_{AA} = \frac{D_{AA}(z_T, E_T^{trig})}{D_{pp}(z_T, E_T^{trig})}$$

$$D^{h_1 h_2}(z_T, p_T^{trig}) = p_T^{trig} \frac{d\sigma_{AA}^{h_1 h_2} / dp_T^{trig} dp_T}{d\sigma_{AA}^{h_1} / dp_T^{trig}}$$

Good agreement  
between theory  
and measurement  
for higher  $p_{T,assoc}$

# First measure of away-side $I_{AA}$ for $\gamma$ -h

## (Direct) $\gamma$ triggers



$$E_{jet} = E_{\gamma} = E_{trig}$$

$$I_{AA} = \frac{D_{AA}(z_T, E_T^{trig})}{D_{pp}(z_T, E_T^{trig})}$$

$$D^{h_1 h_2}(z_T, p_T^{trig}) = p_T^{trig} \frac{d\sigma_{AA}^{h_1 h_2} / dp_T^{trig} dp_T}{d\sigma_{AA}^{h_1} / dp_T^{trig}}$$

Good agreement  
between theory  
and measurement  
for higher  $p_{T,assoc}$

Suppression similar level to inclusives in central collisions